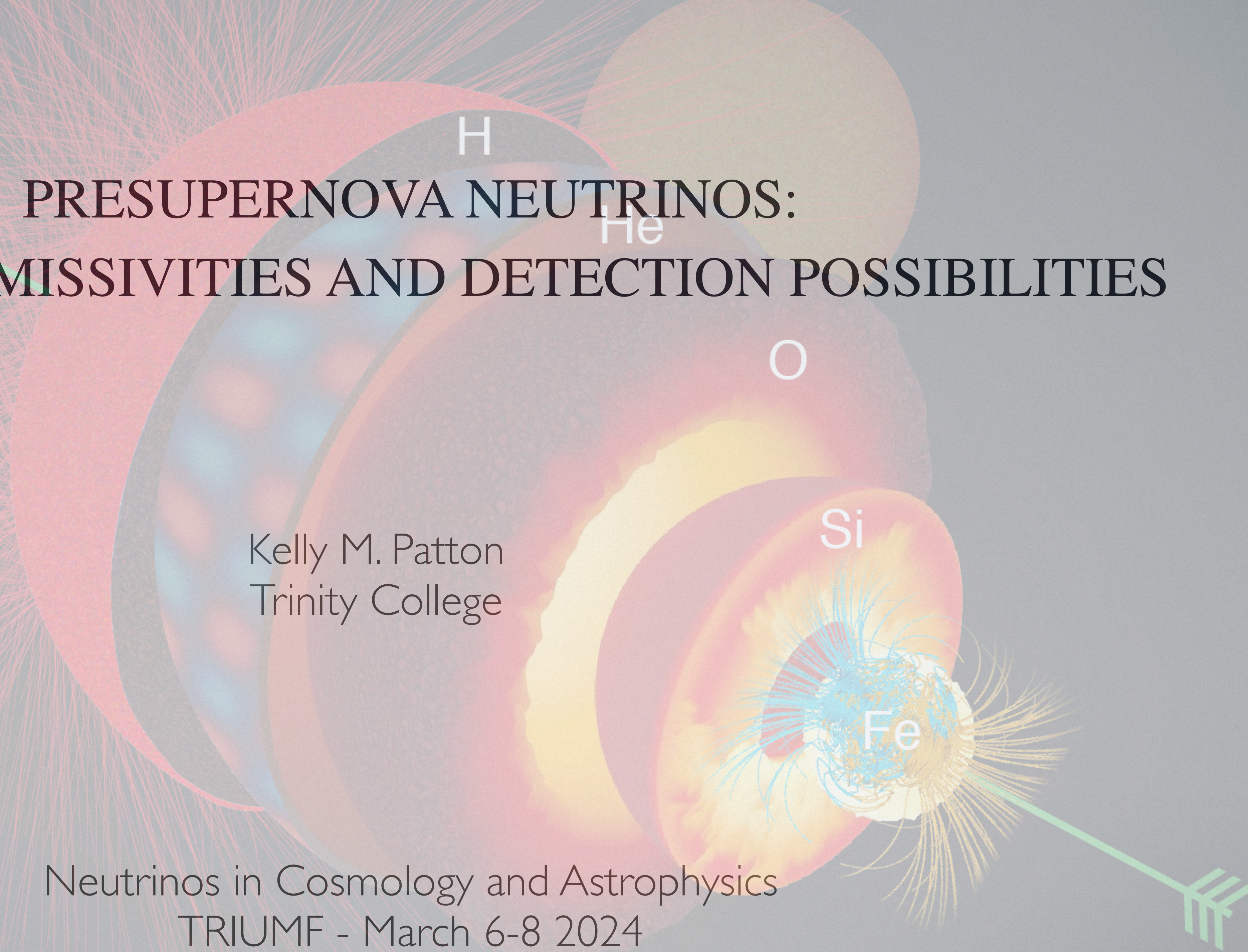




PRESUPERNOVA NEUTRINOS: REALISTIC EMISSIVITIES AND DETECTION POSSIBILITIES



Kelly M. Patton
Trinity College

Neutrinos in Cosmology and Astrophysics
TRIUMF - March 6-8 2024



EVOLUTION OF A MASSIVE STAR

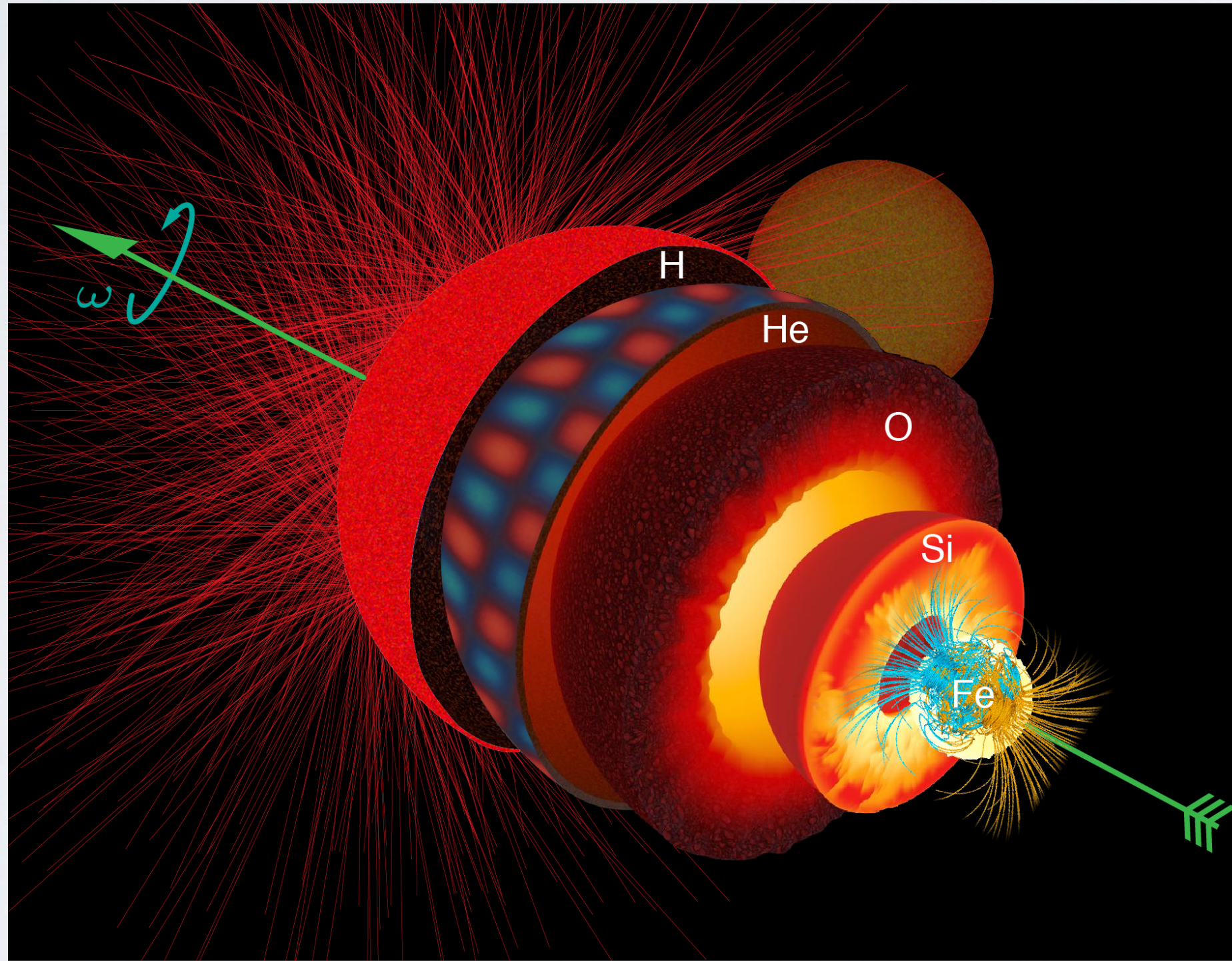
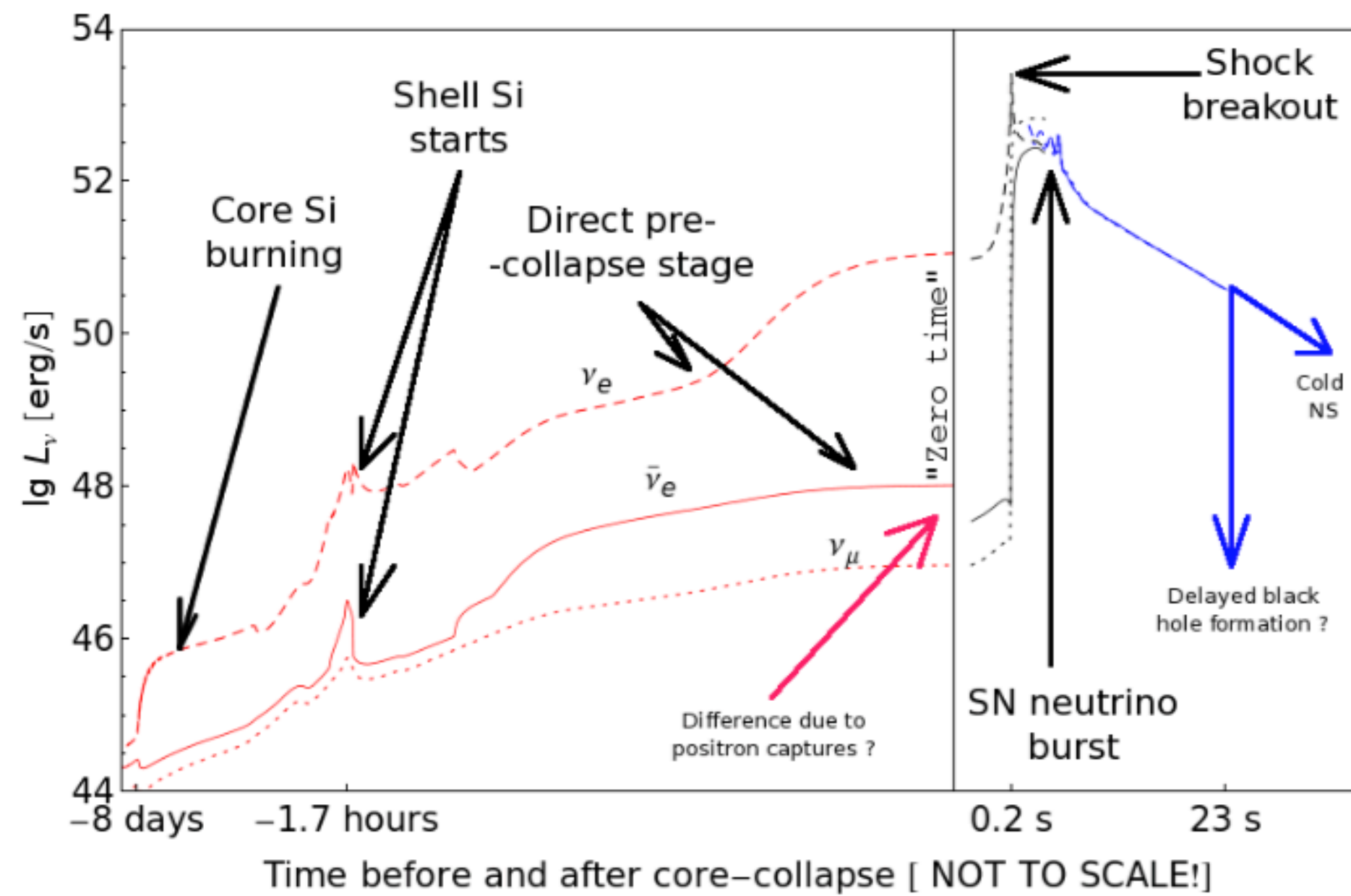


