# Description of transfer reactions with coupled-channels Born approximation

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# Introduction

### Description of transfer reactions (conventional approach)

✓ The transition matrix for the A(a, b)B reaction within the **distorted-wave Born approximation (DWBA)**.

$$T_{\text{DWBA}} = \left\langle \Psi_{\beta}^{(-)} \middle| V_{xb} \middle| \Psi_{\alpha}^{(+)} \right\rangle$$

- ✓ The optical potential  $U_{aA}(U_{bB})$  for the a + A(b + B) 2-body system generates the distorted wave.
- ✓ One-step transition induced by the residual interaction  $V_{xb}$  ( $V_{xA}$ ) for the post (prior) form is assumed.

# Introduction

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✓ The transition matrix for the A(a, b)B reaction within the **distorted-wave Born approximation (DWBA)**.

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- ✓ The optical potential  $U_{(U_{a})}$  for the a + A(b + B) 2 body system generates the diagonal of C an DWBA describe the reaction
- ✓ **One-step tran** even if *a* or *B* is **loosely bound system**? for the post (prior) form is assumed.  $c_b(V_{xA})$

# Model

### Beyond DWBA

M. Kamimura *et al.*, Prog. Theor. Phys. Suppl. No. 89, 1 (1986).
N. Austern *et al.*, Phys. Rep. **154**, 125 (1987).
M. Yahiro *et al.*, Prog. Theor. Exp. Phys. **2012**, 01A209 (2012).

 ✓ Coupled-channels Born approximation (CCBA) with the continuum-discretized coupled-channels (CDCC) method.

$$T_{\text{CCBA}} = \left\langle \Psi_{\beta(\text{CDCC})}^{(-)} \middle| V_{xb} \middle| \Psi_{\alpha(\text{CDCC})}^{(+)} \right\rangle$$

- ✓ The optical potential  $U_{xA}(U_{bA})$  for the subsystem x + A(b + A) generates the distorted wave based on the 3-body model.
- ✓ The CDCC wave functions both in the initial and final channels.
  - $\rightarrow$  **Remnant term** is canceled out exactly.
  - $\rightarrow$  **Rearrangement component** is involved implicitly.

# Model

### Breakup process

 $\checkmark\,$  Decomposition of the transition matrix

$$T_{CCBA} = \underline{T}_{\beta(el),\alpha(el)} + \underline{T}_{\beta(el),\alpha(br)} + \underline{T}_{\beta(br),\alpha(el)} + \underline{T}_{\beta(br),\alpha(br)}$$

$$A = \begin{bmatrix} b & & & \\ b & & & \\ c &$$

# Model

### Breakup process

 BC is implicitly taken into account in DWBA as "absorption".

✓ Decomposition of the transition matrix
 ✓ BT is never involved in DWBA.

$$T_{\rm CCBA} = T_{\beta(\rm el),\alpha(\rm el)} + T_{\beta(\rm el),\alpha(\rm br)} + T_{\beta(\rm br),\alpha(\rm el)} + T_{\beta(\rm br),\alpha(\rm br)}$$



## Result 1 T. Fukui *et al.*, Phys. Rev. C 91, 014604 (2015).



- ✓ The <sup>8</sup>B(d, n)<sup>9</sup>C reaction is paid attention with astrophysical interest.
- ✓ Significant breakup effect (58%) can be seen at the forward angles of the angular distribution of the cross section.

#### Breakup effects of each path



 $\checkmark$  The BC is weak and the ET result can be regarded as that of DWBA.

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- ✓ Strong interferences between the ET and the BT in each channel enhance the cross section. → Never involved in DWBA.

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- ✓ Strong interferences between the ET and the BT in each channel enhance the cross section. → Never involved in DWBA.
- $\checkmark$  The BT among continuum states is negligible.

#### Dynamical change of transferred angular momentum *l*



### **Breakup effect on** $S_{18}$ **of** $^{8}B(p, \gamma)^{9}C$



### Future work

- (1) Inclusion of the 3-body configuration in  ${}^{9}C (p + p + {}^{7}Be)$ .
- (2) The CCBA analysis of the mirror reaction  ${}^{8}\text{Li}(d, p){}^{9}\text{Li}$ .

[12] B. Guo et al., Nucl. Phys. A761, 162 (2005).

#### In the second second



N. Anantaraman *et al.*, Nucl. Phys. **A313**, 445 (1979). F. D. Becchetti *et al.*, Nucl. Phys. **A303**, 313 (1978).

#### In <u>16O(6Li, d)<sup>20</sup>Ne(g.s.) to search surface manifestation of cluster</u>



#### Improvement

- (1) Diffraction pattern of the 1<sup>st</sup> and 2<sup>nd</sup> peaks
- (2) Reasonable values of the normalization factors
  - $\rightarrow$  Governed by reliabilities of both the  $\alpha$ -<sup>16</sup>O WF and OMP



N. Anantaraman *et al.*, Nucl. Phys. **A313**, 445 (1979). F. D. Becchetti *et al.*, Nucl. Phys. **A303**, 313 (1978).

#### Breakup effects of <sup>6</sup>Li



 Decomposition of the CDCC distorted wave into elastic and breakup channels.

$$\chi_{ ext{cDCC}}(oldsymbol{r}_i) = \chi_0(oldsymbol{r}_i) + \chi_c(oldsymbol{r}_i)$$

- ✓ Full ~ Elastic transfer (ET) ≠ No back coupling (BC)
  - → Breakup transfer (BT) is negligible. Only the BC (CC due to off-diagonal potentials) is essential.

$$\begin{array}{ccc} K_i + U_{00} - E_0 & U_{0c} \\ U_{c0} & K_i + U_{cc} - E_c \end{array} \begin{pmatrix} \chi_0 \\ \chi_c \end{pmatrix} = 0$$

→ DWBA can provide reasonable results, if an appropriate <sup>6</sup>Li-OMP, in which BC is implicitly taken into account as its imaginary part, is given.

