# Steady state holographic turbulence

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### Turbulence

### Recall:

# $\vec{v} + \vec{v} \cdot \vec{\nabla} \vec{v} = - \vec{\nabla} p + \nu \nabla^2 \vec{v} + \vec{f}$ $\vec{\nabla} \cdot \vec{v} = 0$

### The Navier Stokes equations describe a multitude of phenomenon:









$$\vec{v} + \vec{v} \cdot \vec{\nabla} \vec{v} = - \vec{\nabla} p + \nu \nabla^2 \vec{v} + \vec{\nabla} \cdot \vec{v} = 0$$

One characteristic of turbulence is the scaling behaviour of the kinetic energy (per unit mass).

### Define

### Then

$$\overline{\hat{\epsilon}} \propto k^{-5/3}$$



### $\hat{\epsilon}(k)dk$ -Amount of kinetic energy between k and k + dk

f

$$\vec{v} + \vec{v} \cdot \vec{\nabla} \vec{v} = - \vec{\nabla} p + \nu \nabla^{2} \vec{v} + \vec{v} \vec{v} \vec{v} = 0$$

## One characteristic of turbulence is the scaling behaviour of the kinetic energy (per unit mass).

$$\overline{\hat{\epsilon}} \propto k^{-5/3}$$



 $+ \vec{f}(k_f)$  $\vec{f}(\not{k}_f) = 0$ 



$$\vec{v} + \vec{v} \cdot \vec{\nabla} \vec{v} = -\vec{\nabla} p + \nu \nabla^2 \vec{v} + \vec{f}(k_f) \qquad \vec{f}(k_f) = 0$$
$$\vec{\nabla} \cdot \vec{v} = 0$$

One characteristic of turbulence is the scaling behaviour of the kinetic energy (per unit mass).

$$\overline{\hat{\epsilon}} \propto k^{-5/3}$$

This is part of a broader set of predictions:

$$\overline{\left((\overrightarrow{v}(\overrightarrow{r})-\overrightarrow{v}(0))\cdot\widehat{r}\right)^n} \propto |r|^{\frac{n}{3}}$$



$$(n = 2)$$

$$\vec{v} + \vec{v} \cdot \vec{\nabla} \vec{v} = - \vec{\nabla} p + \nu \nabla^2 \vec{v} + \vec{\nabla} \cdot \vec{v} = 0$$

Thisrisad as the too factor set of predictions:

$$\overline{\left(\left(\overrightarrow{v}(\overrightarrow{r})-\overrightarrow{v}(0)\right)\cdot\widehat{r}\right)^{n}} \propto |r|^{\frac{n}{3}}$$

For n=3,

$$F_3 = \overline{\left( (\overrightarrow{v}(\overrightarrow{r}) - \overrightarrow{v}(0)) \cdot \widehat{r} \right)^3} \propto |r|$$



 $\vdash \vec{f}(k_f) \qquad \overline{\vec{f}(k_f)} = 0$ 



 $\vec{v} + \vec{v} \cdot \vec{\nabla} \vec{v} = -\vec{\nabla} p + \nu \nabla^2 \vec{v} +$  $\overrightarrow{\nabla} \cdot \overrightarrow{v} = 0$ 

A broad set of predictions

$$\overline{\left((\overrightarrow{v}(\overrightarrow{r})-\overrightarrow{v}(0))\cdot\widehat{r}\right)^n} \propto |r|^{\frac{n}{3}}$$



$$\vec{f}(k_f) \qquad \overline{\vec{f}(k_f)} = 0$$



Holographic turbulence  $R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = 0$ 



•Maldacena, 1997 •Maldacena, 1997  $ds^2 = r^2(-f(r)dt^2 + (dx^1)^2 + (dx^2)^2) + \frac{dr^2}{r^2 f(r)}$ •Bhattacharyya et. al. 2007



Holographic turbulence  $R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 0$ 



### •Maldacena, 1997

- •Witten, 1998
- •Bhattacharyya et. al. 2007

### $\nabla_{\mu}T^{\mu\nu} = 0$ $T^{\mu\nu} = (\epsilon + P) u^{\mu} u^{\nu} + P \eta^{\mu\nu} + \dots$



Holographic turbulence  $R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 0$ 



- •Maldacena, 1997
- •Witten, 1998
- •Bhattacharyya et. al. 2007
- •Adams et. al. 2013

