



Institute for Theoretical Particle Physics

New Physics in $b \rightarrow c \tau \nu$: Impact of Polarisation Observables and $B_c \rightarrow \tau \nu$

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Motivation: the $R_{D^{(*)}}$ anomalies

Test of lepton flavour universality in $b
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u$

$$R_{D^{(*)}} = \frac{\mathcal{BR}(B \to D^{(*)}\tau\nu)}{\mathcal{BR}(B \to D^{(*)}\ell\nu)}$$

- theoretically clean, since hadronic uncertainties largely cancel in ratio
- measured by BaBar, LHCb, Belle →2019: Semileptonic tagging [arXiv:1904.08794]
- $R_{\mu/e}$ agrees with SM \Rightarrow disagreement in τ -channel





| Correlations between observables | | Conclusions |
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Parametrisation of new physics



New physics lies above the scale m_B , so we can parametrise it in terms of four-fermion interactions

$$\begin{aligned} \mathcal{H}_{\text{eff}} =& 2\sqrt{2}G_F V_{\text{cb}}[(1+C_V^L)(\bar{c}\gamma^{\mu}P_Lb)(\bar{\tau}\gamma_{\mu}P_L\nu_{\tau}) + & \bullet \text{ no light } \nu_{\text{R}} \\ &+ C_S^R(\bar{c}P_Rb)(\bar{\tau}P_L\nu_{\tau}) + C_S^L(\bar{c}P_Lb)(\bar{\tau}P_L\nu_{\tau}) & \bullet \text{ NP in } \tau \text{ only} \\ &+ C_T(\bar{c}\sigma^{\mu\nu}P_Lb)(\bar{\tau}\sigma_{\mu\nu}P_L\nu_{\tau})] \end{aligned}$$

Procedure: perform a fit of the Wilson coefficients

- including all available data on the vertex $(\bar{c}\Gamma b)(\bar{\tau}\Gamma\nu_{\tau})$
- restricting to single-particle scenarios

One particle scenarios



| $(C_V^L, C_S^L = -4C_T)$ | Scalar Leptoquark S_1 , $SU(2)$ singlet |
|--------------------------|--|
| (C_S^R, C_S^L) | Charged Higgs |
| (C_V^L, C_S^R) | Vector Leptoquark U_1 , $SU(2)$ singlet |
| $C_S^L = 4 C_T$ | Scalar Leptoquark S ₂ , SU(2) doublet |

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$(\bar{c}\Gamma b)(\bar{\tau}\Gamma \nu_{\tau})$ – Fit



Observables available for the fit

• \mathcal{R}_D

• \mathcal{R}_{D^*}

•
$$au$$
 polarisation in $B \to D^*$:
 $P_{\tau}(D^*) = \frac{\Gamma(\tau^{\lambda=+1/2}) - \Gamma(\tau^{\lambda=-1/2})}{\Gamma(\tau^{\lambda=+1/2}) + \Gamma(\tau^{\lambda=-1/2})}$

•
$$D^*$$
 polarisation: $F_L(D^*) = \frac{\Gamma(D_L^*)}{\Gamma(D^*)}$

Predicted observables

•
$$P_{\tau}(D) = \frac{\Gamma(\tau^{\lambda=+1/2}) - \Gamma(\tau^{\lambda=-1/2})}{\Gamma(\tau^{\lambda=+1/2}) + \Gamma(\tau^{\lambda=-1/2})}$$

• $\mathcal{R}(\Lambda_c) = \frac{\mathrm{BR}(\Lambda_b \to \Lambda_c \tau \nu)}{\mathrm{BR}(\Lambda_b \to \Lambda_c \ell \nu)}$

B_c

 $BR(B_c \rightarrow \tau \nu)$ not measured. We perform the fit requiring

- BR($B_c \rightarrow \tau \nu$) < 10% [Akeroyd, Chen (2017)]
- BR($B_c \rightarrow \tau \nu$) < 30% [Alonso, Grinstein, Martin Camalich (2016)]
- BR($B_c \rightarrow \tau \nu$) < 60% [Conservative limit]