Image provided by F.X. Timmes

- Stars are born burning H in He in the core
- When H is exhausted, the star will burn He into O
- And so on, until an Fe core is created
- This leaves an onion structure behind, and the shells continue burning

PRESUPERNOVA NEUTRINOS



Odrzywolek and Heger *Acta Physica Polonica B* **41**, 1611 (2010)

- As the star gets hotter and denser, neutrino production ramps up
- After collapse we have a huge “burst” of neutrinos
- Notice that the pre-collapse luminosity is only a couple orders of magnitude below the “burst”
- We know we can see the burst (SNI 987A), so the question is, can we see the presupernova?

SNI 987A Image from <https://chandra.harvard.edu/photo/2017/sn1987a/>

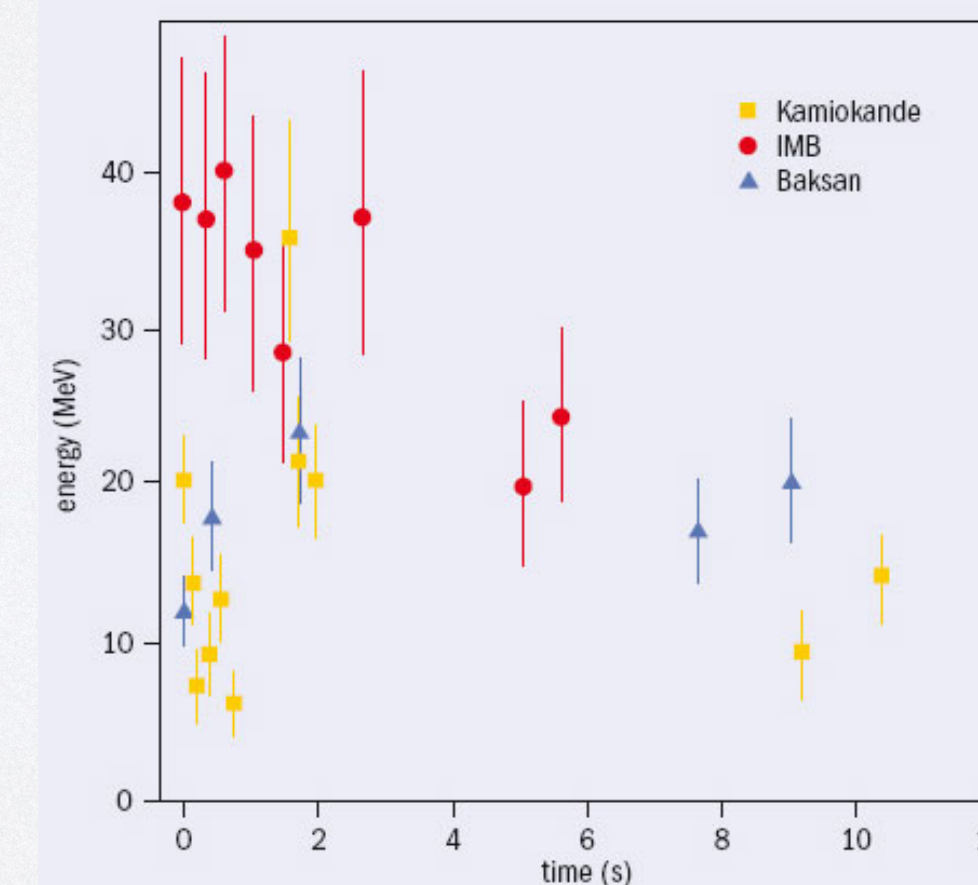
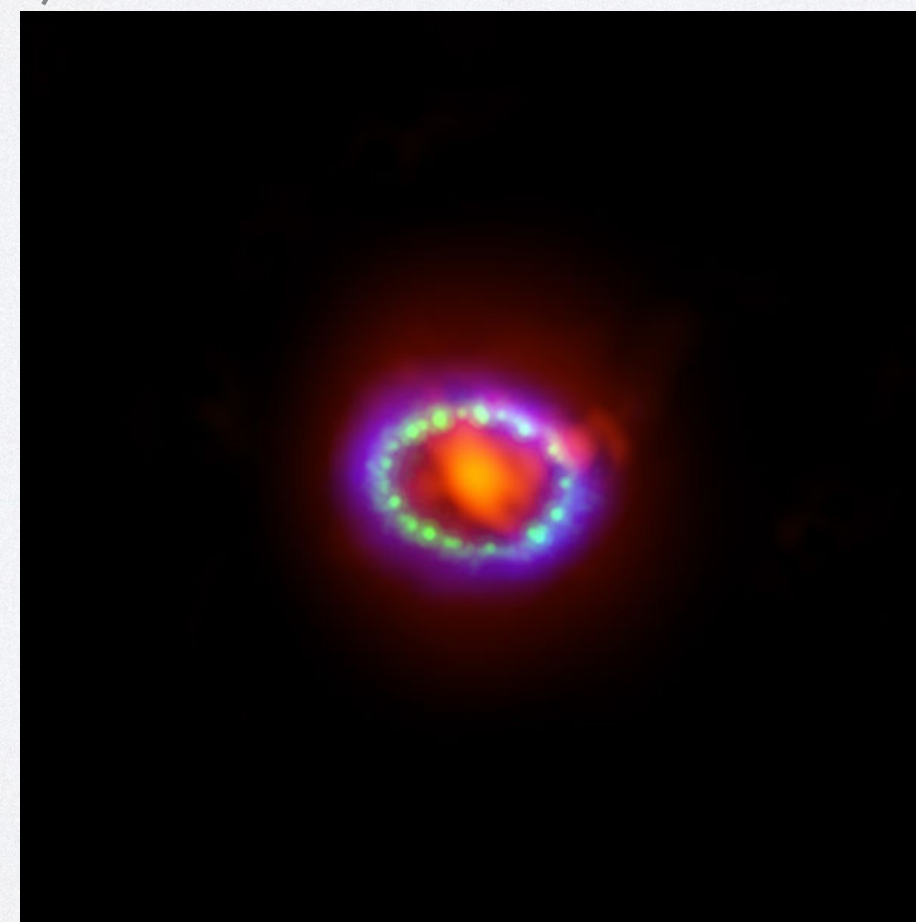
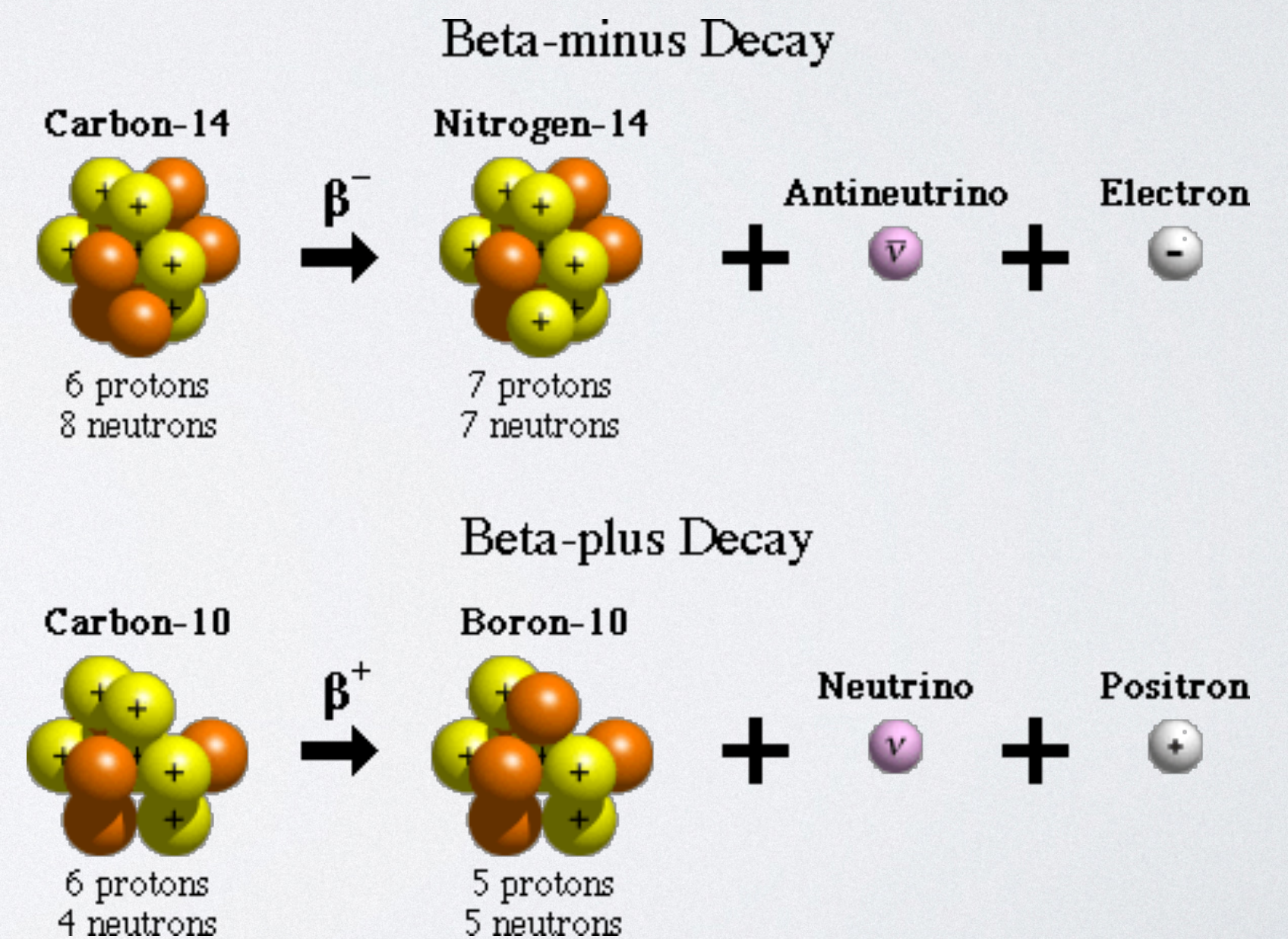


