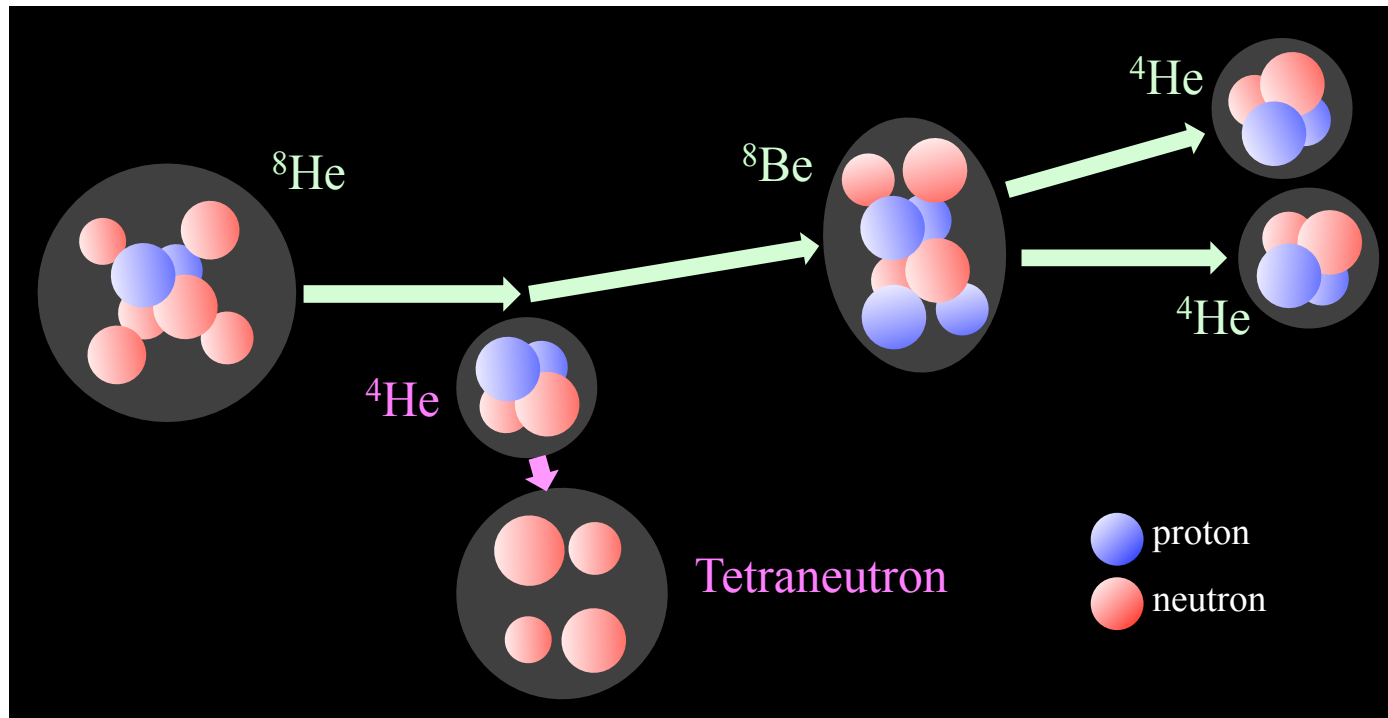




Tetra-neutron states populated by $^4\text{He}(^8\text{He}, ^8\text{Be})$ reaction

Exotic system populated by exotic reaction



S. Shimoura

CNS, University of Tokyo

K. Kisamori et al., Phys. Rev. Lett. 116, 052501



Tetra-neutron

- Multi-neutron System
 - Neutron cluster (?) in fragmentation of ^{14}Be
PRC65, 044006 (2002)
 - NN, NNN, NNNN interactions
 - Neutron-Neutron interaction
 - T=3/2 NNN force
 - > 3-body force in neutron matter
 - Ab initio type calculations
 - Multi-body resonances
 - Correlations in multi-fermion scattering / resonant (?) states



Historical Review

~ search for a bound state of 4n~

1960s

fission of Uranium

- No evidence for particle stable state of tetra-neutron

J. P. Shiffer Phys. Lett. 5, 4, 292 (1963)

1980s

$^4\text{He}(\pi^-, \pi^+)$ reaction

- Only upper limit of cross section was decided.

J. E. Unger, et al., Phys. Lett. B 144, 333 (1984)

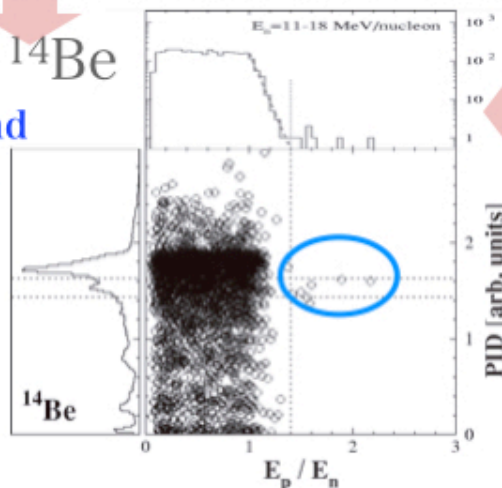
Bound state: No clear evidence.

2000s

Breakup of ^{14}Be

- Candidates of **bound tetra-neutron** were observed.

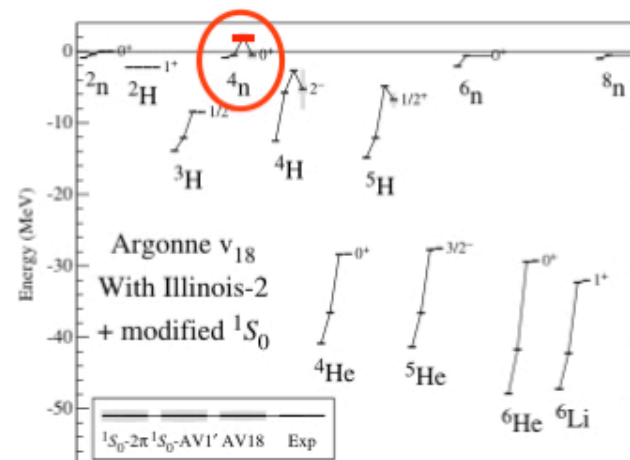
F. M. Marques, et al,
Phys. Rev. C 65,
044006 (2002)



2000s

Theoretical work

- ab-initio calculation
NN, NNN interaction



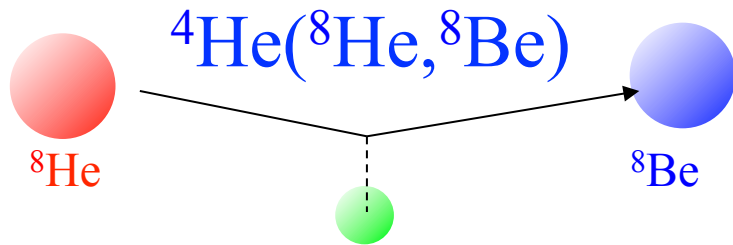
S. C. Piper, Phys. Rev. Lett. 90, 252501 (2003)

- **Bound $4n$ cannot exist**
- **Possible resonance state ~ 2 MeV**

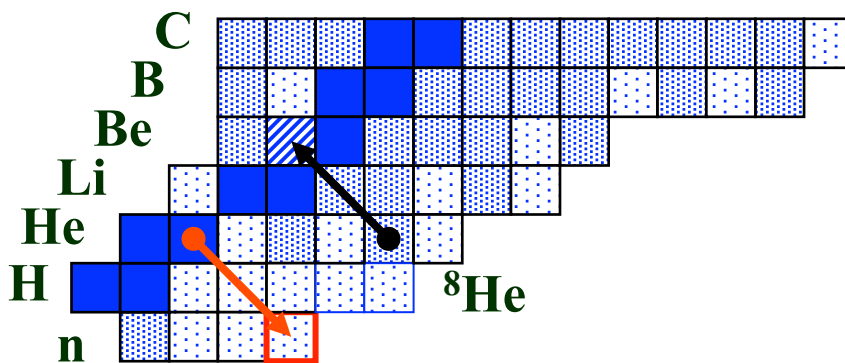
Resonance state : Possibility of the state is still an open and fascinating question.