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Fit results



| Mediator | p-value (%) | $\mathcal{R}(D)$ | $\mathcal{R}(D^*)$ | $F_L(D^*)$ | $P_{\tau}(D^*)$ | $P_{\tau}(D)$ | $\mathcal{R}(\Lambda_c)$ |
|--|-------------|------------------|------------------------|-----------------------|------------------------|---------------|--------------------------|
| Charged Higgs _{60%} | 77.4 | 0.333 0.0 σ | 0.299 +0.1 σ | 0.54 -0.7 σ | -0.27 +0.2 <i>σ</i> | 0.38 | 0.38 |
| Charged Higgs _{30%} | 29.9 | 0.348 +0.4 σ | 0.280 -1.2 <i>σ</i> | 0.51 -1.0 <i>σ</i> | -0.35 0.0 <i>σ</i> | 0.41 | 0.37 |
| Charged Higgs _{10%} | 3.2 | 0.360 +0.8 σ | 0.263 -2.2 <i>σ</i> | 0.48 −1.4 <i>σ</i> | -0.44 -0.1 σ | 0.43 | 0.36 |
| Scalar LQ <i>S</i> _{2;60,30%} | 25.0 | 0.333 0.0 σ | 0.297 0.0 σ | 0.45 -1.7 <i>σ</i> | -0.41 -0.1 σ | 0.40 | 0.38 |
| Scalar LQ S _{2;10%} | 7.1 | 0.326 -0.2 σ | 0.276 -1.4 <i>σ</i> | 0.46 -1.6 <i>σ</i> | -0.44 -0.1 σ | 0.38 | 0.36 |

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Correlation: BR($B_c \rightarrow \tau \nu$) and $\mathcal{R}(D^{(*)})$



If the charged Higgs or the scalar LQ S_2 are responsible for the anomaly, we expect BR($B_c \to \tau \nu) > 10\%$

| Mediator | p-value (%) | $\mathcal{R}(D)$ | $\mathcal{R}(D^*)$ | $F_L(D^*)$ | $P_{\tau}(D^*)$ | $P_{\tau}(D)$ | $\mathcal{R}(\Lambda_c)$ |
|------------------------------|-------------|------------------|--------------------|---------------|-----------------|---------------|--------------------------|
| Charged Higgs | 77.4 | 0.333 | 0.299 | 0.54 | -0.27 | 0.38 | 0.38 |
| | | 0.0 σ | $+0.1\sigma$ | -0.7σ | $+0.2\sigma$ | | |
| Charged Higgs _{30%} | 29.9 | 0.348 | 0.280 | 0.51 | -0.35 | 0.41 | 0.37 |
| | 20.0 | $+0.4 \sigma$ | -1.2σ | -1.0σ | 0.0 σ | | |
| Charged Higgs | 3.2 | 0.360 | 0.263 | 0.48 | -0.44 | 0.43 | 0.36 |
| Unarged mgg310% | 0.2 | $+0.8\sigma$ | -2.2σ | -1.4σ | -0.1σ | | |
| Scalar I O S | 25.0 | 0.333 | 0.297 | 0.45 | -0.41 | 0.40 | 0.38 |
| Ocalal LQ 02;60,30% | 20.0 | 0.0 σ | 0.0 σ | -1.7σ | -0.1σ | | |
| Scalar I O S | 7 1 | 0.326 | 0.276 | 0.46 | -0.44 | 0.38 | 0.36 |
| Scalar L& 32;10% | 7.1 | -0.2σ | -1.4σ | -1.6σ | -0.1σ | | |

