



Neutrinos, Dark Matter and the Standard Model

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Outline

Neutrinos, Dark Matter and the Standard Model

C. Francisco,
T. Duchesne,
T.Oliveira,
L. Santos,
C. Almeida

Neutrinos

The God Particle

Discovery of a Higgs Boson

CERN's Large Hadron Collider

The LHC Tunnel

Cost Statistical Analyses

Simulation of Collision

Decay Channels of the Higgs

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- 2 The God Particle
- 3 Discovery of a Higgs Boson
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WNPPC 2018

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Winter Nuclear and Particle Physics Conference (WNPPC)



Mont Tremblant, 15-18 February 2018 - Québec - Canada



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The 3 types of neutrinos: “Les enfants terribles”.

 Symmetry magazine
4/2 às 10:12 - 

Detector detectives: find out how detectors in Fermilab's Short-Baseline Neutrino Program are searching for the mysterious fourth neutrino flavor.



Sterile neutrino sleuths

Meet the detectors of Fermilab's Short-Baseline Neutrino Program, hunting for signs of a possible fourth type of neutrino.

SYMMETRYMAGAZINE.ORG



The “God Particle”

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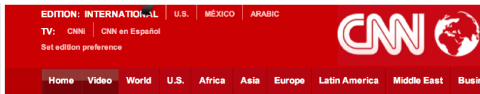
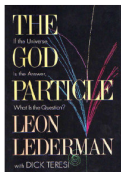
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The “God Particle” in the Popular Press:



Higgs and the holy grail of physics

By Lawrence M. Krauss, Special to CNN
July 6, 2012 -- Updated 1507 GMT (2307 HKT)

This boson is so central to the state of physics today, so crucial to our final understanding of the structure of matter, yet so elusive, that I have given it a nickname: the God Particle.

— Leon Lederman in “The God Particle”¹

¹He goes on to say that the name “Goddamn Particle” might be a more appropriate, “given its villainous nature and the expense it is causing” (published in 1993).



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Discovery of a Higgs Boson:

The Higgs Boson

ARTICLE

Journey in the Search for the Higgs Boson: The ATLAS and CMS Experiments at the Large Hadron Collider

M. Della Negra,¹ P. Jenni,² T. S. Virdee^{1*}

— *Science*, 21 Dec 2012, Vol 338

- July 2012: Simultaneously announce discovery of “a new boson”
- P-values less than 1 in 3 million
- Behavior broadly matches predictions of Standard Models for the Higgs boson.
- March 2013: The new boson is promoted to “a Higgs boson”



CERN's Large Hadron Collider

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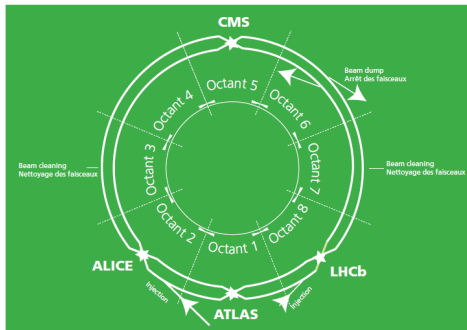
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CERN's Large Hadron Collider:

Proton-Proton Collisions:



LHC's 27km tunnel straddles the Swiss-French border.



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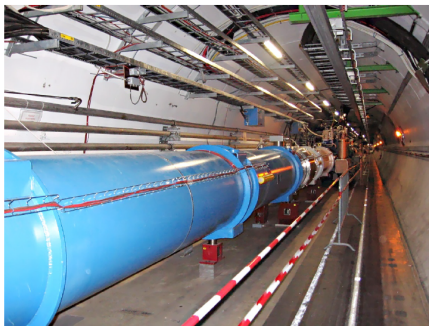
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The LHC Tunnel:

Forbes: Finding The Higgs Boson Cost \$13.25 Billion





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Cost Meaning for Statistical Analyses:

Researchers really don't want:

- 1 conflicting results from different analyses
- 2 change their minds about the choice of statistical methods
- 3 in any way be wishy-washy.