Tetra-neutron system produced by exothermic double-charge exchange reaction

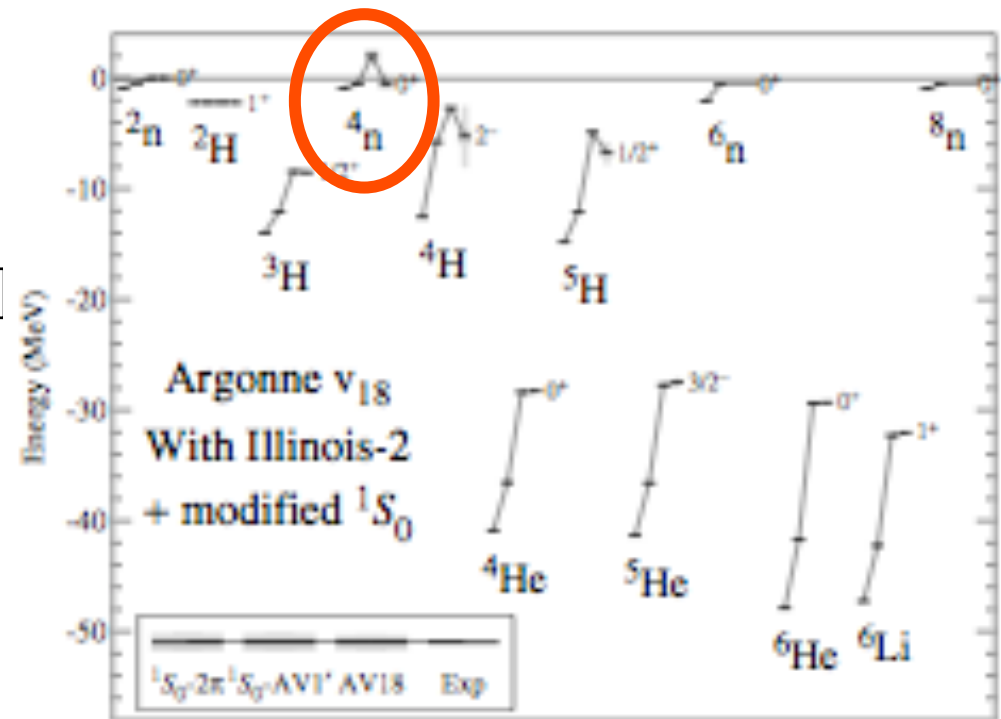


Recoil-less 4n system via DCX using internal energy of ${}^8\text{He}$



Multi-Neutron

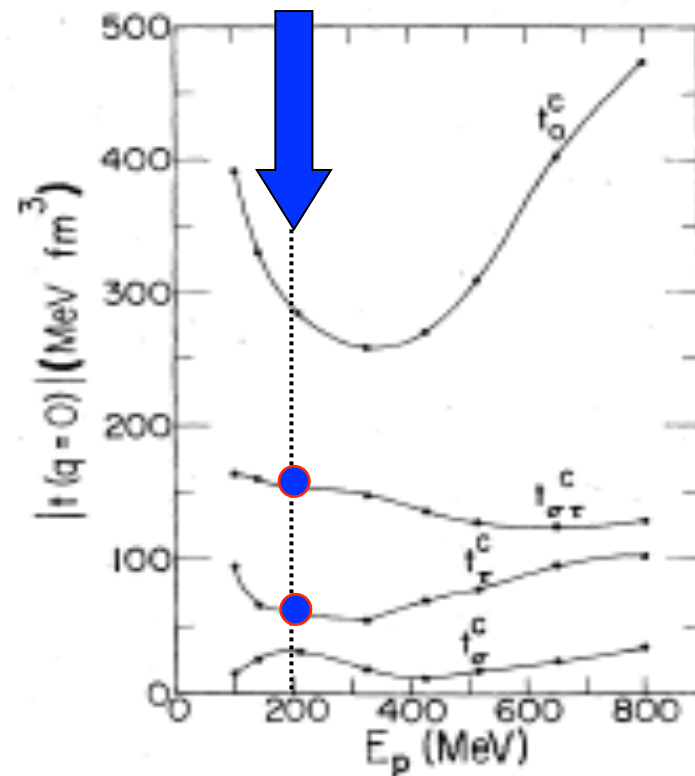
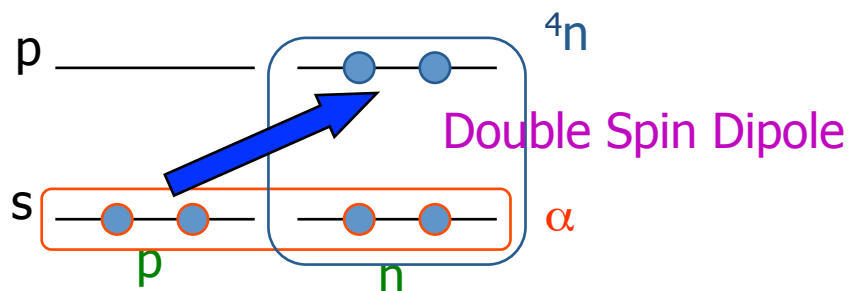
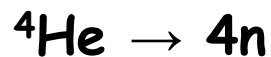
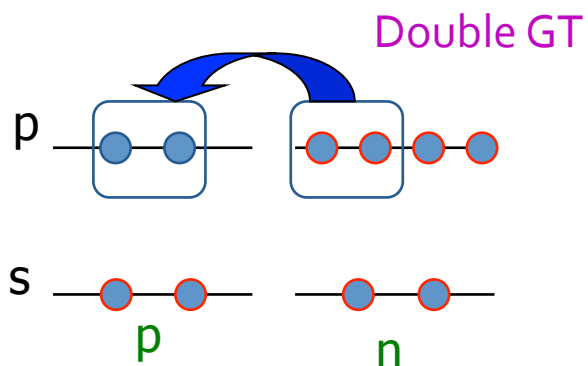
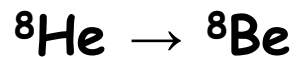
Almost recoil-less condition with ${}^4\text{He}({}^8\text{He}, {}^8\text{Be})4n$ reaction around 200 A MeV



S.C. Pieper et al., PRL 90, 252501 (2003)