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$$\vec{v} + \vec{v} \cdot \vec{\nabla} \vec{v} = - \vec{\nabla} p + \nu \nabla^2 \vec{v}$$
$$\vec{\nabla} \cdot \vec{v} = 0$$



t = 2496

t = 3001





Holographic turbulence  $R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 0$ 



- •Maldacena, 1997
- •Witten, 1998
- •Bhattacharyya et. al. 2007
- •Adams et. al. 2013

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Holographic turbulence  $R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 0$ 



- •Maldacena, 1997
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- •Adams et. al. 2013



























![](_page_14_Figure_3.jpeg)

![](_page_14_Figure_4.jpeg)

**Btdobædigtavitylendeurbulence**  $R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 0$  $r \to \infty$  $\bigcirc$  $\overline{g_{mn}} = \eta_{mn}$ 

![](_page_15_Picture_1.jpeg)

![](_page_15_Figure_3.jpeg)

![](_page_15_Figure_4.jpeg)

Stochastic gravity and turbulence  $W_{\mu\nu} = 0$  $r \to \infty$  $\sum_{mn}^{\infty} \frac{1}{2} Rg_{mn} - \frac{12}{2}g_{mn} = 0$  $r = r_h$  such that at t<0 we have  $ds^{2} = r^{2}(-f(r)dt^{2} + (dx^{1})^{2} + (dx^{2})^{2})$ 

### At t>0 we have

 $g_{mn} \xrightarrow[r \to \infty]{} g^{(0)}_{\mu\nu}$ 

 $g^{(0)}_{\mu\nu}dx^{\mu}dx^{\nu} =$ 

![](_page_16_Figure_5.jpeg)

$$f(r) = \left(1 - \left(\frac{r_0}{r}\right)^3\right)$$

$$= -(1+Q)dt^2 + (dx^1)^2 + (dx^2)^2$$

We wish to solve

$$R_{mn} - \frac{1}{2} Rg_{mn} - \frac{12}{\ell^2} g_{mn} = 0$$

At t>0 we have

$$g_{mn} \xrightarrow[r \to \infty]{} g_{\mu\nu}^{(0)} \qquad g_{\mu\nu}^{(0)} dx^{\mu} dx^{\nu} :$$

where  

$$Q = q \qquad \dot{q} = -\frac{q}{\tau} + \frac{\xi}{\tau}$$

$$\overline{\xi(t, \vec{x})} = 0 \qquad \overline{\xi(t, \vec{x})}\xi(t', \vec{x'}) =$$

![](_page_17_Figure_6.jpeg)

### $= -(1+Q)dt^2 + (dx^1)^2 + (dx^2)^2$

 $= D(\overrightarrow{x} - \overrightarrow{x'})\delta(t - t')$ 

 $\hat{D}(\vec{k}) = \delta(|\vec{k}| - k_f)$ 

![](_page_17_Picture_10.jpeg)

We wish to solve

$$R_{mn} - \frac{1}{2} Rg_{mn} - \frac{12}{\ell^2} g_{mn} = 0$$

At t>0 we have

$$g_{mn} \xrightarrow[r \to \infty]{} g^{(0)}_{\mu\nu} \qquad g^{(0)}_{\mu\nu} dx^{\mu} dx^{\nu} =$$

The energy momentum tensor,  $T^{\mu\nu}$ , can be read off of the metric. After averaging we obtain  $\overline{T^{\mu\nu}}$ .

![](_page_18_Figure_6.jpeg)

 $= -(1+Q)dt^2 + (dx^1)^2 + (dx^2)^2$ 

We wish to solve

$$R_{mn} - \frac{1}{2} Rg_{mn} - \frac{12}{\ell^2} g_{mn} = 0$$

At t>0 we have

$$g_{mn} \xrightarrow[r \to \infty]{} g^{(0)}_{\mu\nu} \qquad g^{(0)}_{\mu\nu} dx^{\mu} dx^{\nu} =$$

In practice, we need to solve numerically.