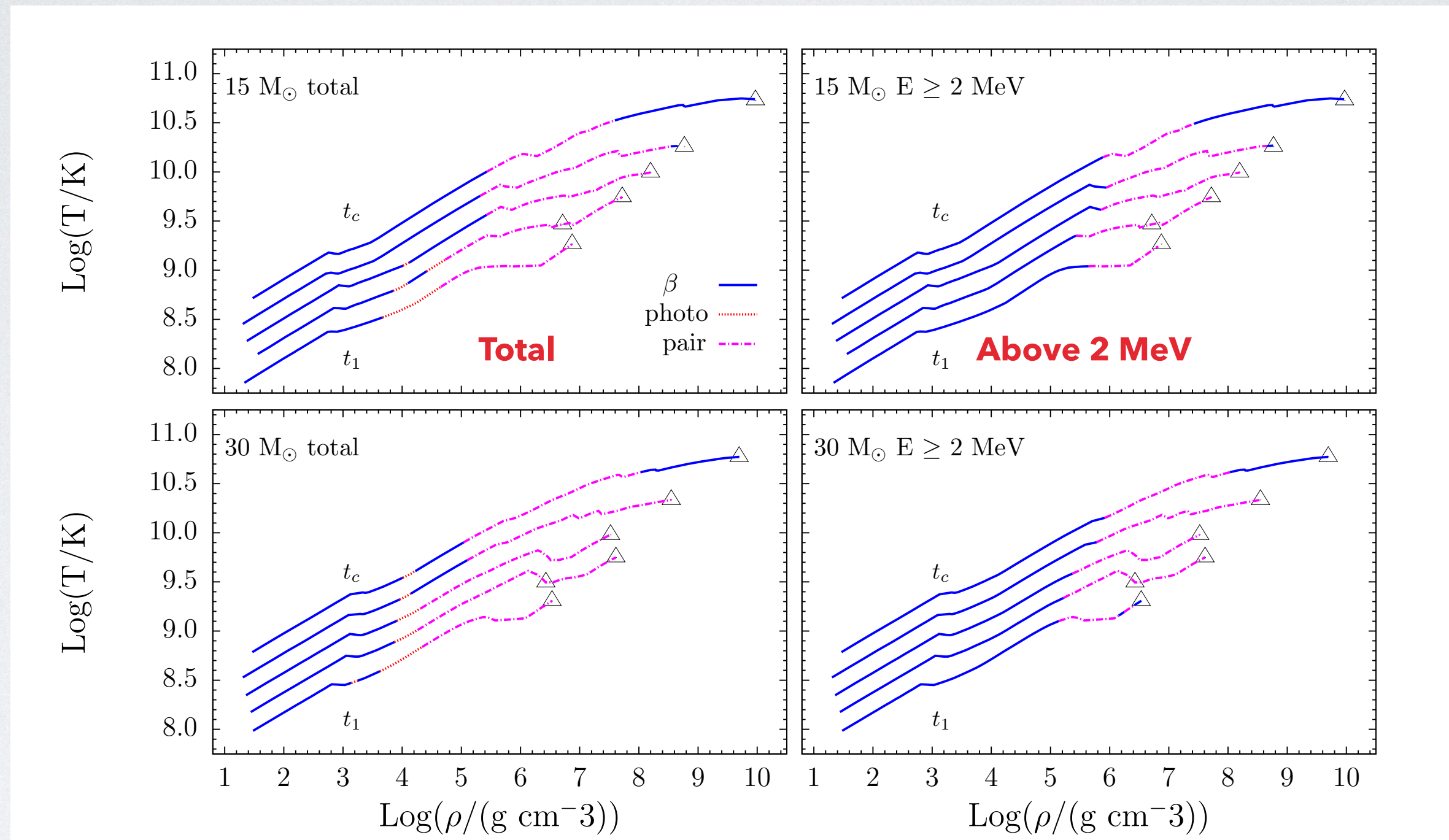
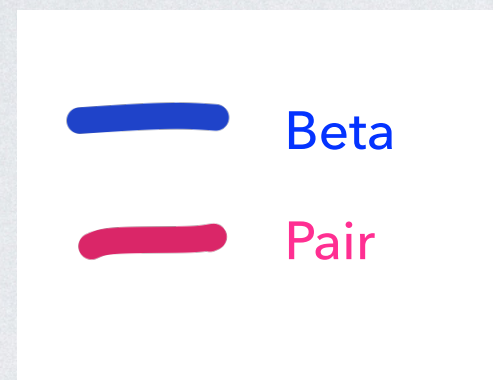
Figure from <https://cerncourier.com/a/sn1987a-heralds-the-start-of-neutrino-astronomy/>

PRESUPERNOVA NEUTRINOS

- The hot dense environment is a neutrino producing factory
- Two methods:
 - Thermal processes governed by T , ρ , and Y_e
 - Dominated by pair annihilation
 - Nuclear processes governed by T , ρ , Y_e , and isotopic composition



PRESUPERNOVA NEUTRINOS: DOMINANT PROCESSES



Curve shifted upward by $0.2 \cdot (n-1)$ for $n = 2 - 5$

- Mainly pair or beta, with a few islands of photoneutrino dominance in total emissivity
- For detectable energies ($E > 2$ MeV), beta dominance is extended
- Beta very important in the core at t_c

PRESUPERNOVA NEUTRINOS: PAIR ANNIHILATION

20 M_{\odot} star at D=1 kpc ~ 3200 lyr

- Odryzwolek et al. in 2004
- Looked at core Si burning stage
- Assumed neutrinos only from pair annihilation
- Calculated the number of events in different detectors
- ~10s of events in current detectors
- ~100s of events in future detectors

Detector	Mass [kton]	Reactions	Number of Targets	Flux at 1 kpc [$cm^{-2} day^{-1}$]	Event rate [day^{-1}]
Borexino	0.3 (C_9H_{12})	$\bar{\nu}_e + p \rightarrow e^+ + n$	$1.80 \cdot 10^{31}$	$2.8 \cdot 10^{11}$	0.34
		$\nu_e + e^- \rightarrow \nu_e + e^-$	$9.92 \cdot 10^{31}$	$2.8 \cdot 10^{11}$	0.49
		$\bar{\nu}_e + e^- \rightarrow \bar{\nu}_e + e^-$	$9.92 \cdot 10^{31}$	$2.8 \cdot 10^{11}$	0.19
		$\nu_{\mu,\tau} + e^- \rightarrow \nu_{\mu,\tau} + e^-$	$9.92 \cdot 10^{31}$	$1.0 \cdot 10^{11}$	0.03
		$\bar{\nu}_{\mu,\tau} + e^- \rightarrow \bar{\nu}_{\mu,\tau} + e^-$	$9.92 \cdot 10^{31}$	$1.0 \cdot 10^{11}$	0.026
KamLAND	0.2 (C_9H_{12})	$\bar{\nu}_e + p \rightarrow e^+ + n$	$8.55 \cdot 10^{31}$	$2.8 \cdot 10^{11}$	1.6
	0.8 ($C_{12}H_{26}$)	$\nu_e + e^- \rightarrow \nu_e + e^-$	$3.43 \cdot 10^{32}$	$2.8 \cdot 10^{11}$	1.7
		$\bar{\nu}_e + e^- \rightarrow \bar{\nu}_e + e^-$	$3.43 \cdot 10^{32}$	$2.8 \cdot 10^{11}$	0.65
		$\nu_{\mu,\tau} + e^- \rightarrow \nu_{\mu,\tau} + e^-$	$3.43 \cdot 10^{32}$	$1.0 \cdot 10^{11}$	0.11
		$\bar{\nu}_{\mu,\tau} + e^- \rightarrow \bar{\nu}_{\mu,\tau} + e^-$	$3.43 \cdot 10^{32}$	$1.0 \cdot 10^{11}$	0.09
SNO	1.7 (H_2O)	$\bar{\nu}_e + p \rightarrow e^+ + n$	$1.14 \cdot 10^{32}$	$2.8 \cdot 10^{11}$	2.2
	1 (D_2O)	$\bar{\nu}_e + d \rightarrow e^+ + n + n$	$6.00 \cdot 10^{31}$	$2.8 \cdot 10^{11}$	0.004
		$\nu_x + d \rightarrow \nu_x + p + n$	$6.00 \cdot 10^{31}$	$3.8 \cdot 10^{11}$	0.038
		$\bar{\nu}_x + d \rightarrow \bar{\nu}_x + p + n$	$6.00 \cdot 10^{31}$	$3.8 \cdot 10^{11}$	0.032
Super-K	32 (H_2O)	$\bar{\nu}_e + p \rightarrow e^+ + n$	$2.14 \cdot 10^{33}$	$2.8 \cdot 10^{11}$	41
UNO	440 (H_2O)	$\bar{\nu}_e + p \rightarrow e^+ + n$	$2.94 \cdot 10^{34}$	$2.8 \cdot 10^{11}$	560
Hyper-K	540 (H_2O)	$\bar{\nu}_e + p \rightarrow e^+ + n$	$3.61 \cdot 10^{34}$	$2.8 \cdot 10^{11}$	687