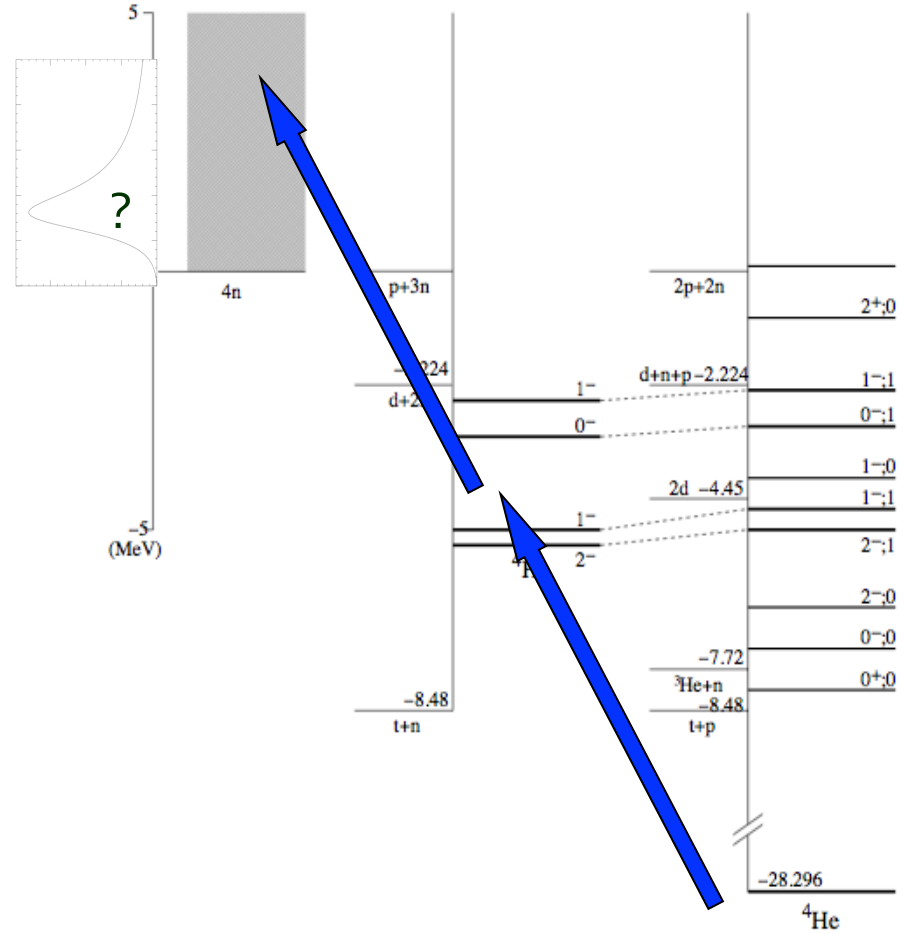
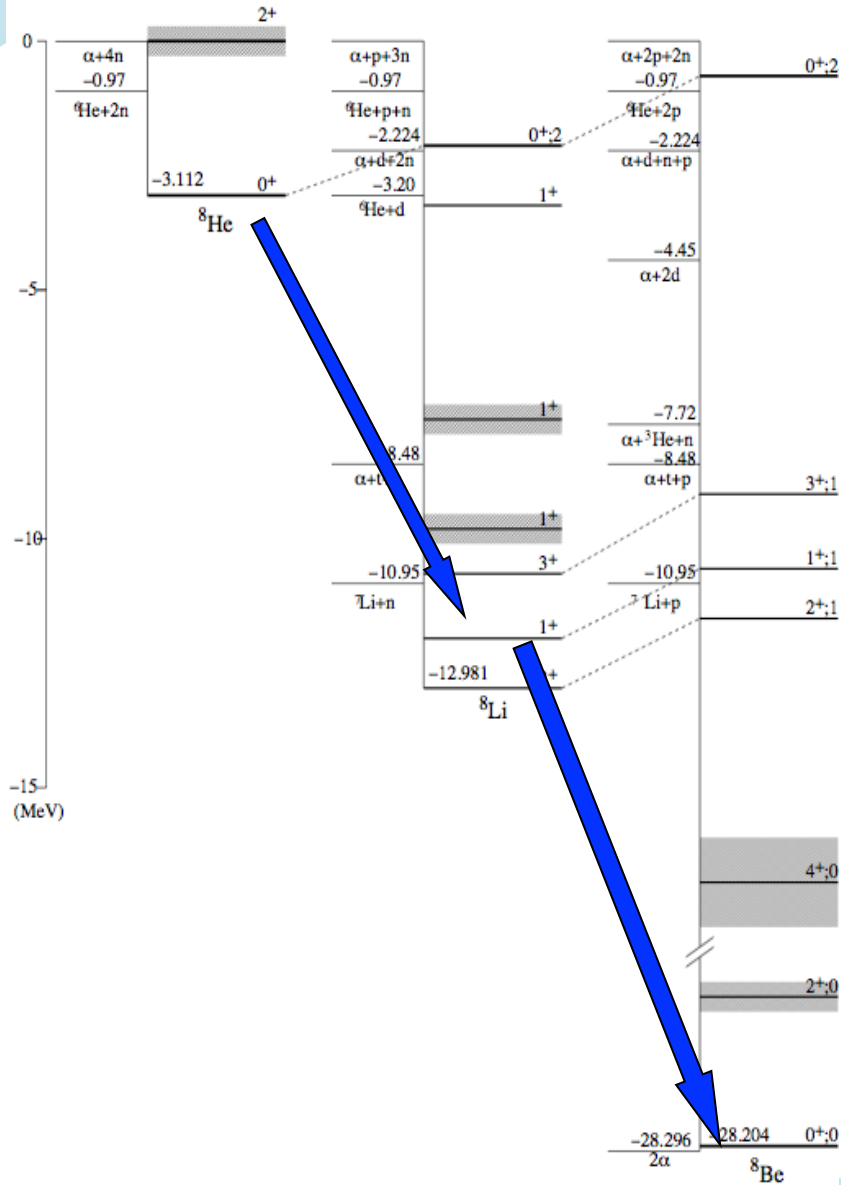
Reaction Mechanism



$$\left[(\vec{\tau}_p \cdot \vec{\tau}_t) (\vec{\sigma}_p \cdot \vec{\sigma}_t) r_t Y_1(\hat{r}_t) \right]^2$$