Committee decides the form of statistical analysis, and will

- 1 be conservative
- 2 be frequentist (subjectivity is bad!)
- 3 set out the analysis protocol in advance.

Probably the most careful frequentist analysis ever conducted.



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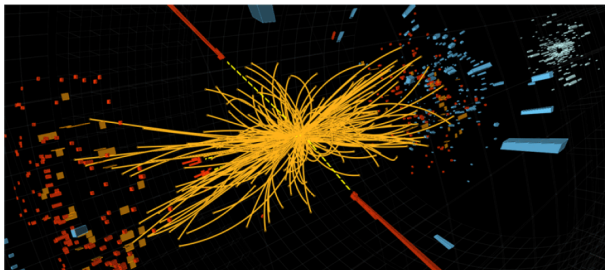
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Simulation of Proton-Proton Collision:



- Collisions produce a shower of new particles
- Unstable particles (e.g., Higgs) decay before detection
- Track trajectory, energy, and momentum to identify particles
- Compute *invariant mass* of particles formed in collisions



Decay Channels of the Higgs

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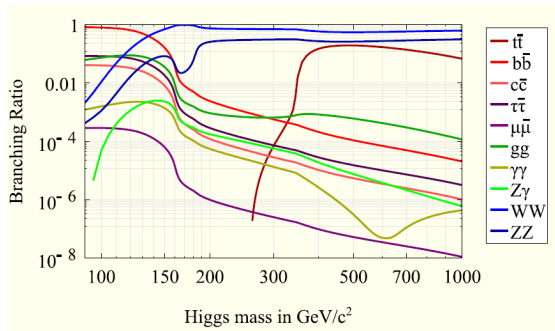
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Decay Channels of the Higgs:



The Standard Model predicts the relative probabilities of Higgs decay channels as a function of its mass.



Decay Channels of the Higgs

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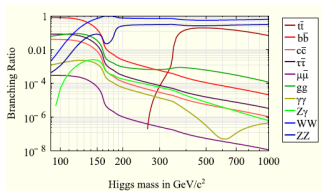
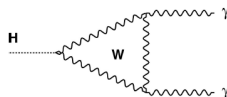


Fig. 4 Feynman diagram illustrating the decay mechanism of a Higgs boson, H, to two photons, γ , via quantum loop process involving a W boson.



The ATLAS Collaboration Bulletin 2012.208 (19th June 2012)

Published by arXiv



■ Primary Decay Channels for the Detection of the Higgs:

- ★ Higgs into two photons ($H \rightarrow \gamma\gamma$)
- ★ Higgs into two Z bosons ($H \rightarrow ZZ$), each decay into electrons or muons

■ Not the most likely decay channels (together less than 1%).

■ Advantageous signal-to-background ratio (Bayes Theorem)



Searching above Background

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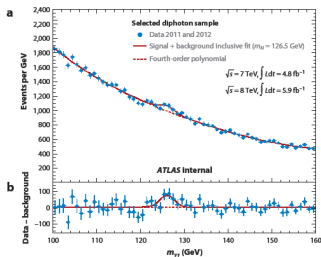
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Searching for the Bump above Background:



- Look at event counts in a number of invariant mass bins.
- Known physical processes result in background distribution
- Expect excess counts at invariant mass of Higgs boson.



Searching above Background

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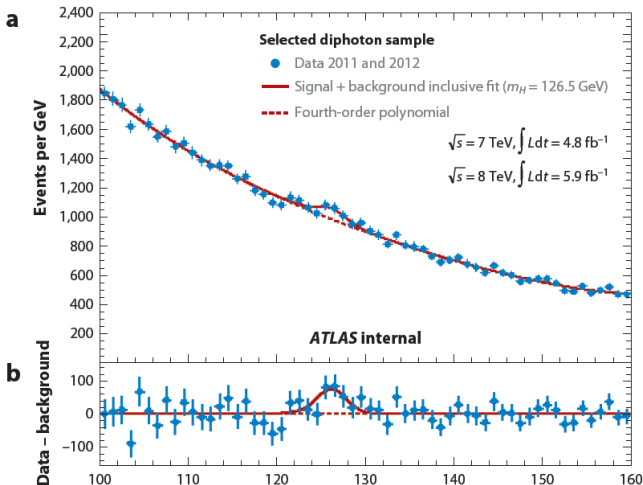
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Searching for the Bump above Background ($H \rightarrow \gamma\gamma$):