PRESUPERNOVA NEUTRINOS: BETA PROCESSES

- Beta processes are also needed
- Many isotopes are present in the late stage of the stellar evolution
- All of these can undergo beta decay or electron/positron capture
- These processes produce only electron flavor neutrinos or antineutrinos

β^-/β^+ Decay

$$A(N, Z) \rightarrow A(N - 1, Z + 1) + e^- + \bar{\nu}_e$$

$$A(N, Z) \rightarrow A(N + 1, Z - 1) + e^+ + \nu_e$$

e^-/e^+ Capture

$$A(N, Z) + e^- \rightarrow A(N + 1, Z - 1) + \nu_e$$

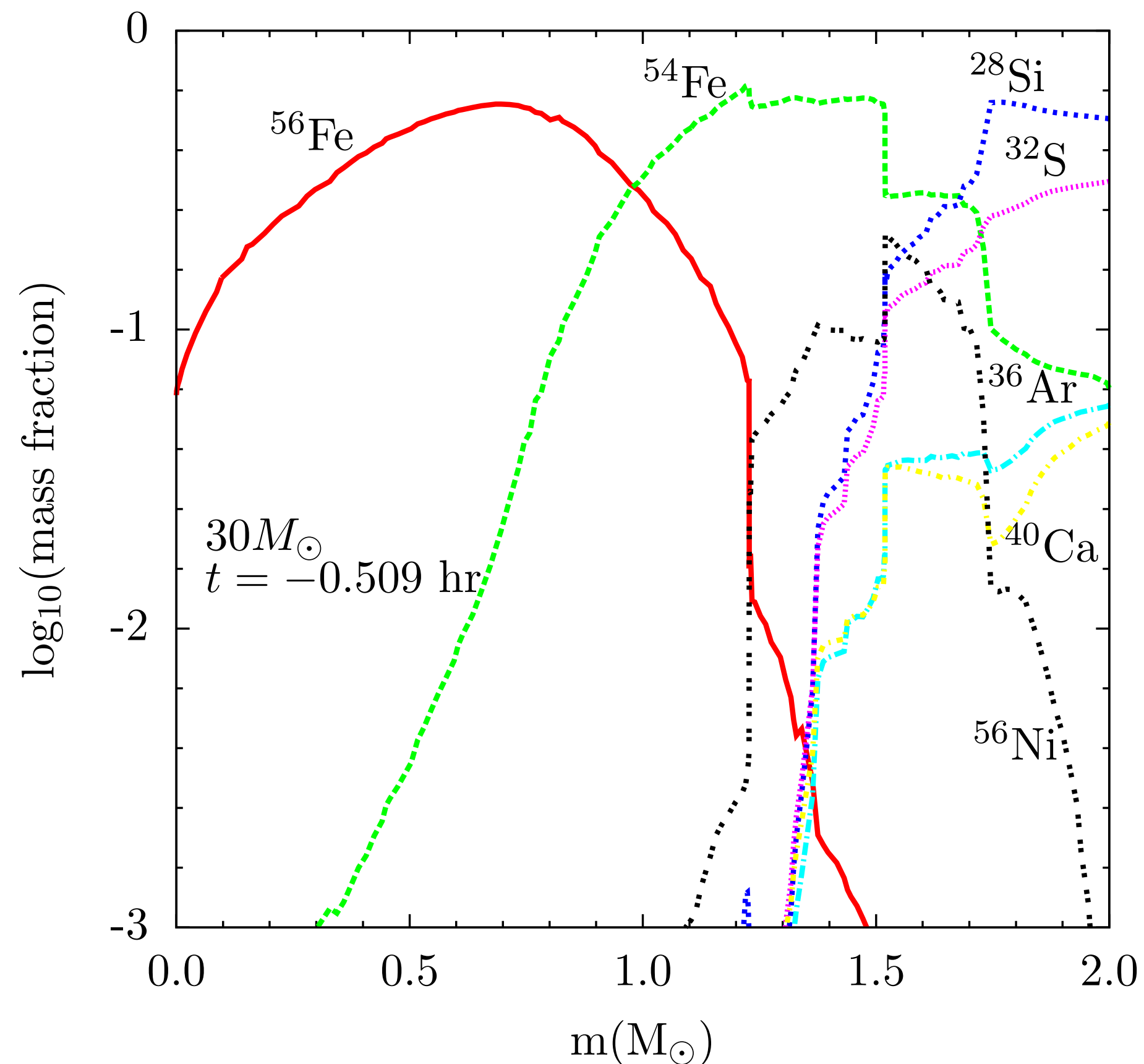
$$A(N, Z) + e^+ \rightarrow A(N - 1, Z + 1) + \bar{\nu}_e$$

N = neutron number

Z = proton number

A = atomic number = $N + Z$

PRESUPERNOVA NEUTRINOS: BETA PROCESSES



Mass fractions for 10 most abundant isotopes

- Since nuclear processes depend so much on which isotopes are present, we have to decide how to calculate the abundances
- We chose to use the stellar evolution code

MESA

- Tracks the abundance of 204 isotopes over the entire lifetime of the star
- Evolves the abundances and nuclear reactions in situ, completely coupled to the hydrodynamics
- Also tracks temperature, density, and electron fraction
- See Farmer et al. arXiv:1611.01207 for details

BETA SPECTRUM

- Shape of spectrum completely determined by phase space of electrons involved
- Depends on chemical potential μ_e , and temperature T , which we get from MESA
- $N_{EC,PC}$ and N_β are normalization factors so our rates match tabulated rates

$$\phi_{EC,PC} = N_{EC,PC} \frac{E_\nu^2 (E_\nu - Q)^2}{1 + \exp((E_\nu - Q - \mu_e)/kT)} \Theta(E_\nu - Q - m_e)$$
$$\phi_\beta = N_\beta \frac{E_\nu^2 (Q - E_\nu)^2}{1 + \exp((E_\nu - Q + \mu_e)/kT)} \Theta(Q - m_e - E_\nu)$$

$$Q_{ij} = M_p - M_d + E_i - E_j$$

- Define “effective Q-value” as the one that reproduces tabulated rates
- Langanke and Martinez-Pinedo, ADNDT **79**, 1 (2001)
- Oda et al., ADNDT **56**, 231 (1994)
- Fuller, Fowler, and Newman, ApJ **252**, 715 (1982)

PRESUPERNOVA NEUTRINOS: BETA PROCESSES

- The total spectrum is found from a weighted sum of all our isotopes

$$\Phi_{\nu, \bar{\nu}} = \sum_k \phi_k n_k = \sum_k X_k \phi_k \frac{\rho}{m_p A_k}$$