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New physics scenarios

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Impact of $F_L(D^*)$



The current value of $F_L(D^*)$ favors the charged Higgs scenario

 $F_L(D^*) = 0.60 \pm 0.08 \pm 0.035$ [Belle, 2018]

| Mediator | p-value (%) | $\mathcal{R}(D)$ | $\mathcal{R}(D^*)$ | $F_L(D^*)$ | $P_{\tau}(D^*)$ | $P_{\tau}(D)$ | $\mathcal{R}(\Lambda_c)$ |
|--|-------------|------------------|------------------------|-----------------------|------------------------|---------------|--------------------------|
| Scalar LQ S ₁ | 31.5 | 0.327 -0.2 σ | 0.300 +0.2 <i>σ</i> | 0.47 -1.5 <i>σ</i> | -0.48 -0.2 <i>σ</i> | 0.21 | 0.38 |
| Charged Higgs _{60%} | 77.4 | 0.333 0.0 σ | 0.299 +0.1 σ | 0.54 -0.7 σ | -0.27 +0.2 <i>σ</i> | 0.38 | 0.38 |
| Vector LQ U ₁ | 25.9 | 0.337 +0.1 σ | 0.296 -0.1 σ | 0.46 -1.6 <i>σ</i> | -0.50 -0.2 <i>σ</i> | 0.29 | 0.38 |
| Scalar LQ <i>S</i> _{2;60,30%} | 25.0 | 0.333 0.0 σ | 0.297 0.0 σ | 0.45 -1.7 <i>σ</i> | -0.41 -0.1 σ | 0.40 | 0.38 |

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Polarisation observables





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Polarisation observables





Polarisation observables distinguish new physics scenarios

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Correlation between $\mathcal{R}(\Lambda_c)$ and $\mathcal{R}(D^{(*)})$





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Correlation between $\mathcal{R}(\Lambda_c)$ and $\mathcal{R}(D^{(*)})$





Fitting the current $\mathcal{R}(D^{(*)})$ central values always implies an increase of $\mathcal{R}(\Lambda_c)$

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$\mathcal{R}(\Lambda_c)$ sum rule



The numerical expressions for $\mathcal{R}(\Lambda_c)$ and $\mathcal{R}(D^{(*)})$ lead to the sum rule

$$rac{\mathcal{R}(\Lambda_c)}{\mathcal{R}_{\mathrm{SM}}(\Lambda_c)} = \ 0.262 rac{\mathcal{R}(D)}{\mathcal{R}_{\mathrm{SM}}(D)} + 0.738 rac{\mathcal{R}(D^*)}{\mathcal{R}^{\mathrm{SM}}(D^*)} + x$$

•
$$x \sim \mathcal{O}(0.1(\frac{\Lambda_{\text{EW}}}{\Lambda_{\text{NP}}})^2)$$

 heavy quark limit: R(Λ_c), R(D^(*)) correspond to the (same) branching ratios at the quark level

Standard Model

 $\mathcal{R}_{SM}(\Lambda_c) = 0.33 \pm 0.01$

[Detmold, Lehner, Meinel 2015]

 $\mathcal{R}_{\text{SM}}(\Lambda_c) = 0.324 \pm 0.004$

[Bernlochner, Ligeti, Robinson, Sutcliffe 2018]

New Physics

 $\mathcal{R}_{NP}(\Lambda_c) = 0.38 \pm 0.02 \pm 0.01$

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$\mathcal{R}(\Lambda_c)$ will serve as cross-check of the $\mathcal{R}(D^{(*)})$ measurements

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New physics scenarios

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Summary



- Update of the $b \rightarrow c \tau \nu$ fit, including $F_L(D^*)$ and new Belle data
- Analysis of correlations between observables:
 - BR($B_c \rightarrow \tau \nu$) and $\mathcal{R}(D^{(*)})$: charged Higgs and scalar Leptoquark S_2 predict BR($B_c \rightarrow \tau \nu$) > 10%
 - polarisation observables crucial in distinguishing new physics scenarios
 - $\mathcal{R}(\Lambda_c)$ will serve as cross-check of $\mathcal{R}(D^{(*)})$