Level diagrams

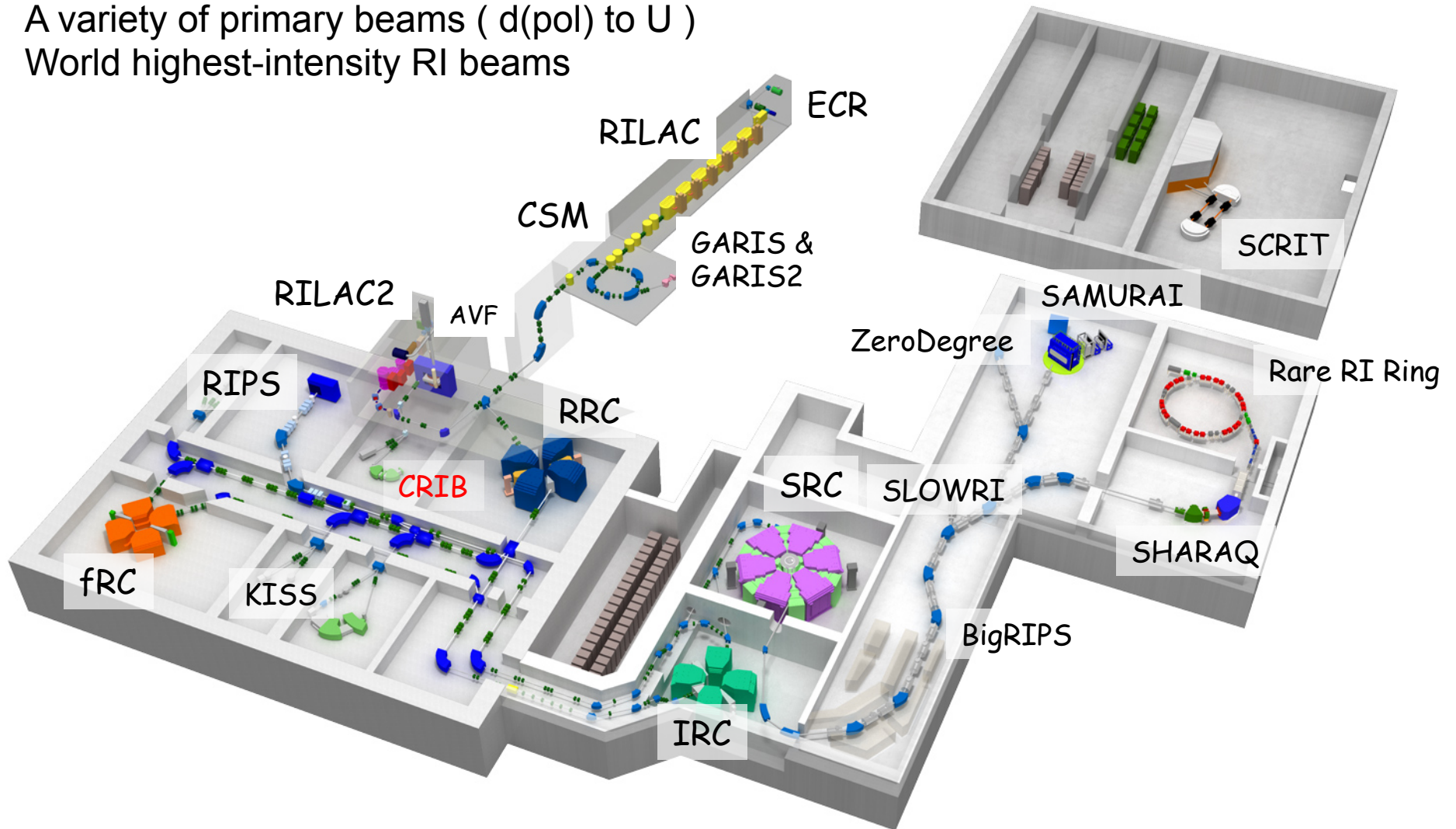


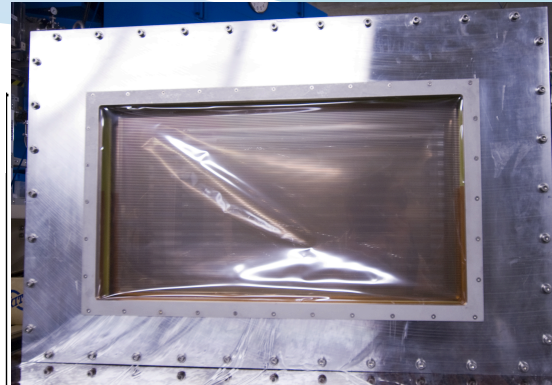
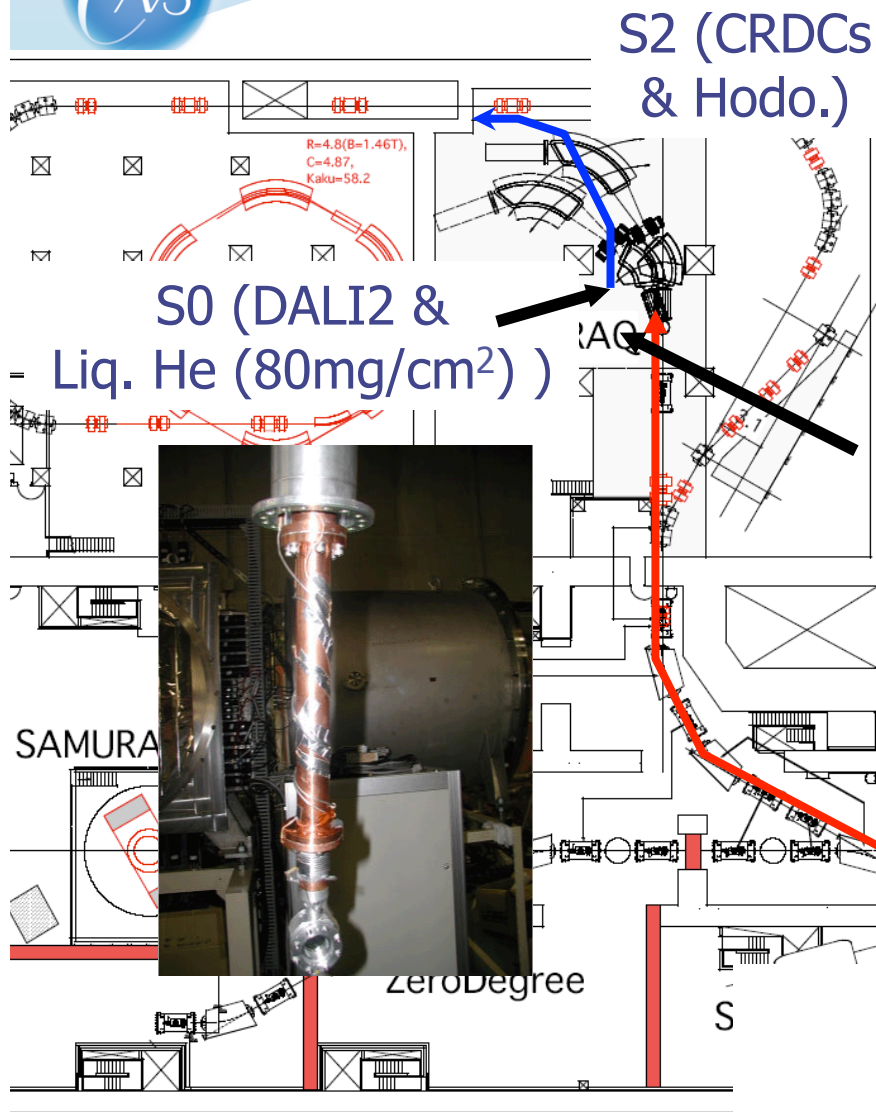
$q_{\min} \sim 10 \text{ MeV}/c$

RI Beam Factory at RIKEN

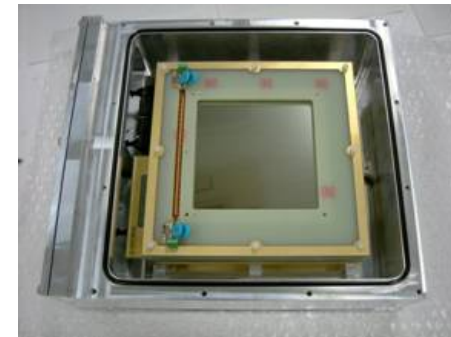
3 injectors + cascade of 4 cyclotrons
⇒ several to 345 MeV/nucleon

A variety of primary beams (d(pol) to U)
World highest-intensity RI beams





FH10 (MWDCs)