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Possible Statistical Outcomes:

- 1 **Discovery of new physical particle:** Conclude data are inconsistent with the null hypothesis of known background physics, but consistent with the hypothesized Higgs boson.
- 2 **Conclude data inconsistent with Higgs boson hypothesis.**
- 3 **Upper limit** on possible Higgs boson signal strength.
- 4 **Determine that experiment is not sensitive** enough to distinguish between new physics and background.

More sophisticated than accepting / rejecting H_0

Can we reject H_A ?



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A Simplified Model:

- Data involve multiple invariant mass bins for each channel.
- To focus statistical issues we consider a single channel:

$$N_m \sim \text{POISSON}(\beta_m + \kappa_m \mu),$$

β_m is the expected background count in bin m

κ_m is the expected Higgs Boson count in bin m

μ is either one or zero.

- To hedge against misspecification, treat μ as continuous
- For the moment we assume the β_m and κ_m are known.



Detection and Upper Limits:

Basic Model: $N_m \sim \text{POISSON}(\beta_m + \kappa_m \mu)$

Can we detect the Higgs? Conduct a hypothesis test that compares $H_0 : \mu = 0$ with $H_A : \mu > 0$.

Upper Limit: In the absence of detection, compute an UL on μ .

- The hypothesis is tested with a Likelihood Ratio Test
- Standard asymptotics do not apply (H_0 on boundary)
- The test is inverted to compute intervals and upper limits
- *Unified Approach* for UL and intervals: maintains frequency coverage (Feldman & Cousins, 1998)



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Exclusion and Sensitivity:

Basic Model: $N_m \sim \text{POISSON}(\beta_m + \kappa_m \mu)$

Can we exclude the Higgs? If the UL is too small to be consistent with $\mu = 1$, exclude possibility of Higgs.

Compute Sensitivity: Is the hypothesis test powerful enough to distinguish between background and Higgs?

Sensitivity is quantified by

- The median of the UL under H_0 or
- Smallest value of μ obtaining a given power for α -level test.

We would like both quantities to be small.



The Null Hypothesis

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Treating the Higgs as the Null Hypothesis:

We can formalize the relationship between detection, exclusion, and sensitivity by comparing two hypothesis tests:

Standard Test: Compare $H_0 : \mu = 0$ with the hypothesis $\mu > 0$

$$p_0 = \Pr(T \geq t_{\text{obs}} | H_0)$$

Reversed Test: Compare $H_A : \mu = 1$ with the hypothesis $\mu < 1$

$$1 - p_A = \Pr(T \leq t_{\text{obs}} | H_A)$$

Small values of $1 - p_A$ are evidence for rejecting H_A and excluding the Higgs boson².

²Here both p-values are defined to be left tail probabilities.



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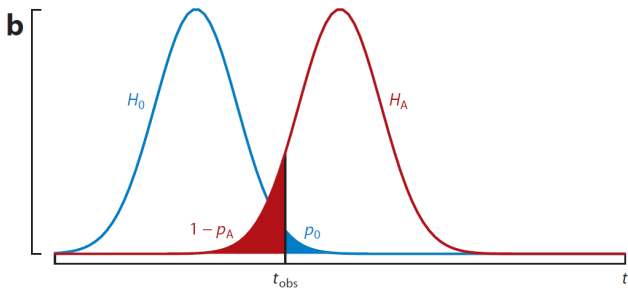
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The Two P-values:





Sensitive Test

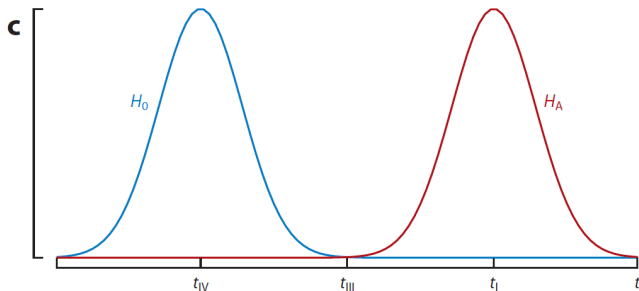
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A Sensitive Test:

Not all values of t_{obs} lead to a clear decision.





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Possible Outcomes:

Four possible outcomes of the two tests:

	$p_0 \leq \alpha_0$	$p_0 > \alpha_0$
$1 - p_A > \alpha_A$	Case I: Detection	Case II: Both models are consistent w/ data
$1 - p_A \leq \alpha_A$	Case III: Neither model is consistent w/ data	Case IV: Exclusion

$\alpha_0 \ll \alpha_A$: Particle physicists are more concerned with false discovery than with missing a signal.



Insensitive Test

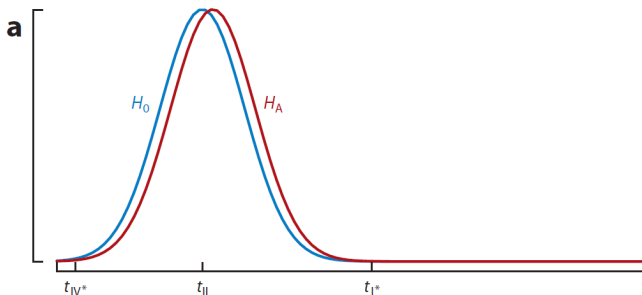
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An Insensitive Test:

Case II: Both models are consistent with data.



Two troublesome special cases: I^* and IV^*



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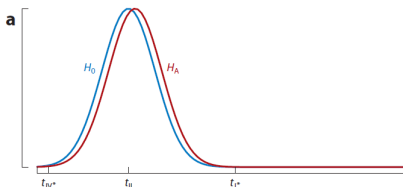
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Conclusion

The Trouble with an Insensitive Test:



Case IV* : $1 - p_A$ is small, but data are unlikely under both models.

- Exclusion questionable: there are fewer counts than expected even without a Higgs boson.
- Can be caused by a downward fluctuation in background.
- Case I* is less problematic, because $\alpha_0 \ll \alpha_A$.



Comparing the Sensitivity of Tests

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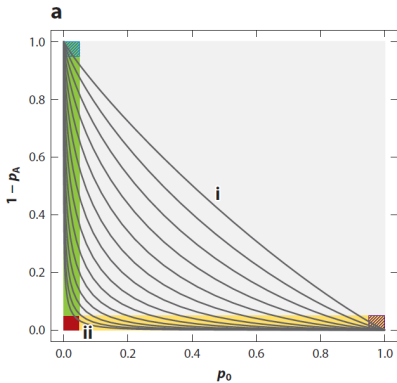
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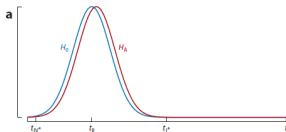


Green Detection
Yellow Exclusion
Red Neither Model
Grey Both Models

i insensitive test
ii sensitive test
 $\alpha_0 = \alpha_1 = 0.05$



Avoiding Exclusion Under an Insensitive Test:



Case IV*

- Exclusion is unwarranted
- Do not exclude $\mu = 1$ if both $1 - \rho_0$ and $1 - \rho_A$ are small

Read (2000) suggested excluding $\mu = \mu_0$ only if

$$CL_S = \frac{1 - \rho_A}{1 - \rho_0} = \frac{\Pr(T < t_{\text{obs}} | \mu = \mu_0)}{\Pr(T < t_{\text{obs}} | \mu = 0)} \leq \alpha_A.$$

CL_S ULs bound the values of μ that cannot be excluded.