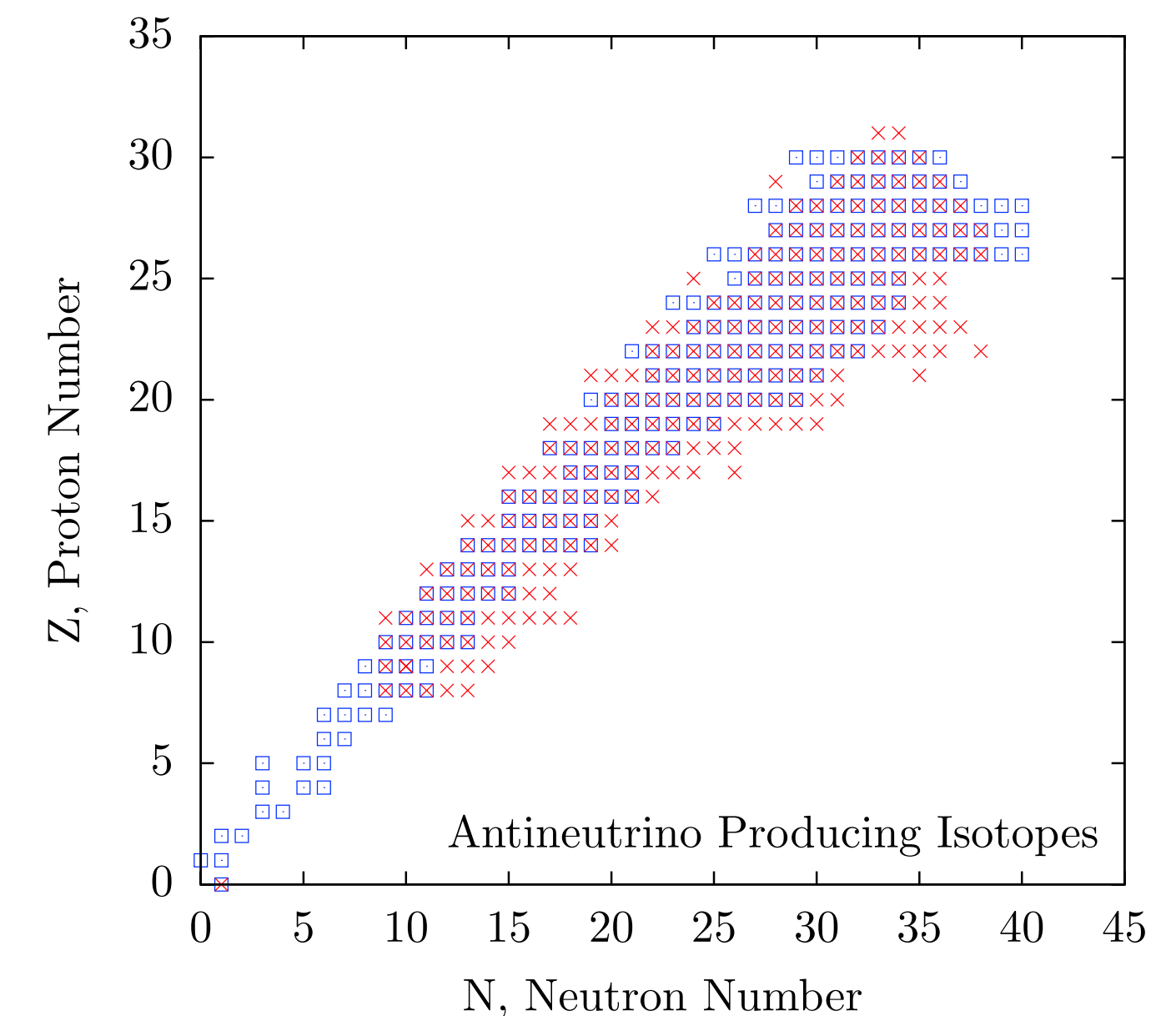
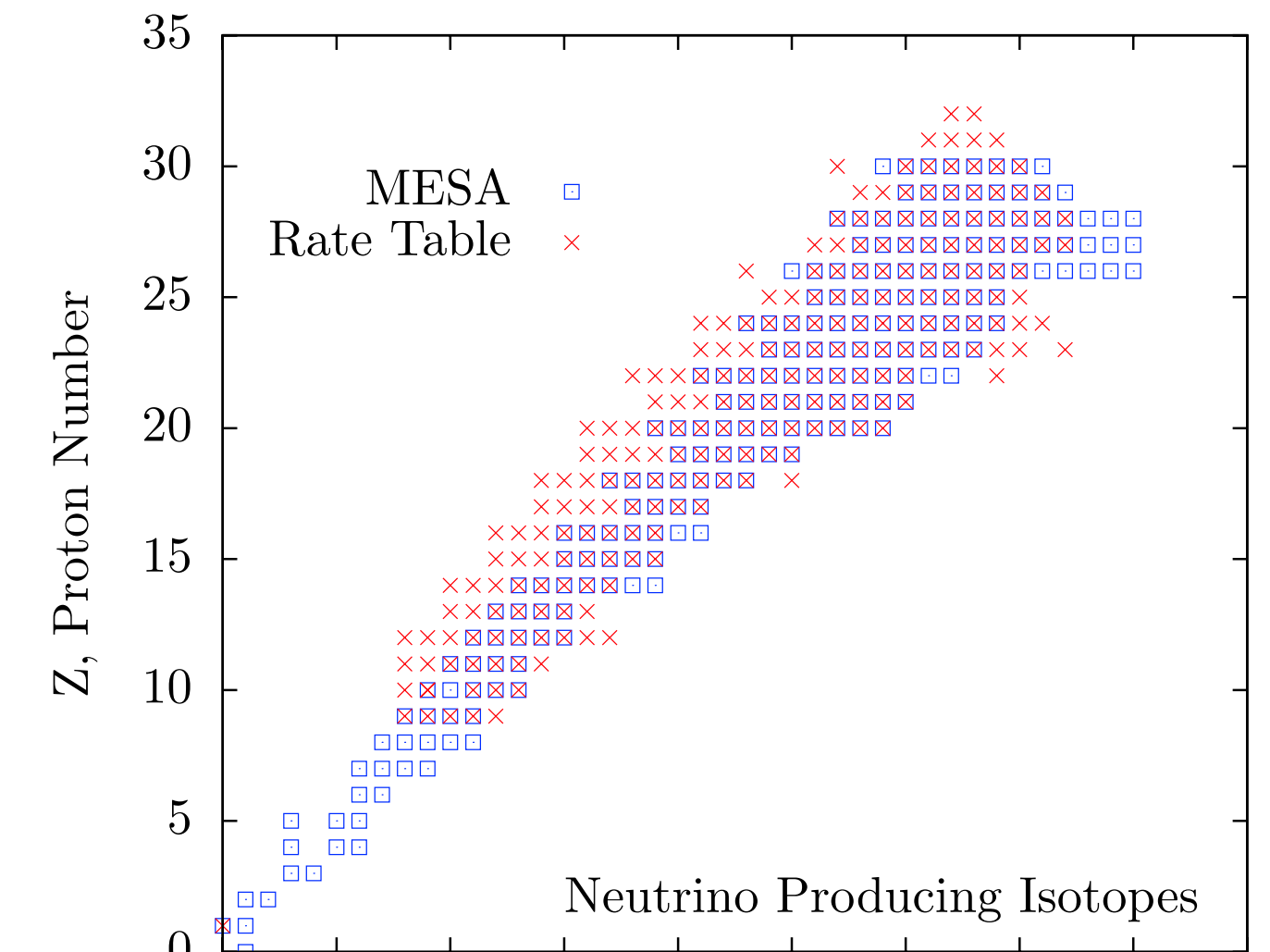
ϕ_k = single isotope spectrum, isotope k

X_k = mass fraction, isotope k

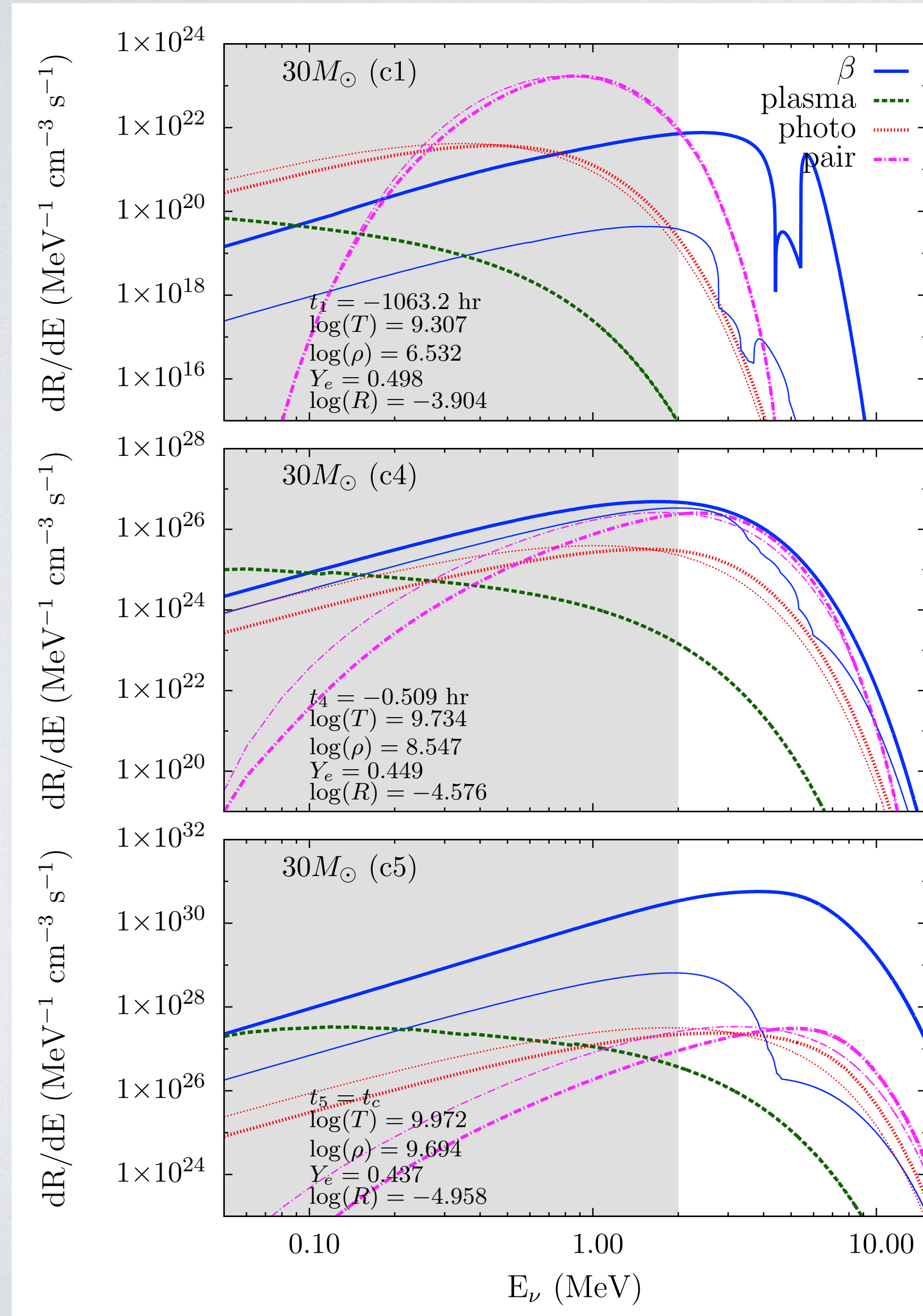
ρ = mass density

m_p = proton mass

A_k = atomic number, isotope k



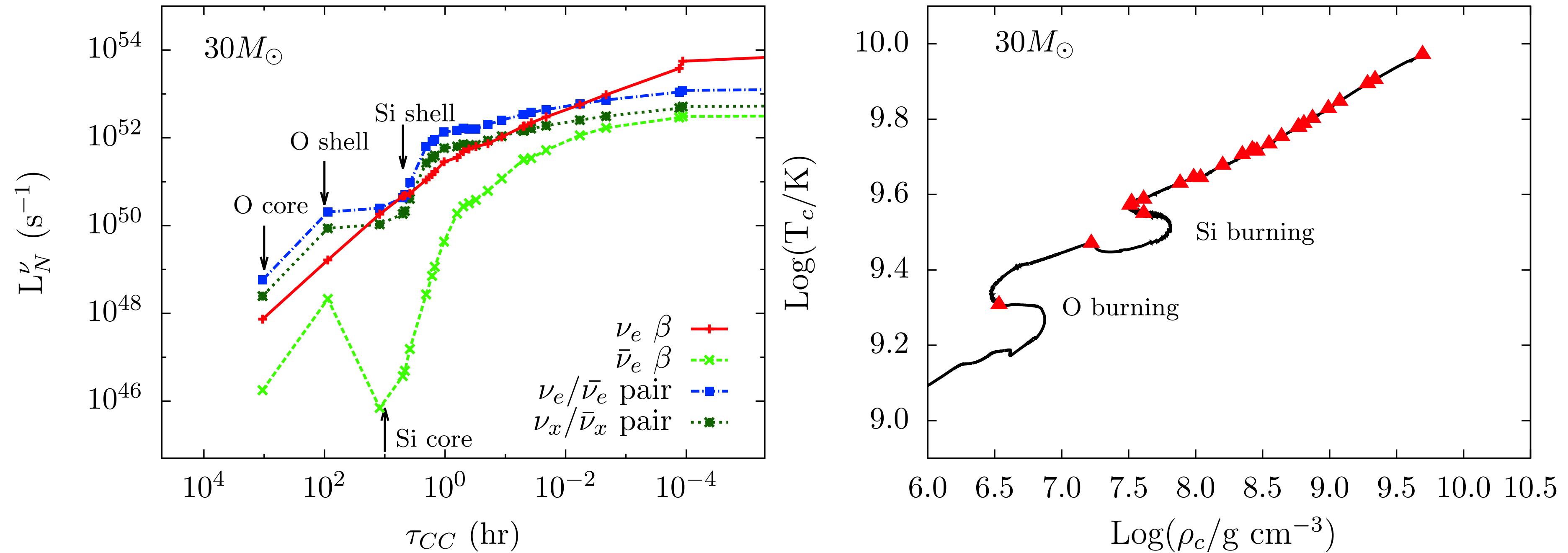
PRESUPERNOVA NEUTRINOS: $30 M_{\odot}$



- Thick lines = ν_e , thin lines = $\bar{\nu}_e$
- Blue = beta, magenta = pair
- Early times, beta spectrum shows a lot of structure and dominates at high (detectable) energies
- Late times, beta spectrum smooths out and dominates overall for neutrinos
- Pair is dominant at high energy for antineutrinos
- Highest contributions from Fe, Co, Mn, and Cr isotopes