F6 (MWDCs)
mom.
tagging

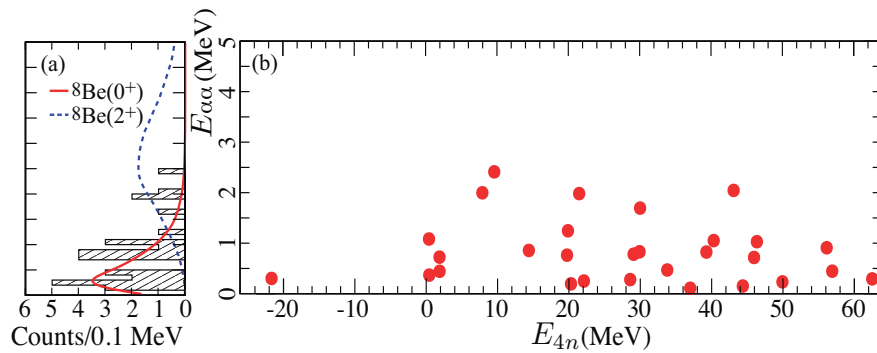
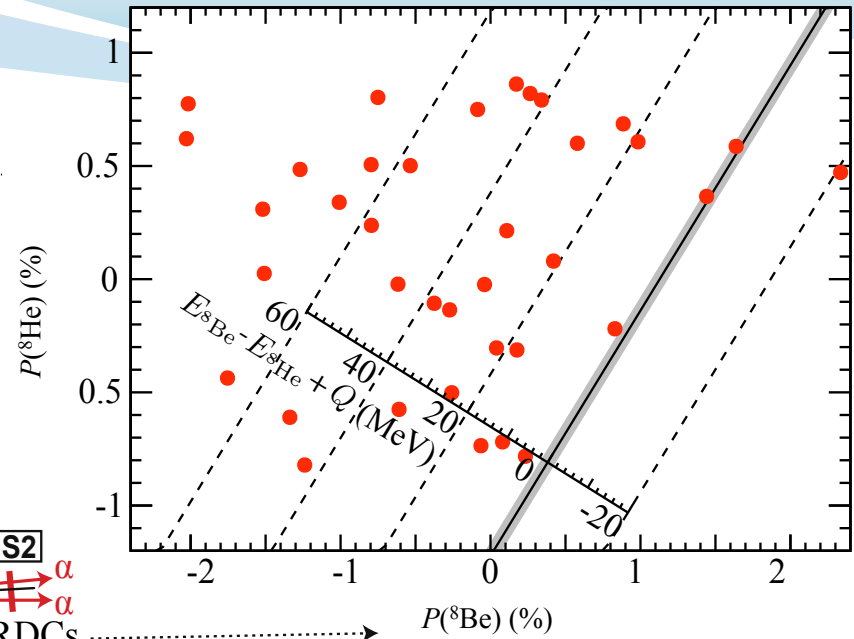
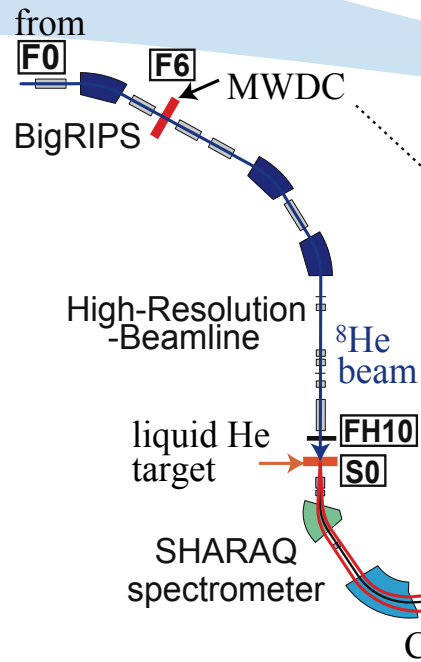
⁸He beam (8.4Tm)
(190 ± 4) A MeV,
2 MHz from 400 pA of ¹⁸O

l. Tgt.)