T. Fukui *et al.*, Prog. Theor. Phys. 125, 1193 (2011)T. Fukui *et al.*, Phys. Rev. C **91**, 014604 (2015).

# Summary

#### CCBA analyses



Why is the breakup effect large?

Why opposite?

→ Explained in detail in T. Fukui *et al.*, Phys. Rev. C 91, 014604 (2015).

# Future work

### Iransfer reaction to unbound state (ex. ${}^{4}\text{He}(d, p){}^{5}\text{He}$ )

✓ The transition matrix of the post-form representation for (d, p) reaction

$$T_{\text{DWBA}}^{(\text{post})} = \left\langle \chi_{\beta}^{(-)} \psi_n \middle| \psi_d \chi_{\alpha}^{(+)} \right\rangle$$

$$= \int d\boldsymbol{r}_{\alpha} \int d\boldsymbol{r}_d \chi_{\beta}^{*(-)}(\boldsymbol{r}_{\alpha}, \boldsymbol{r}_d) \psi_n^*(\boldsymbol{r}_{\alpha}, \boldsymbol{r}_d) V_{pn}(\boldsymbol{r}_d) \psi_d(\boldsymbol{r}_d) \chi_{\alpha}^{(+)}(\boldsymbol{r}_{\alpha}).$$

$$= \int d\boldsymbol{r}_{\alpha} \int d\boldsymbol{r}_d \chi_{\beta}^{*(-)}(\boldsymbol{r}_{\alpha}, \boldsymbol{r}_d) \psi_n^*(\boldsymbol{r}_{\alpha}, \boldsymbol{r}_d) V_{pn}(\boldsymbol{r}_d) \psi_d(\boldsymbol{r}_d) \chi_{\alpha}^{(+)}(\boldsymbol{r}_{\alpha}).$$

$$= \int d\boldsymbol{r}_{\alpha} \int d\boldsymbol{r}_d \chi_{\beta}^{*(-)}(\boldsymbol{r}_{\alpha}, \boldsymbol{r}_d) \psi_n^*(\boldsymbol{r}_{\alpha}, \boldsymbol{r}_d) V_{pn}(\boldsymbol{r}_d) \psi_d(\boldsymbol{r}_d) \chi_{\alpha}^{(+)}(\boldsymbol{r}_{\alpha}).$$

#### ✓ The prior form

$$T_{\text{DWBA}}^{(\text{prior})} = \left\langle \tilde{\chi}_{\beta}^{(-)} \psi_{n} \middle| V_{n\alpha} \middle| \psi_{d} \chi_{\alpha}^{(+)} \right\rangle$$

$$p = \int d\mathbf{r}_{\alpha} \int d\mathbf{r}_{d} \tilde{\chi}_{\beta}^{*(-)}(\mathbf{r}_{\alpha}, \mathbf{r}_{d}) \psi_{n}^{*}(\mathbf{r}_{\alpha}, \mathbf{r}_{d}) V_{n\alpha}(\mathbf{r}_{\alpha}, \mathbf{r}_{d}) \psi_{d}(\mathbf{r}_{d}) \chi_{\alpha}^{(+)}(\mathbf{r}_{\alpha}).$$
oscillate attenuate attenuate
$$\mathbf{r}_{d} \quad \mathbf{r}_{\beta} \quad \mathbf{r}_{p}$$

$$\mathbf{r}_{a} \quad \mathbf{r}_{n}$$

$$5\text{He} \quad \alpha$$

# Future work

n

### **Transfer reaction to unbound state (ex.** <sup>4</sup>**He**(*d*, *p*)<sup>5</sup>**He**)

✓ The transition matrix of the post-form representation for  $(d \ p)$  reaction

 $T_{\rm DW}^{\rm (p)}$  The distorted wave  $\tilde{\chi}_{\beta}^{(-)}$  should be exact.

 $\rightarrow \text{The CCBA approach is necessary for the final channel.} \begin{array}{c} t \text{ converge.} \\ T_{\text{DWBA}}^{(\text{prior})} \rightarrow T_{\text{CCBA}}^{(\text{prior})} & \stackrel{+)}{}_{\chi}(\boldsymbol{r}_{\alpha}). \end{array}$ 



$$= \int d\boldsymbol{r}_{\alpha} \int d\boldsymbol{r}_{d} \tilde{\chi}_{\beta}^{*(-)}(\boldsymbol{r}_{\alpha}, \boldsymbol{r}_{d}) \psi_{n}^{*}(\boldsymbol{r}_{\alpha}, \boldsymbol{r}_{d}) V_{n\alpha}(\boldsymbol{r}_{\alpha}, \boldsymbol{r}_{d}) \psi_{d}(\boldsymbol{r}_{d}) \chi_{\alpha}^{(+)}(\boldsymbol{r}_{\alpha})$$



oscillate attenuate attenuate

These respectively attenuate for two independent coordinates. → The integration does converge.