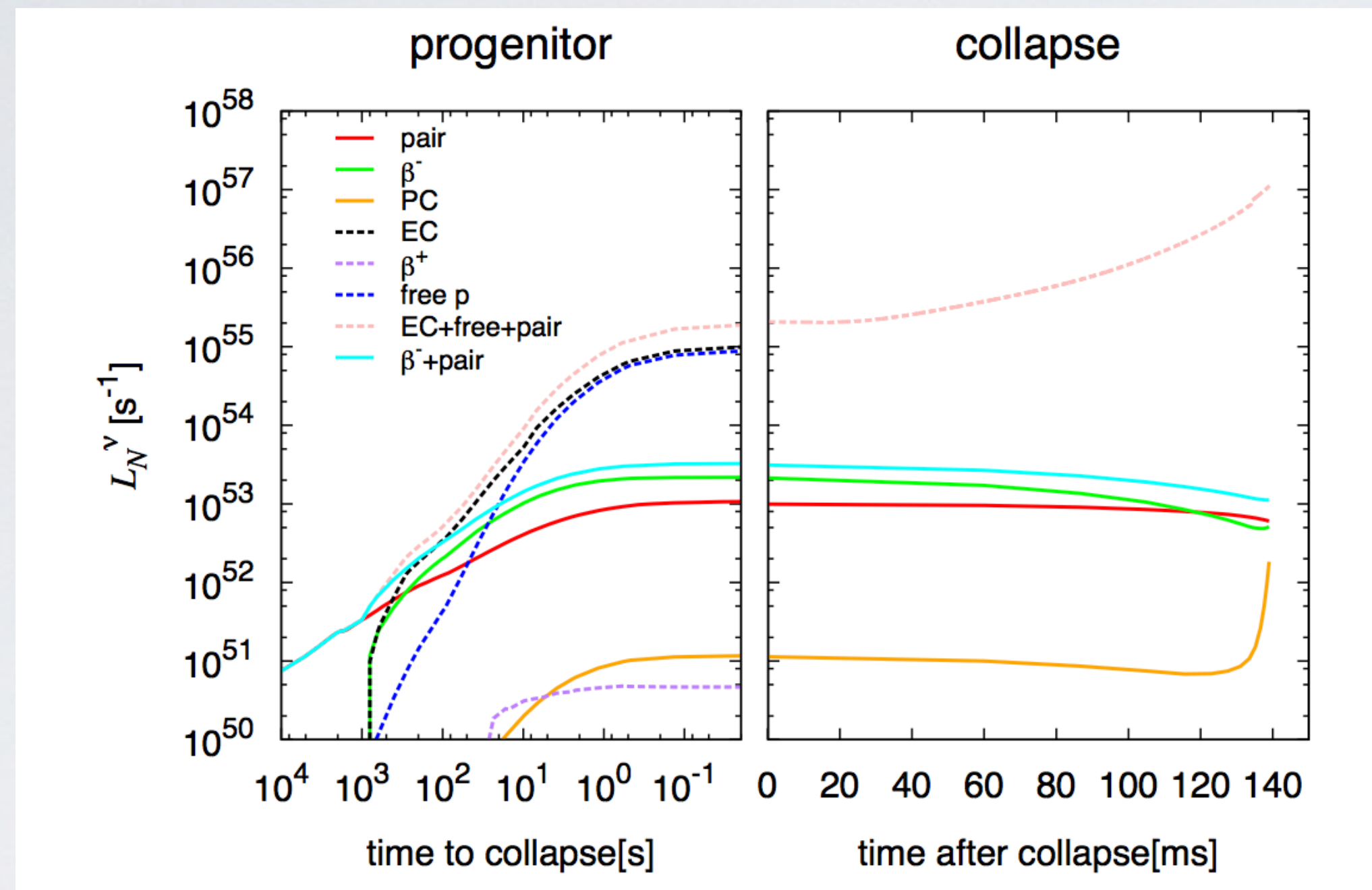
Figure Adapted from KMP, C. Lunardini and R. J. Farmer *ApJ* **840**, 2 (2017)

PRESUPERNOVA NEUTRINOS: IMPORTANCE OF BETA PROCESSES

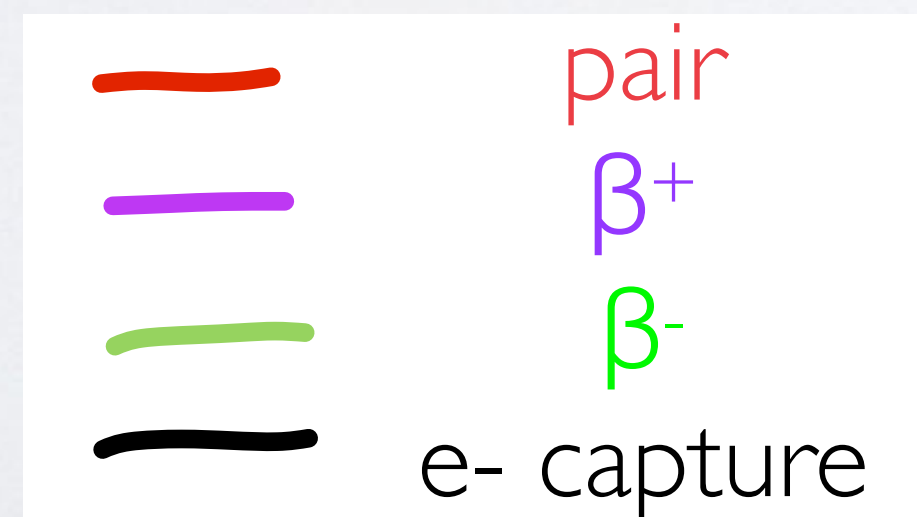


- At early times, pair process dominates
- Later on, in the neutrino channel beta processes rule
- Beta processes for antineutrinos are always subdominant

COMPARISON TO KATO ET AL. (2017)



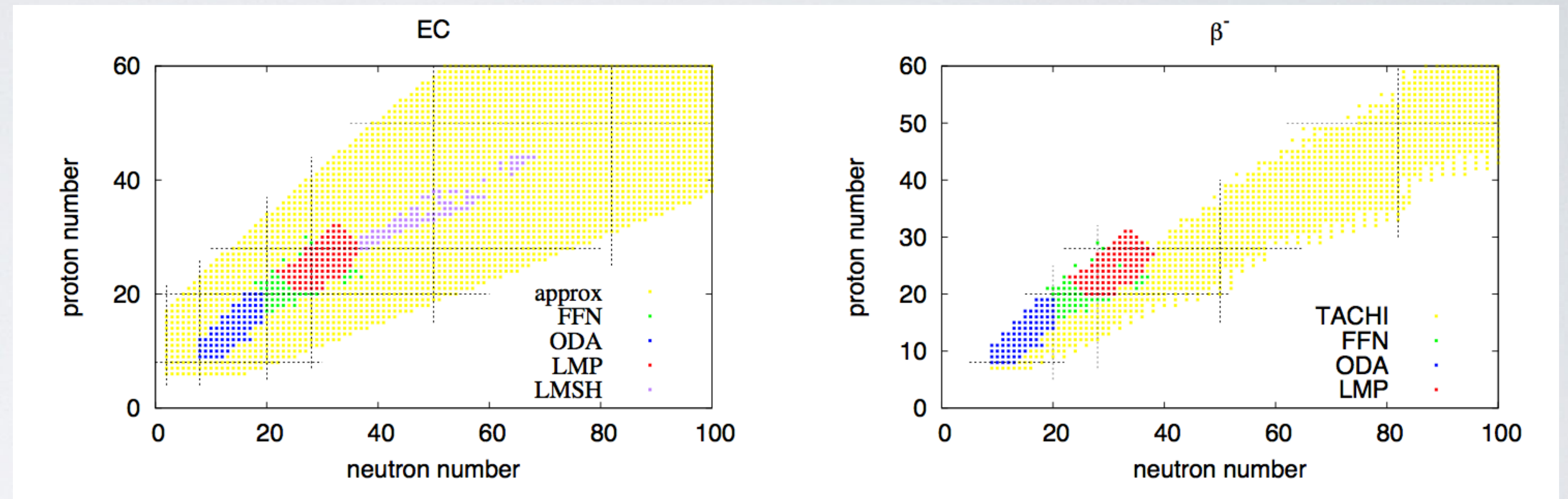
Kato et al. arXiv:1704:05480 (2017)