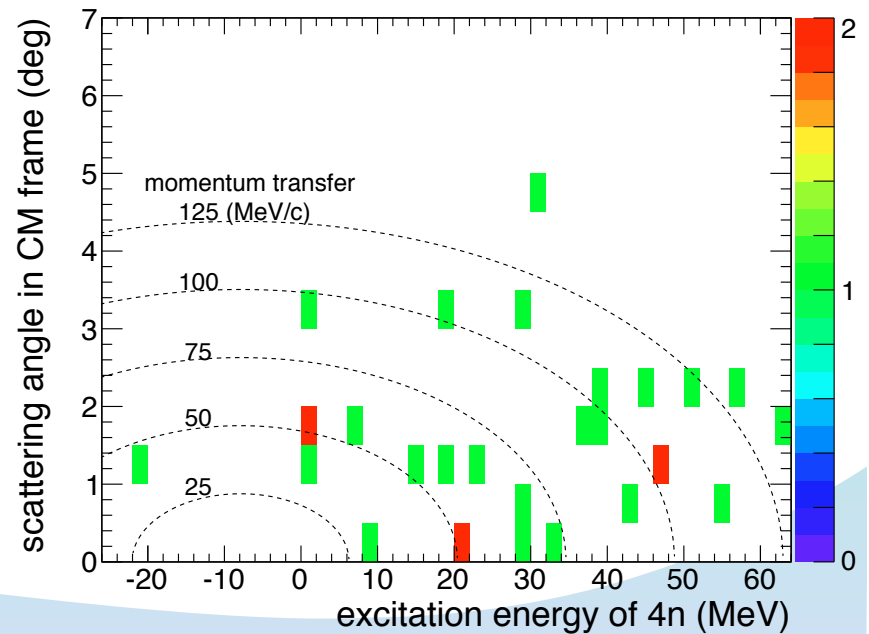
Experimental setup



Experimental Results

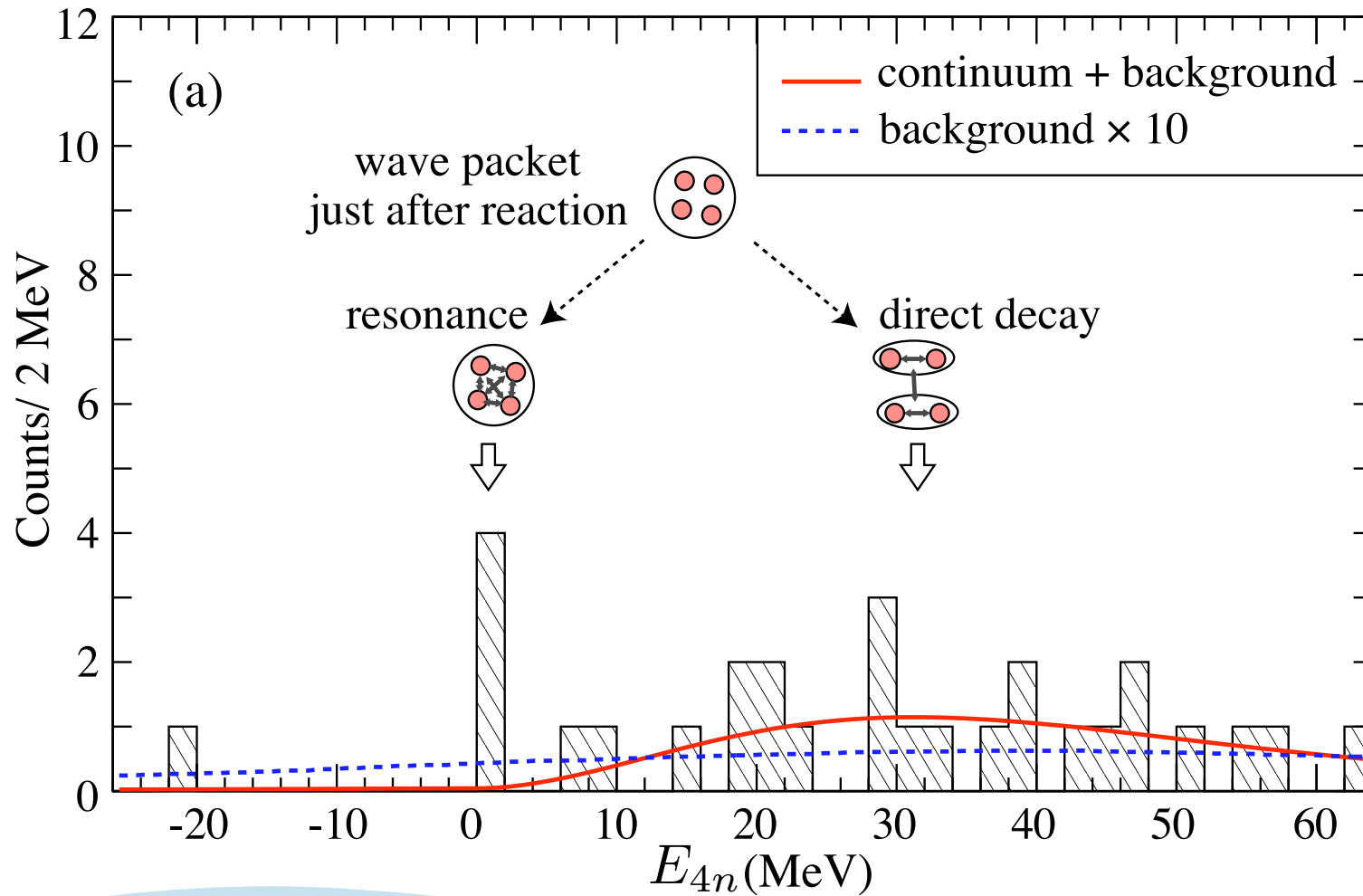


Acceptance for ${}^8\text{Be}(2^+)$ was 13 % of that for ${}^8\text{Be}(0^+)$
A few events could be from ${}^8\text{Be}(2^+)$.





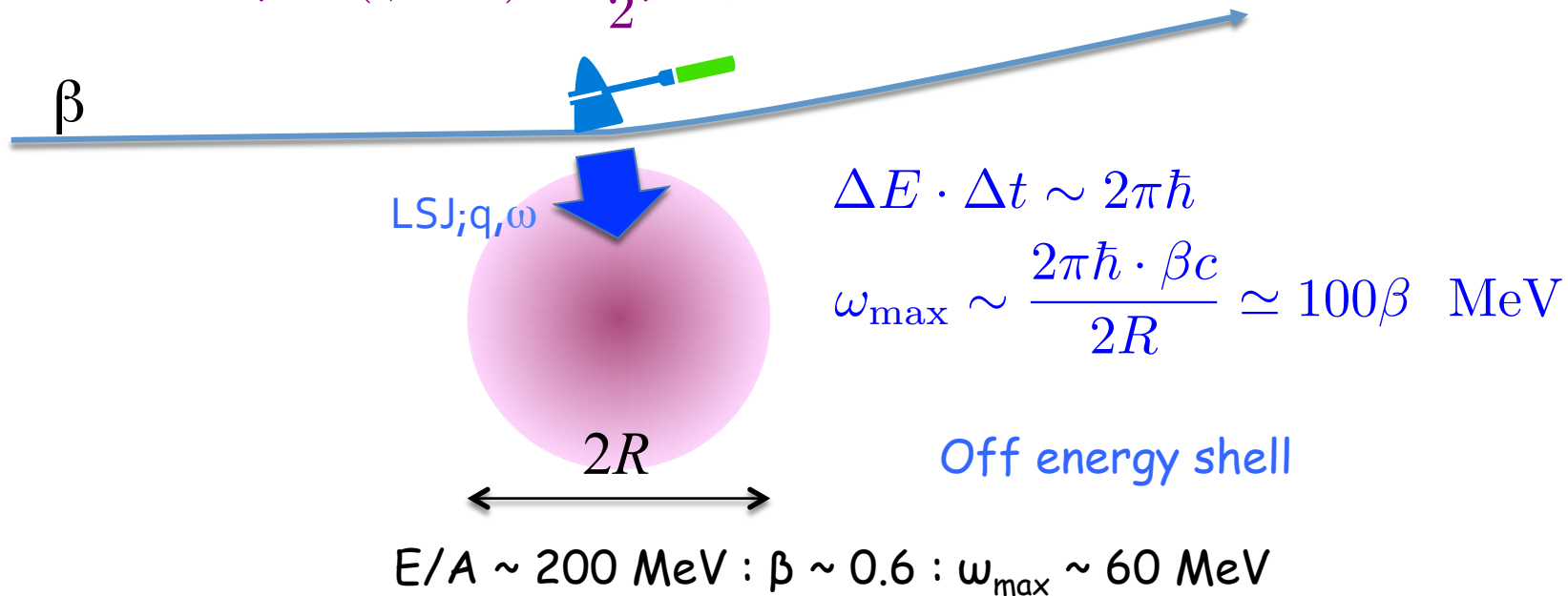
Experimental Results





Reaction time & excitation energy for intermediate-energy “inelastic-type scattering”

$$\omega \ll \mu c^2 (\gamma - 1) \simeq \frac{1}{2} \mu c^2 \beta^2$$



$$O(ls j \tau; \xi) |E_i J_i \pi_i T_i; \xi_i\rangle = \int M_{if}(E_f) |E_f J_f \pi_f T_f; \xi_f\rangle \text{ Response}$$

$|M_{if}(E_f)|^2$: Energy Spectrum



Direct Part

c.f. Continuum spectrum with n-n FSI

L.V. Grigorenko, N.K. Timofeyuk, M.V. Zhukov, Eur. Phys. J. A 19, 187 (2004)

$$\mathcal{A}\Phi_0(\mathbf{r}_{12}, \mathbf{r}_{34}, \mathbf{r}_\alpha) \sim$$

$$\left[\left(\frac{r_{12}^2}{a^2} - \frac{3}{2} \right) - \left(\frac{r^2}{a^2} - \frac{3}{4} \right) \right] \exp \left[-\frac{r^2}{a^2} - \frac{r_{12}^2}{2a^2} - \frac{r_{34}^2}{2a^2} \right] \chi(1,2)\chi(3,4)$$

$$\left[\left(\frac{r_\alpha^2}{(a/\sqrt{2})^2} - \frac{3}{2} \right) - \frac{2\vec{r}_{12} \cdot \vec{r}_{34}}{a^2} \right] \exp \left[-\frac{r_\alpha^2}{a^2} - \frac{r_{12}^2}{2a^2} - \frac{r_{34}^2}{2a^2} \right] \chi(1,3)\chi(4,2)$$

$$\left[\left(\frac{r_\alpha^2}{(a/\sqrt{2})^2} - \frac{3}{2} \right) + \frac{2\vec{r}_{12} \cdot \vec{r}_{34}}{a^2} \right] \exp \left[-\frac{r_\alpha^2}{a^2} - \frac{r_{12}^2}{2a^2} - \frac{r_{34}^2}{2a^2} \right] \chi(1,4)\chi(2,3)$$

$$\vec{r}_\alpha = \frac{\vec{r}_1 + \vec{r}_2}{2} - \frac{\vec{r}_3 + \vec{r}_4}{2} \quad \chi(i,j) = \frac{1}{\sqrt{2}} (\uparrow(i)\downarrow(j) - \downarrow(i)\uparrow(j))$$