Exclude μ if it makes $T < t_{\text{obs}}$ much less likely than does $\mu = 0$



Effect of Criterion

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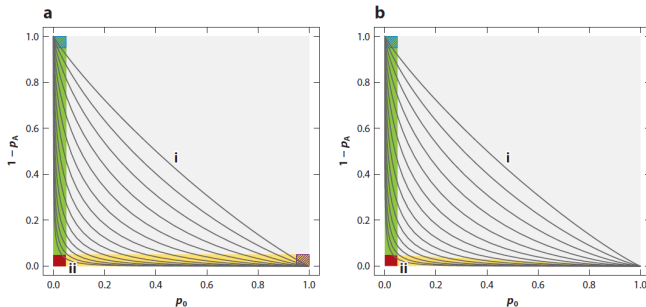
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Effect of the CL_S Criterion:



Troublesome purple region is almost eliminated with CL_S .



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5σ Detection Threshold:

P-values are converted to SDs from 0, under std normal.

- P-value of 0.025 corresponds to a “ 1.96σ ” result.
- July 2012 detections were 6σ and 5σ (ATLAS/CMS)

5σ is required for “discovery”

- High profile false discoveries led to conservative threshold
- Treat Higgs mass as known (multiple-testing)
- Calibration, systematic errors, and model misspecification
- Of course cranking down α_0 does not address these issues

“In particle physics, this criterion has become a convention ... but should not be interpreted literally³.”

³Glossary in the *Science* review of the 2012 CMS and ATLAS discoveries.



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Data Collection:

- LHC produces millions of proton collisions per second, most are uninteresting
- Fast online *triggers* save only ~ 100 events per second
- As few as 1 in 10^8 stored events may involve Higgs boson
- Further *cuts* aim to prune data, reduce background, and focus analysis on potential new physics
- Cuts are based on values of individual variables and supervised learning algorithms trained on simulations and manually classified events

Counts in each Category that survive the triggers and cuts are used in analysis.



Data Processing:

Within each channel, event counts are stratified into relatively homogeneous strata called categories.

- Categories with similar signal-to-noise increases power.
- Often based on flavors of decay, e.g.,
 - ★ $H \rightarrow ZZ$, 4 categories based on Z into electrons or muons.
- Sometimes based on boosted decision trees
 - ★ Aim to separate Higgs from background events.
 - ★ Use simulated data.
 - ★ Predictors: Momenta, energy, and presence of particles.
 - ★ Cut into categories based on the fitted probability of Higgs.

The search is based on the Stratified Sample within each Channel.



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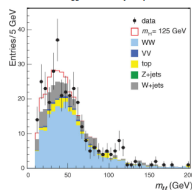
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Background Models:

Parametric models are used for background events

- Owing to cuts and stratification, different background distribution in each category-channel pair
- Functional forms set by simulation with same cuts/strata
- Various background models are considered
- E.g., exponential of polynomials (order ≤ 3), Bernstein polynomials (order ≤ 7), etc.
- The fitted parameters are determined with real data

Fig. 5 Distribution of the invariant mass of lepton pairs for the zero-jet $\mu\mu$ category in the search at 8 TeV for the SM Higgs boson decay to a pair of W bosons.



The CMS Collaboration Science 2012,238:1369-1376

Reprinted by AAAA





Mass-by-Mass Analysis:

Search conducted separately on fine grid of Higgs masses, m_H

- Integrated search may overlook masses w/ low sensitivity.
- Mass-by-mass search allows sensitivity, ULs, and p-values to be computed for each potential mass.

Let m index **mass** bins, s **categories** (strata), and c **channels**.