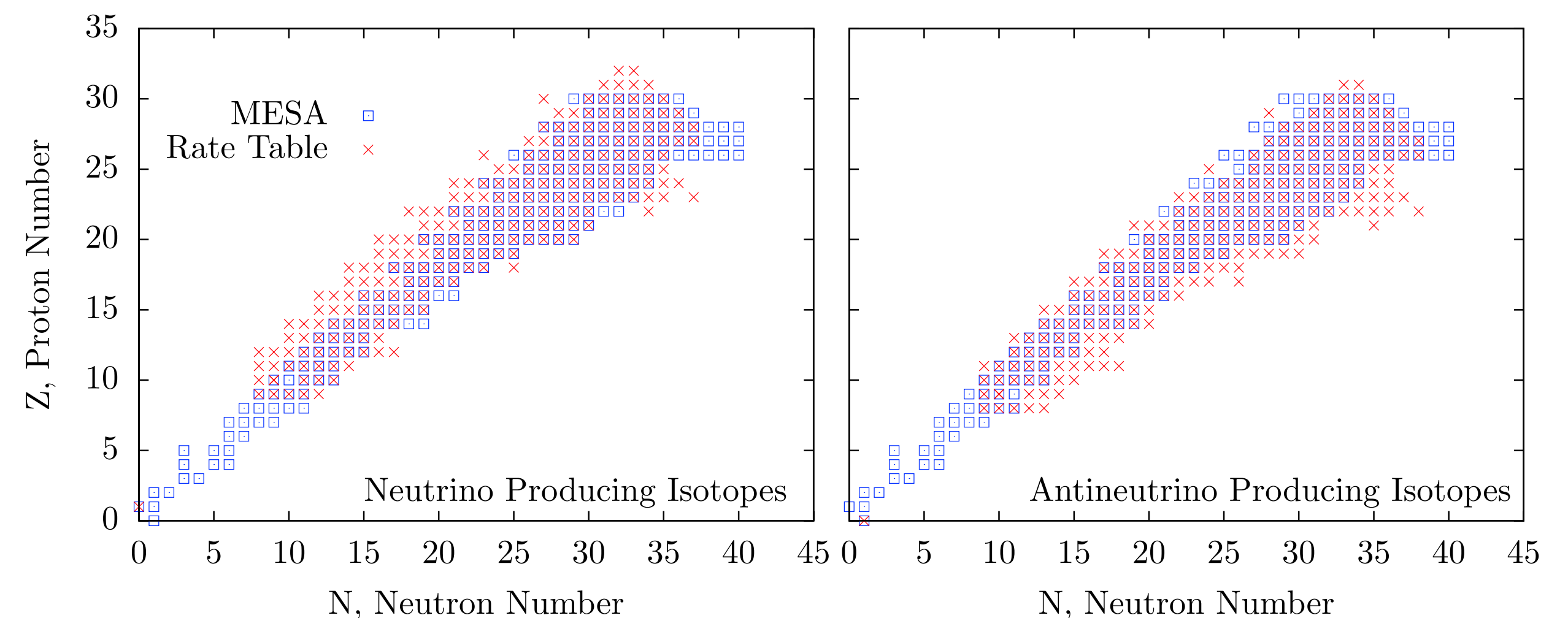
- Kato et al. looked at a $15M_{\odot}$ star
- They found a similar pattern with pair annihilation dominating until shortly before collapse
- One difference: antineutrinos from β^- decay overtake pair neutrinos

COMPARISON TO KATO ET AL. (2017)

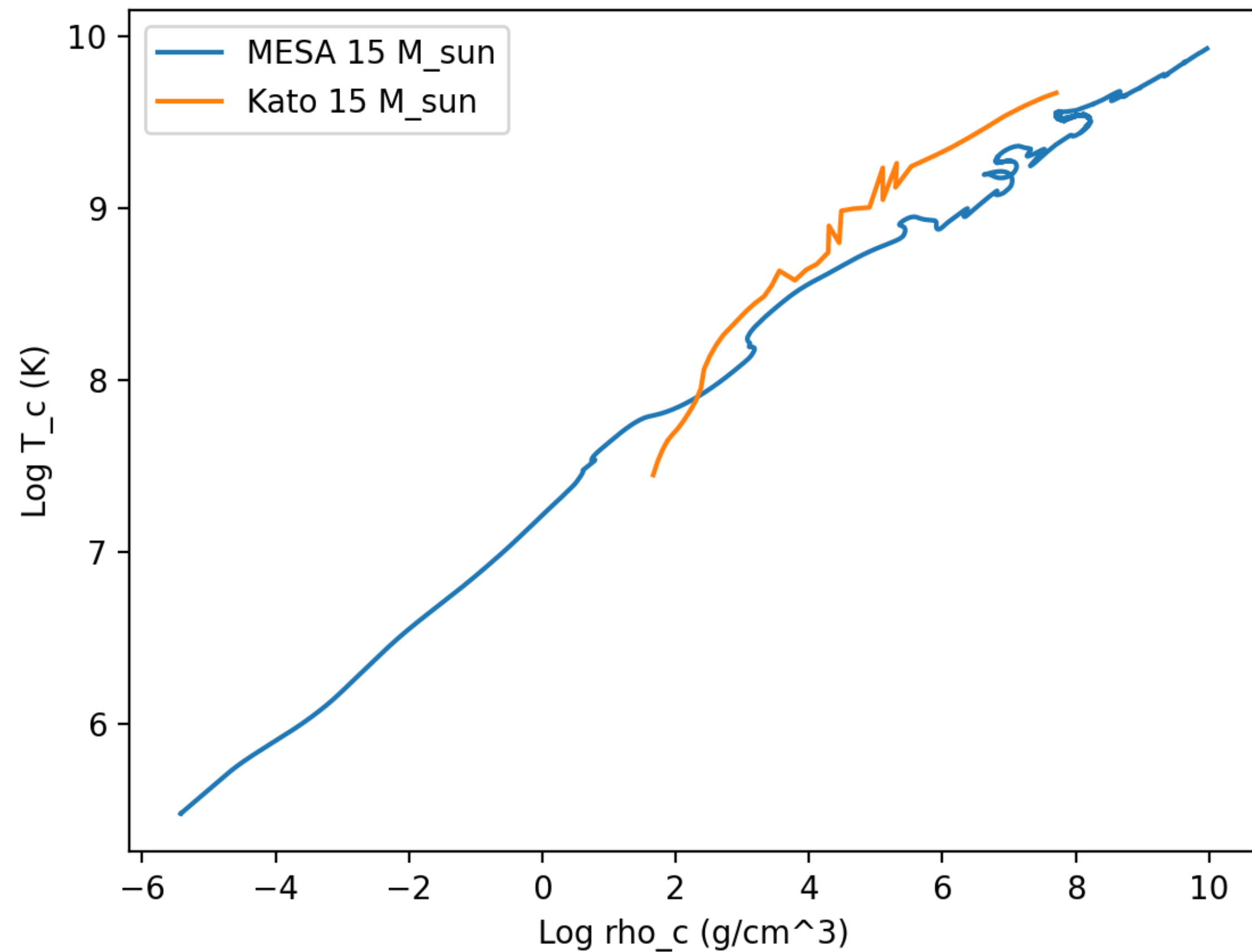
- Expanded isotope table
- Yellow isotopes have terrestrial rates that are incorporated using a suppression factor $1 - f_e(\langle E_e \rangle)$
- More neutron rich isotopes = more electron antineutrinos



Kato et al. arXiv:1704:05480 (2017)



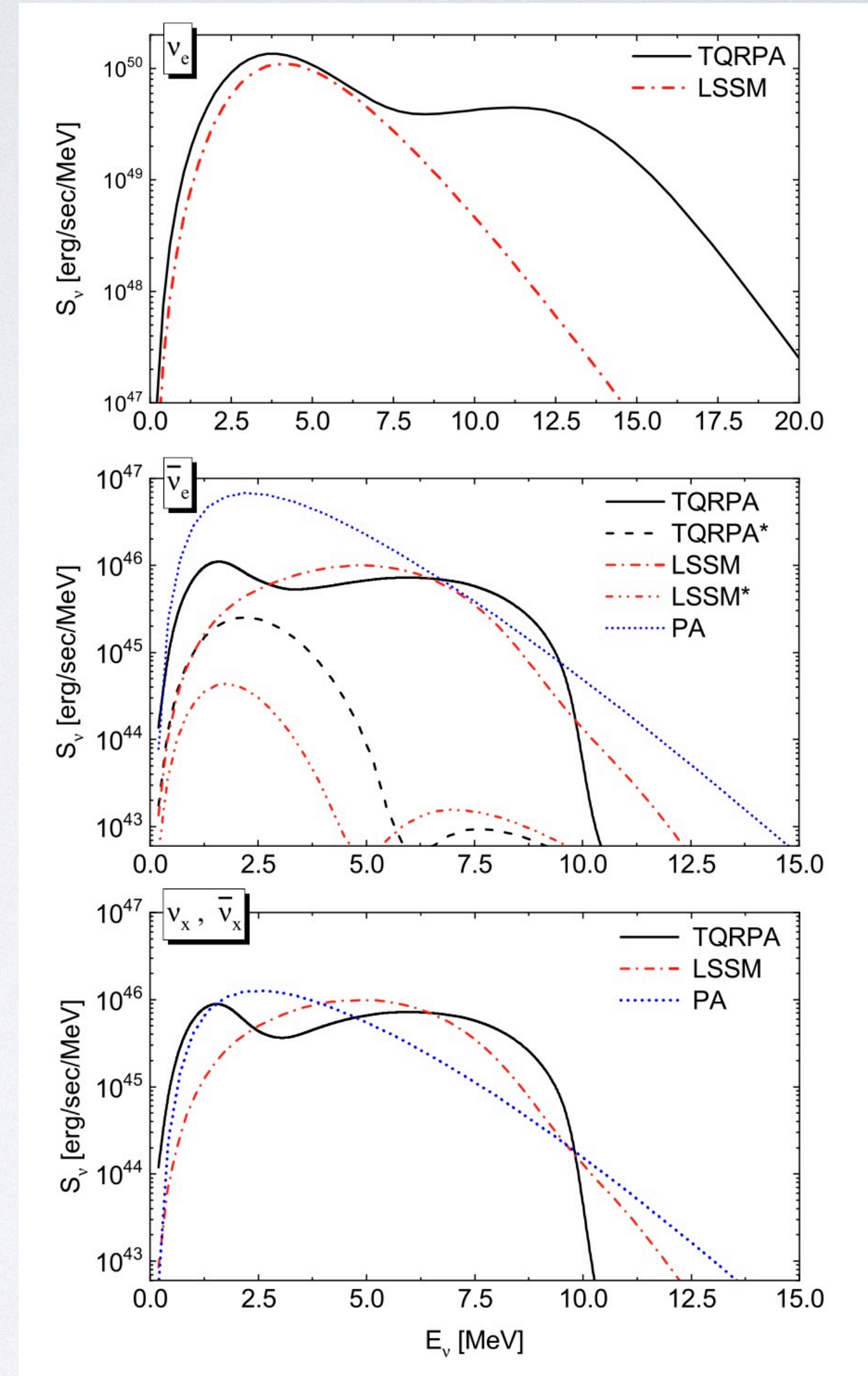
COMPARISON TO KATO ET AL. (2017)