Solving in the right order allows us to rewrite the Einstein equations as a set of ordinary stochastic differential equations. (Chesler, Yaffe, 2013)

![](_page_19_Figure_7.jpeg)

- $= -(1+Q)dt^{2} + (dx^{1})^{2} + (dx^{2})^{2}$

We solved

$$R_{mn} - \frac{1}{2} Rg_{mn} - \frac{12}{\ell^2} g_{mn} = 0$$

Did this many times, and then computed the average:

$$\overline{T_{\mu
u}}$$

![](_page_20_Figure_5.jpeg)

![](_page_21_Figure_2.jpeg)

 $\overline{T_{\mu\nu}}$ 

![](_page_22_Figure_2.jpeg)

![](_page_22_Figure_3.jpeg)

![](_page_22_Picture_4.jpeg)

Recall:

$$\hat{\epsilon} = \int \frac{1}{2} \rho |\hat{v}|^2 k d\theta_k \propto k^{-\frac{5}{3}}$$

Define:

$$T^{\mu}_{\ \nu}u^{\nu} = -\epsilon u^{\nu}$$

 $\overline{T_{\mu
u}}$ 

with

$$u^{\mu} = \gamma\left(1, \overrightarrow{v}\right)$$

![](_page_23_Figure_7.jpeg)

![](_page_23_Figure_8.jpeg)

![](_page_24_Figure_0.jpeg)

![](_page_24_Picture_2.jpeg)

**g**<sub>01</sub>

![](_page_25_Figure_2.jpeg)

There's an apparent horizon at

 $0.9 \le \rho = \rho_h(t, x_1, x_2) \le 1.1$ 

We would like to find a geometric quantity that encodes

$$\overline{\left((\overrightarrow{v}(\overrightarrow{r})-\overrightarrow{v}(0))\cdot\widehat{r}\right)^n} \propto |r|^{\zeta_n}$$

(Recall that  $\zeta_n = n/3$  for Kolmogorov theory)

![](_page_26_Figure_6.jpeg)

We would like to find a geometric quantity that encodes

$$\overline{\left((\overrightarrow{v}(\overrightarrow{r})-\overrightarrow{v}(0))\cdot\widehat{r}\right)^n} \propto |r|^{\zeta_n}$$

(Recall that  $\zeta_n = n/3$  for Kolmogorov theory)

An alternate expression which encodes  $\zeta_n$  is  $\epsilon_{\mathbf{R}}(\mathbf{x}) = \frac{1}{V_{\mathbf{R}}} \int_{|\mathbf{x}-\mathbf{x}'| < \mathbf{R}} \left(\partial_i v_j + \partial_j v_i\right)^2 d^d x'$ 

![](_page_27_Figure_5.jpeg)

We would like to find a geometric quantity that encodes

$$\overline{\left((\overrightarrow{v}(\overrightarrow{r})-\overrightarrow{v}(0))\cdot\widehat{r}\right)^n} \propto |r|^{\zeta_n}$$

theory)

500 400 (Recall that  $\zeta_n = n/3$  for Kolmogorov 300  $\mathbf{x}^2$ 200 An alternate expression which encodes  $\zeta_n$  is 100  $\epsilon_{\mathbf{R}}(\mathbf{x}) = \frac{1}{\mathbf{V}_{\mathbf{R}}} \int_{|\mathbf{x} - \mathbf{x}'| \le \mathbf{R}} \left(\partial_i v_j + \partial_j v_i\right)^2 d^d \mathbf{x}'$ 100 200 300 400 500 0

$$\overline{(\epsilon_R(x))^n} \sim R^{\zeta_n - \frac{n}{3}}$$

![](_page_28_Picture_7.jpeg)

(Recall that  $\zeta_n = n/3$  for Kolmogorov theory)

An alternate expression which encodes <sup>0</sup>  $\zeta_{n} \text{ is } \epsilon_{R}(x) = \frac{1}{V_{R}} \int_{|x-x'| \le R} \left( \partial_{i} v_{j} + \partial_{j} v_{i} \right)^{2} d^{d} x' \qquad ^{0.5}$   $\overline{\left( \epsilon_{R}(x) \right)^{n}} \sim R^{\zeta_{n} - \frac{n}{3}} \qquad \rho$ 

As it turns out, the extrinsic curvature of the horizon,  $K_{ii}$ , is proportional to the shear. Thus,

$$\epsilon_{R}(x) \sim e_{R}(x)$$
$$e_{R}(x) = \frac{1}{D_{R}} \int_{|x-x'| \leq R} K_{i}^{j} K_{j}^{i} d^{d} x'$$

![](_page_29_Figure_5.jpeg)

![](_page_29_Figure_6.jpeg)