For a given Higgs mass, m_H ,

$$N_{msc} \sim \text{POISSON} \left(\beta_{sc}(\theta_{sc}, m) + \kappa_{sc}(\phi_{sc}, m) \mu \right).$$

- Models and their parameters vary among channel/categories
- Source model specifies the expected bump above background



CL_S Upper Limits on μ as a Function of m_H

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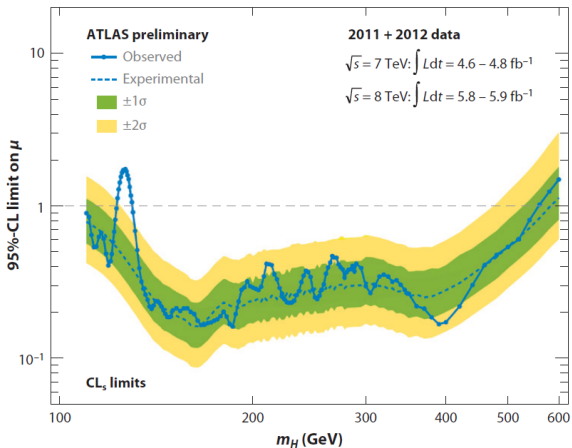
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CL_S Upper Limits on μ as a Function of m_H :





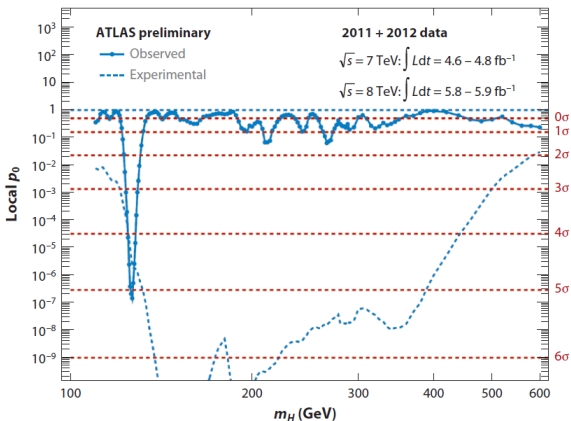
Local P-values as a function of m_H

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Local P-values as a function of m_H :



But what about Multiple Testing?



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The Look Elsewhere Effect:

- Let $T(m_H) \sim \chi_s^2$ be the LRT for the m_H -specific test.
- Davies (1987) shows

$$\Pr\left(\max_{m_H} T(m_H) > c\right) \leq \Pr(\chi_s^2 > c) + E(M(c) | H_0),$$

- $M(c)$ = number of times $T(m_H)$ increases above c as $m_H \uparrow$
- To avoid MC evaluation of $E(M(c)|H_0)$, let $c_0 \ll c$ and use⁴

$$E(M(c) | H_0) = E(M(c_0) | H_0) \left(\frac{c}{c_0}\right)^{(s-1)/2} \exp\left(-\frac{(c-c_0)}{2}\right),$$

- 6 σ / 5 σ significances reduce to 5.1 σ / 4.6 σ (ATLAS/CMS)

⁴Gross and Vitells (2010)



Estimating the Higgs Mass

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Estimating the Higgs Mass:

- The Higgs mass is estimated via a unified analysis.
- Source model, κ_{SC} depends on m_H :

$$N_{msc} \sim \text{POISSON} \left(\beta_{SC}(\theta_{SC}, m) + \kappa_{SC}(\phi_{SC}, m, m_H) \mu \right).$$

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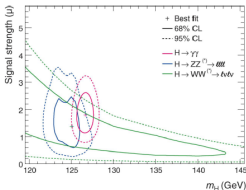
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Fig. 12 Confidence intervals comparing mass and signal strength for the $H \rightarrow \gamma\gamma$, $H \rightarrow ZZ \rightarrow 4\ell\ell$, and $H \rightarrow WW \rightarrow \ell\nu\ell\nu$ channels, including all systematic uncertainties.

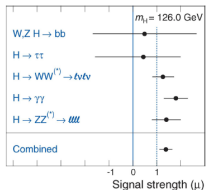


The ATLAS Collaboration Science 2012:338:1576-1582



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Fig. 13 Measurements of the signal strength parameter μ for $m_H = 126$ GeV for the individual channels and their combination.