- Comparison of the central temperature and density to collapse also shows differences
- The star used by Kato et al. is generally hotter for a given density
- The MESA model stretches to a higher temperature and density before collapse

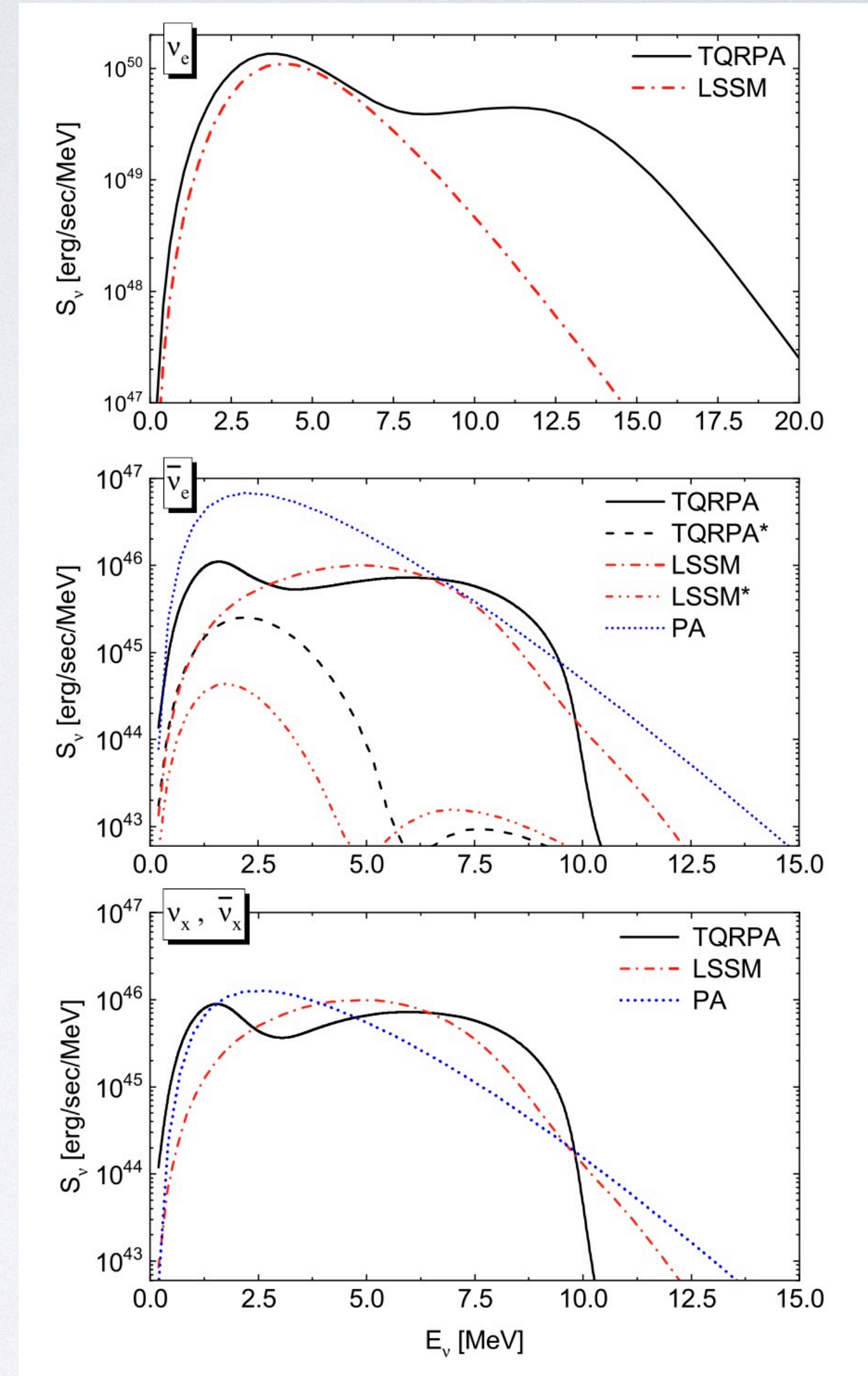
ANOTHER METHOD AND ANOTHER PROCESS: TQRPA AND NEUTRAL CURRENT DE-EXCITATION

- The “effective q-value” method used here could miss important features in the spectrum
- Nuclei could emit neutrino pairs by a neutral current de-excitation process
- Expect these neutrino pairs to be higher energy than from beta processes
- Could we be missing potentially detectable neutrinos?



ANOTHER METHOD AND ANOTHER PROCESS: TQRPA AND NEUTRAL CURRENT DE-EXCITATION

- TQRPA - Thermal Quasiparticle Random Phase Approximation - is another way to compute the beta spectra
- Tends to produce higher energy peaks and sometimes double peaks where “effective q-value” (LSSM) produces one
- Neutral current de-excitation also produces high energy neutrinos
- Shown here only for ^{56}Fe immediately pre-collapse, but an important process to add in!



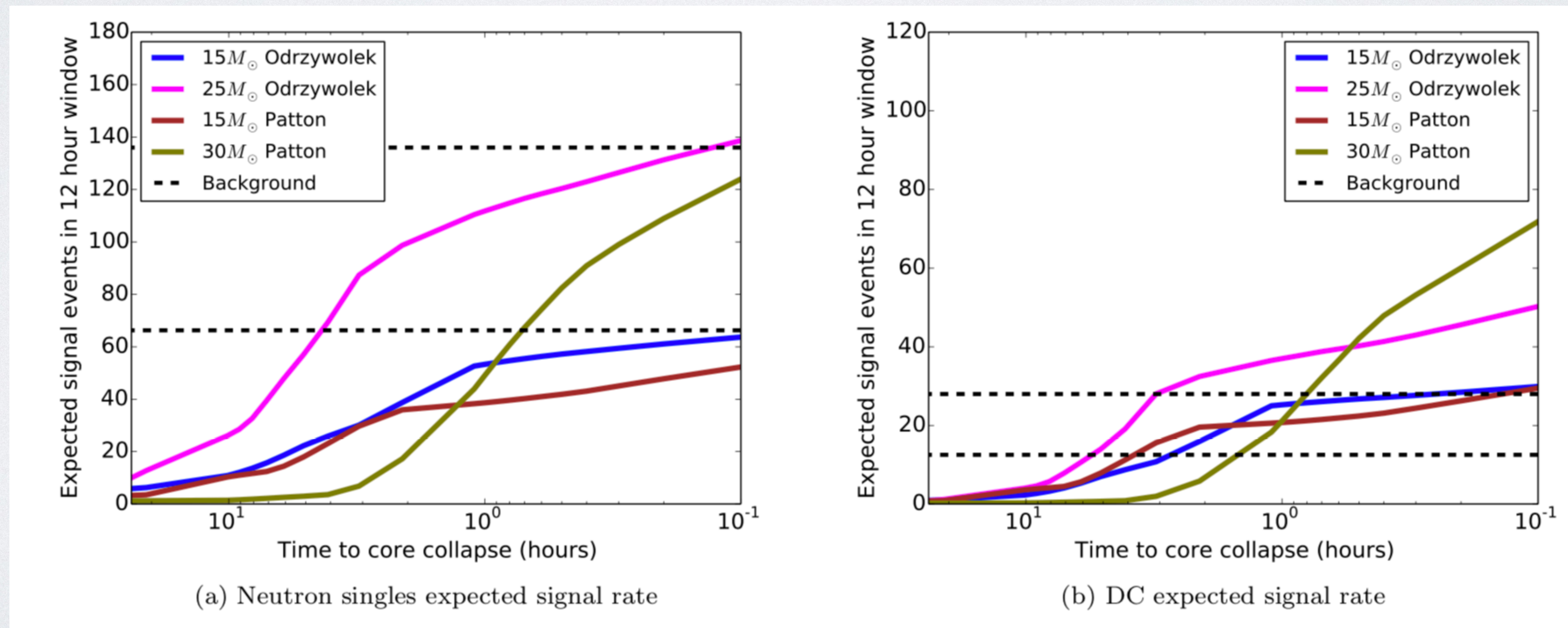
DETECTION OF PRESUPERNOVAE

detector	composition	mass	interval	N_{β}^{CC}	N_{β}^{el}	N^{CC}	N^{el}	$N^{tot} = N^{CC} + N^{el}$
JUNO	C_nH_{2n}	17 kt	$E_e \geq 0.5$ MeV	1.83 [0.05]	4.40 [9.47]	40.1 [13.1]	32.1 [42.7]	72.3 [55.9]
SuperKamiokande	H_2O	22.5 kt	$E_e \geq 4.5$ MeV	0.063 [0.00]	0.053 [0.13]	2.27 [0.78]	0.098 [0.20]	2.37 [0.98]
DUNE	LAr	40 kt	$E \geq 5$ MeV	0.05 [0.76]	0.04 [0.09]	0.19 [1.1]	0.06 [0.13]	0.25 [1.2]

- Consider our $30 M_{\odot}$ star at a distance of $D = 1$ kpc, over the last 2 hours before collapse
- JUNO has the best chance with around 70 events
- If we consider Betelgeuse ($D = 200$ pc), event numbers increase by a factor of 25

DETECTION: SUPER-K WITH GADOLINIUM

- SK-Gd increases detection efficiency because of Gd high neutron capture rate
- Detection coincidence of IBD e^+ and γ -ray cascade after neutron capture on Gd
- Also single neutron events (no IBD e^+): higher background but lower threshold



200 pc
Normal Hierarchy

FUTURE

- More stars!
- Kato et al. have shown that different mass progenitors have different neutrino spectra
- Currently a work in progress
- Add in neutral current de-excitation
- Further comparison with TQRPA and/or update beta calculations with more detailed methods

THANK YOU!

- Thanks to the organizers for the invite
- Thanks to my collaborators: Cecilia Lunardini, Rob Farmer, Frank Timmes
- Thanks to ASU, INT/UW, Colby College, and Trinity College
- Thanks to you for listening!