### Holographic and turbulence (Recall that $\zeta_n = n/3$ for Kolmogorov 5. × 10<sup>-4</sup> $\overline{e_r}$ theory) $1. \times 10^{-4}$ 5. × 10<sup>-5</sup> An alternate expression which encodes $1. \times 10^{-5}$ $\zeta_n$ is $5. \times 10^{-6}$ $\epsilon_R(x) = \frac{1}{V_R} \int_{|x-x'| < R} \left( \partial_i v_j + \partial_j v_i \right)^2 d^d x'$ $1. \times 10^{-6}$ 0.00 0.05 $\left(\epsilon_R(x)\right)^n \sim R^{\zeta_n - \frac{n}{3}}$ $1. \times 10^{-7}$ 5.×10<sup>-8</sup> As it turns out, the extrinsic curvature of the $1. \times 10^{-8}$ horizon, $K_{ii}$ , is proportional to the shear. Thus, $5. \times 10^{-9}$ $1. \times 10^{-9}$ 5. × 10<sup>-10</sup>

$$\epsilon_R(x) \sim e_R(x)$$
$$e_R(x) = \frac{1}{D_R} \int_{|x-x'| \le R} K_i^{j} K_j^{i} d^d x'$$

![](_page_30_Figure_5.jpeg)

### Summary

![](_page_31_Picture_1.jpeg)

![](_page_31_Figure_2.jpeg)

![](_page_31_Figure_3.jpeg)

![](_page_32_Figure_0.jpeg)

![](_page_32_Picture_2.jpeg)

![](_page_32_Picture_3.jpeg)

![](_page_32_Picture_4.jpeg)

 $|F_{\vec{n}}|$ 

1.5

1.0

0.5

 $\overline{T_{\mu\nu}} = \overline{g_{\mu\nu}^{(3)}}$ 

![](_page_34_Figure_2.jpeg)

![](_page_34_Figure_3.jpeg)

 $T_{01} = \frac{1}{L} \sum_{\overrightarrow{n}} F_{\overrightarrow{n}} e^{i\frac{2\pi\overrightarrow{n}}{L}\overrightarrow{x}}$ 

![](_page_34_Picture_5.jpeg)

|F<sub>n</sub>|

1.5

1.0

0.5

 $\overline{T_{\mu\nu}} = \overline{g_{\mu\nu}^{(3)}}$ 

![](_page_35_Figure_2.jpeg)

![](_page_35_Figure_3.jpeg)

 $T_{01} = \frac{1}{L} \sum_{\overrightarrow{n}} F_{\overrightarrow{n}} e^{i\frac{2\pi \overrightarrow{n}}{LL}\overrightarrow{x}}$ 

![](_page_35_Picture_5.jpeg)

|F<sub>n</sub>|

1.5

1.0

0.5

 $\overline{T_{\mu\nu}} = \overline{g_{\mu\nu}^{(3)}}$ 

![](_page_36_Figure_2.jpeg)

![](_page_36_Figure_3.jpeg)

 $T_{01} = \frac{1}{L} \sum_{\vec{n}} F_{\vec{n}} e^{i \frac{2\vec{n}\vec{n}}{LL} \vec{x}}$ 

![](_page_36_Picture_5.jpeg)

![](_page_37_Figure_2.jpeg)

 $|F_{\vec{n}}|$ 

1.5

1.0

0.5

 $\overline{T_{\mu\nu}} = \overline{g_{\mu\nu}^{(3)}}$ 

![](_page_38_Figure_2.jpeg)

![](_page_38_Figure_3.jpeg)

$$T_{01} = \frac{1}{L} \sum_{\overrightarrow{n}} F_{\overrightarrow{n}} e^{i \frac{2\pi \overrightarrow{n}}{L} \overrightarrow{x}}$$

![](_page_38_Figure_5.jpeg)

![](_page_38_Picture_6.jpeg)

### Stochastic gravity and turbulence 70 $r \rightarrow \infty$ 70 $\overline{T_{\mu\nu}} = \overline{g_{\mu\nu}^{(3)}}$ $\bigcirc$ 11 30 $r = r_h$ $|F_{\vec{n}}|$ 20 $\hat{F}(|\overrightarrow{n}|) = \int F_{\overrightarrow{n}} |\overrightarrow{n}| \, d\theta$ 10 1.5 $n_2$ Ê(n) 1000<sub>₣</sub> 1.0 <u>— Single run</u> -10 Ensemble average 100 0.5 10 -20 -30 0.10 -20 20 -30 -10 10 30 0 0.01 10 20 2 5

![](_page_39_Figure_2.jpeg)

![](_page_39_Picture_3.jpeg)