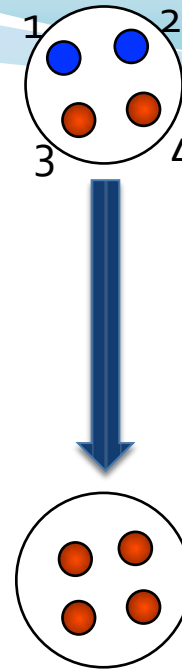


Fourier Transform: $\mathcal{A}\Phi_0(\mathbf{r}_{12}, \mathbf{r}_{34}, \mathbf{r}_\alpha) \rightarrow \mathcal{A}\tilde{\Phi}_0(\mathbf{k}_{12}, \mathbf{k}_{34}, \mathbf{k})$

$$\int |\mathcal{A}\tilde{\Phi}_0|^2 d^3k d^3k_{12} d^3k_{34} \delta(E - \epsilon - \epsilon_{12} - \epsilon_{34}) \propto X^{11/2} \exp(-X)$$

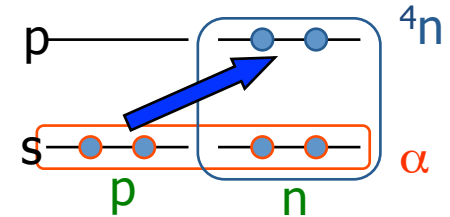
Peak at $X = 11/2$; $E \sim 60$ MeV

$$X = E/\epsilon_a \quad \epsilon_a = \frac{\hbar^2}{m_N a^2} = 11 \text{ MeV}$$



${}^4\text{He} \sim \Phi[(0s)^4]$

DCX



$q \ll 200$ MeV/c

4n wave packet just after DCX

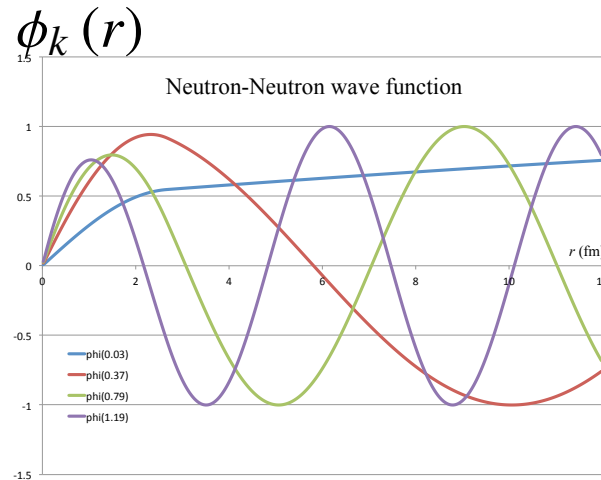
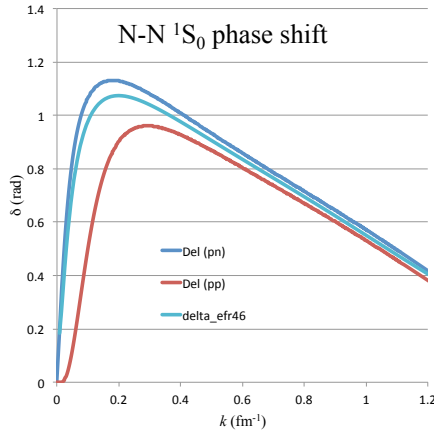
$\Phi_0 \sim \mathbf{r}_1 \cdot \mathbf{r}_2 \Phi[(0s)^4]$

Fourier Transform is expansion with plane waves \Rightarrow correlated scattering waves for FSI



NN FSI

Continuum spectrum with n-n FSI



Density of State

$$D_{ns}(\epsilon_{nn}) = \frac{|\hat{A}_{ns}(k)|^2}{k} \quad (\text{for } n = 1, 2) ; \quad \epsilon_{nn} = \frac{\hbar^2 k^2}{m_N}$$

$$\hat{A}_{1s}(k) = \int_0^\infty dr r \psi_{1s}(r) \phi_k(r) = 2 \left(\frac{1}{\sqrt{\pi} a^3} \right)^{1/2} k A_{1s}(k)$$

$$\hat{A}_{2s}(k) = \int_0^\infty dr r \psi_{2s}(r) \phi_k(r) = 2 \sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{\pi} a^3} \right)^{1/2} k A_{2s}(k)$$

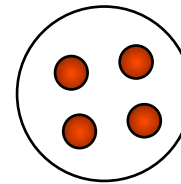
Expand $\mathcal{A}\Phi_0$ with correlated n-n scattering wave $\phi_k(r)$
 $A(k)$'s are used instead of Fourier component



${}^4\text{He} \sim \Phi[(0s)^4]$

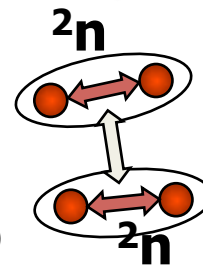
DCX

$q > 15 \text{ MeV}/c$



4n wave packet just after DCX

$\Phi_0 \sim r_1 \cdot r_2 \Phi[(0s)^4]$



Two correlated neutron pairs with weakly correlated

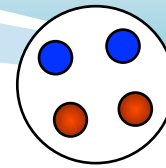
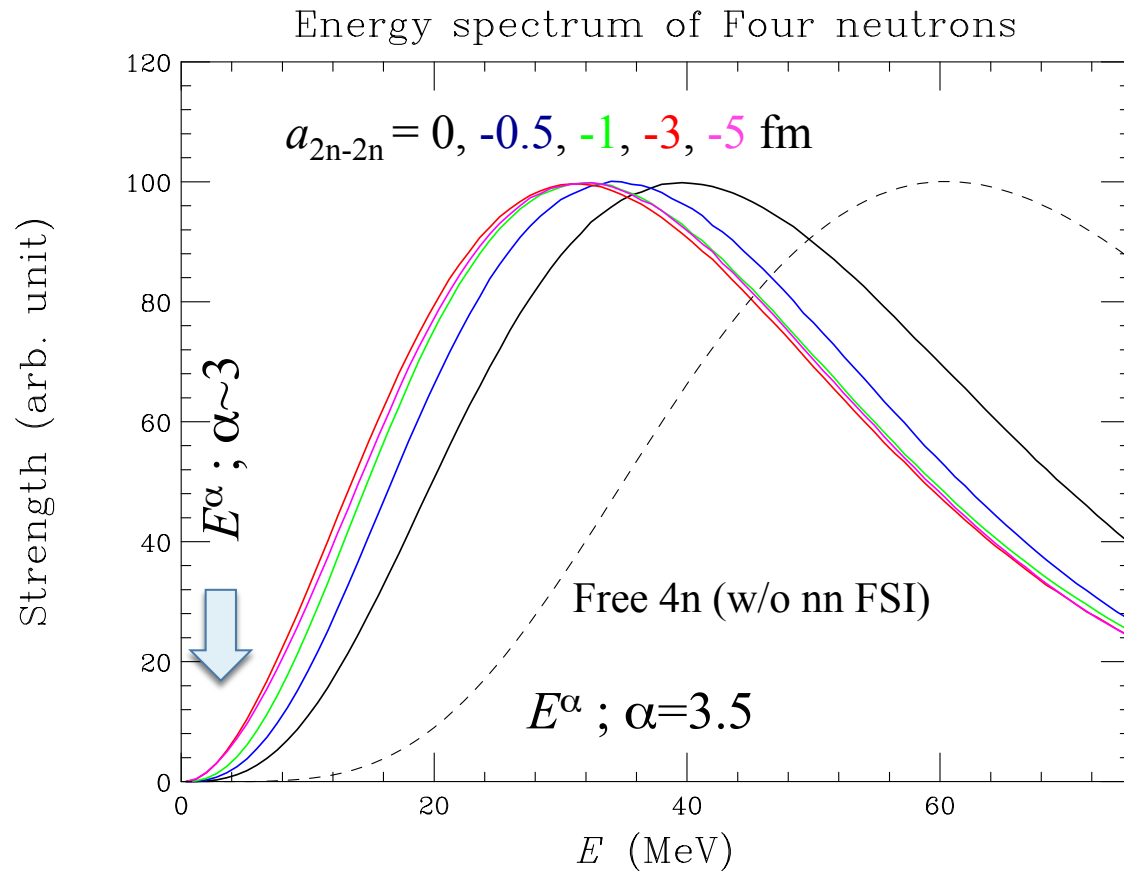