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Bayesian Analysis:

Is a coherent analysis of the full model possible?

$$N_{msc} \sim \text{POISSON}\left(\beta_{sc}(\theta_{sc}, m) + \kappa_{sc}(\phi_{sc}, m, m_H) \mu\right).$$

Consider the LRT for $\mu = 0$ vs $\mu > 0$:

- 1 The null space is on the boundary.
- 2 Worse, m_H is unidentified under H_0 .
- 3 Still worse, a sharp null: p-value can vastly overstate evidence for alternative⁵, Jeffrey-Lindley paradox.

How about Bayes Factors:

- 1 Highly dependent on choice of prior for μ and m_H .

⁵E.g., Berger & Delampady, *Testing Precise Hypotheses*, Stat. Sci., 1987



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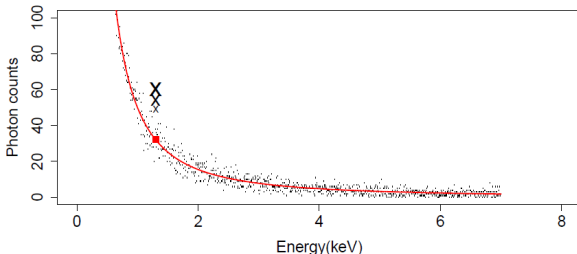
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A Similar Example in High-Energy Astrophysics:

$$Y_i^{\text{ind}} \sim \text{POISSON}(\alpha E_i^{-\beta} + \omega I_{\{i=\mu\}})$$

Is there sufficient evidence to conclude that $\omega > 0$??

- Fix α and β throughout
- The “true” emission line is at $\mu = 1.3$ keV.





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Comparing Bayes Factors with P-values:

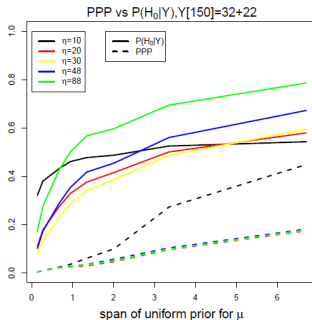
Prior on spectral line: $\omega \sim U(0, \eta)$ and $\mu \sim U(1.3 \pm \kappa)$.

$P(H_0|Y)$ under 50-50 prior.

P-values vastly overstate evidence for line.

Prior on μ

- *let's us decide where to look.*
- *penalty for many looks.*
- *i.e., look elsewhere effect.*
- *Sensitivity of Bayes Factor to prior for μ is sensible.*
- *Prior on line intensity: look for weak or strong lines.*





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R developments:

Physicists at the Large Hadron Collider would like your contribution in studying the Higgs Boson. It's true that they found the Higgs boson which is responsible for giving all particles their mass nearly two years ago, but its exact behavior is still mysterious. Now, the scientists are asking coders to develop algorithms that can reveal the Higgs' properties.



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R code:

Physicists have to hunt through the noisy mess of other particles to see the Higgs' weak decay signal. Bellow is R Code designed to help in this task:

```
ww <- sum(train$Weight)*Weight/sum(Weight)      ## normalization of Weight vector
ss <- sum(ww[which(pp>thresh & gt=="s")])        ## true positive rate
bb <- sum(ww[which(pp>thresh & gt=="b")])        ## false positive rate
br <- 10.0                                       ## regulator
AMS <- sqrt(2.0*((ss+bb+br)*log(1+ss/(bb+br))-ss)) ## AMS score

AMS = function(pred,real,weight)
{
  #a = table(pred,real)
  pred_s_ind = which(pred==1)
  real_s_ind = which(real==1)
  real_b_ind = which(real==0)
  s = sum(weight[intersect(pred_s_ind,real_s_ind)])
  b = sum(weight[intersect(pred_s_ind,real_b_ind)])

  b_tau = 10
  ans = sqrt(2*((s+b+b_tau)*log(1+s/(b+b_tau))-s))
  return(ans)
}
```



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Thank You!

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