Direct Part

Continuum spectrum with n-n FSI

c.f.

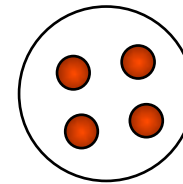
L.V. Grigorenko, N.K. Timofeyuk, M.V. Zhukov, Eur. Phys. J. A 19, 187 (2004)



${}^4\text{He} \sim \Phi[(0s)^4]$

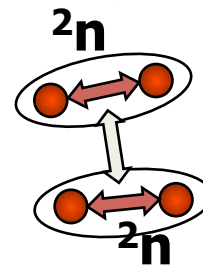
DCX

$q \ll 200$ MeV/c



4n wave packet just after DCX

$\Phi_0 \sim r_1 \cdot r_2 \Phi[(0s)^4]$

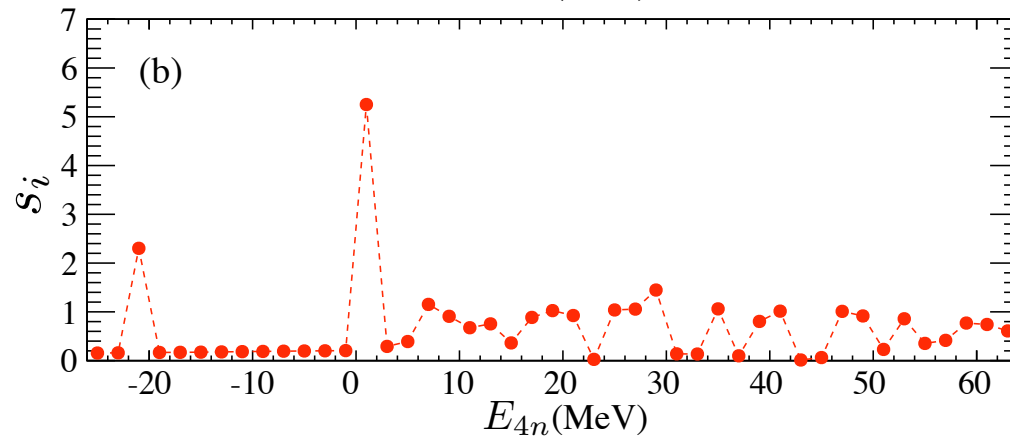
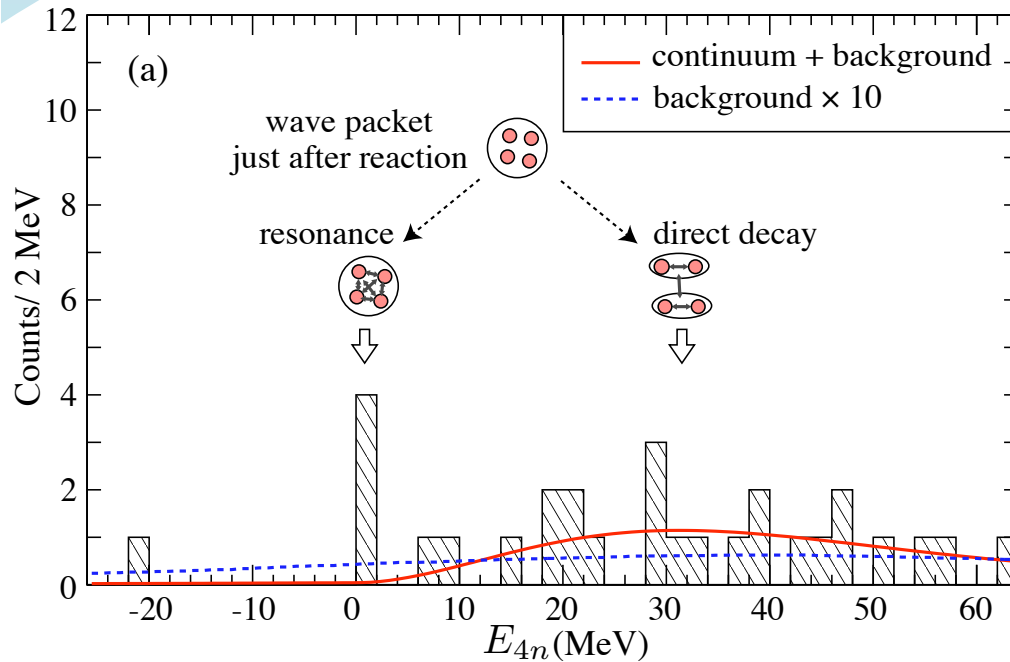


Two correlated neutron pairs with weakly correlated

Correlation is taking into account for 2n-2n relative motion by using scattering length



Fit with direct component & BG



Energy spectrum is expressed by the continuum from the direct decay and (small) experimental background except for four events at $0 < E_{4n} < 2$ MeV

The Four events suggest a possible resonance at $0.83 \pm 0.65(\text{stat.}) \pm 1.25(\text{sys.})$ MeV with width narrower than 2.6 MeV (FWHM). [4.9 σ significance]

Integ. cross section $\theta_{\text{cm}} < 5.4\text{deg}$: $3.8^{+2.9}_{-1.8}$ nb

• likelihood ratio test

$$\chi^2_\lambda = -2 \ln [L(\mathbf{y}; \mathbf{n}) / L(\mathbf{n}; \mathbf{n})]$$

• Significance:

$$s_i = \sqrt{2[y_i - n_i + n_i \ln(n_i/y_i)]}$$

n_i : num. of events in the i -th bin

y_i : trial function in the i -th bin

• Look Elsewhere Effect

$$\mu^n e^{-\mu} / n! \simeq 10^{-6} \text{ for } \mu = 0.07, n = 4$$



Summary

- ${}^4\text{He}({}^8\text{He}, {}^8\text{Be})4n$ has been measured at 190 A MeV at RIBF-SHARAQ
- Missing mass spectrum with very few background
- Although statistics is low (27 evs), spectrum looks two components (continuum + peak)
- Continuum is consistent with direct breakup process from $(0s)^2(0p)^2$ wave packet
- Four events just above 4n threshold is statistically beyond prediction of continuum + background (4.9σ significance)
 - candidate of 4n resonance
 - at $0.83 \pm 0.65(\text{stat.}) \pm 1.25(\text{sys.})$ MeV; $\Gamma < 2.6$ MeV
- Constraint to nucleon forces : n-n two body; T=3/2 three-body force; non-central; off-energy shell ...