FPCP, Victoria, 8th May. 2019

The impact of flavour on high energy searches

Tevong You



The Branco Weiss Fellowship Society in Science

Introduction

- Flavour physics one of the most likely places for indirect signs of new physics to show up
- Complementary to **direct** searches at high energy
- Top-down: flavour structure of models motivated by solutions to SM problems
- **Bottom-up**: models motivated by solutions to flavour anomalies

Top-down flavour implications

Top-down

- Flavour and cosmology See talk by McKeen
- Flavour and dark sectors See talks by Shuve and Robertson
- Neutrino physics See talks by Harris, Ochoa-Ricoux, McFarland, Dunne, Hallin
- Origin of the Higgs sector See talk by Demers
- Top physics See talk by Kareem

Top-down

- Flavour and cosmology See talk by McKeen
- Flavour and dark sectors See talks by Shuve and Robertson
- Neutrino physics See talks by Harris, Ochoa-Ricoux, McFarland, Dunne, Hallin
- Origin of the Higgs sector

See talk by Demers

Top physics

See talk by Kareem

etc.

Composite Higgs model

- Higgs as a (pseudo) Nambu-Goldstone boson
- Confinement of new strong sector at scale f ~ TeV
- Flavour structure from linear mixing with composite operators

$$\mathcal{L} \supset \epsilon_Q \bar{Q}_L \mathcal{O}_Q + \epsilon_U \bar{U}_R \mathcal{O}_U + \epsilon_D \bar{D}_R \mathcal{O}_D. \implies Y_{U(D)})_{ij} \sim \epsilon_Q^i \epsilon_U^j \epsilon_U^j$$

- Requires new sector to be coloured and electroweak-charged ⇒ TeV-scale leptoquark resonances! Gripaios 0910.1789
- O(1) top Yukawa ⇒ composite top couples strongest to new sector
- Minimal composite Higgs model with 4d UVcompletion: SO(6)/SO(5) ⇒ Higgs + singlet scalar

Mъ

GUT

FCNC top decays to new scalar

• *Suppressed* FCNC in **SM**

Banerjee, Chala, Spannowsky 1806.02836

• *Dim-6* Top FCNC in **SMEFT**: $t \rightarrow c S, u S \gg t \rightarrow c h, u h$



BP1(2,3): $\tilde{g} = 1.0(1.0, 0.1), f = 2(10, 2) \text{ TeV} \implies \mathcal{B}(t \to Sc) \sim 10^{-3(4,5)} - 10^{-2(3,4)}.$

FCNC top decays to new scalar

• *Suppressed* FCNC in **SM**

Banerjee, Chala, Spannowsky 1806.02836

• Dim-6 Top FCNC in **SMEFT**: $t \to c S, u S \gg t \to c h, u h$



BP1(2,3): $\tilde{g} = 1.0(1.0, 0.1), f = 2(10, 2) \text{ TeV} \implies \mathcal{B}(t \to Sc) \sim 10^{-3(4,5)} - 10^{-2(3,4)}.$

FCNC top decays to new scalar

• *Suppressed* FCNC in **SM**

Banerjee, Chala, Spannowsky 1806.02836

• *Dim-6* Top FCNC in **SMEFT**: $t \to c S, u S \gg t \to c h, u h$



BP1(2,3): $\tilde{g} = 1.0(1.0, 0.1), f = 2(10, 2) \text{ TeV} \implies \mathcal{B}(t \to Sc) \sim 10^{-3(4,5)} - 10^{-2(3,4)}.$

Bottom-up flavour implications

Bottom-up

See talks by Alonso, Patel, Sevior, Carli, Robinson, Vaquero, Klaver, Waheed, D. Kumar, J. Kumar, Moscati, Malinsky, Wong, Jaeger,

• Anomalies in charged ($B \rightarrow D^{(*)}\mu\nu$) and neutral ($B \rightarrow K^{(*)}\mu^+\mu^-$) current B decays



• *If true,* tremendous implications for **future colliders**

Bottom-up

See talks by Alonso, Patel, Sevior, Carli, Robinson, Vaquero, Klaver, Waheed, D. Kumar, J. Kumar, Moscati, Malinsky, Wong, Jaeger,

• Anomalies in charged ($B \rightarrow D^{(*)}\mu\nu$) and neutral ($B \rightarrow K^{(*)}\mu^+\mu^-$) current B decays



• *If true,* tremendous implications for **future colliders**

Bottom-up

See talks by Alonso, Patel, Sevior, Carli, Robinson, Vaquero, Klaver, Waheed, D. Kumar, J. Kumar, Moscati, Malinsky, Wong, Jaeger,

• Anomalies in charged ($B \rightarrow D^{(*)} \mu \nu$) and neutral ($B \rightarrow K^{(*)} \mu^+ \mu^-$) current B decays



• Focus on **neutral current** B decays for projections

Neutral current B anomalies

 Points towards new physics parametrised by a fourfermion effective operator



New physics behind B anomalies?

• Z' or leptoquarks (at tree-level)



Motivation for future colliders

• Can we *definitely* discover directly the source of the anomalies at higher energies?

80 TeV unitarity limit = **no general no-lose theorem** at FCC-hh (Di Luzio, Nardecchia [1706.01868])

• Consider sensitivity to most **pessimistic** scenario: only include **minimal couplings** required to explain $b \rightarrow s\mu^+\mu^-$ anomalies



• More realistic models will typically be *easier* to discover



• Extrapolate current 13 TeV di-muon search:





• Extrapolate current 13 TeV di-muon search:



Z'

 \overline{s}

Z' Sensitivity

• Extrapolate current 13 TeV di-muon search:



6

5

4

 $g_L^{\mu\mu}$



Fat width







• Extrapolate current 13 TeV di-muon search:



• Extrapolate current 13 TeV di-muon search:



• 100 TeV can cover almost **all** (narrow width) parameter space of most *pessimistic* scenario



• Extrapolate current 13 TeV di-muon search:



• 100 TeV can cover almost **all** (narrow width) parameter space of most *pessimistic* scenario



• Extrapolate current 13 TeV di-muon search:



• 100 TeV can cover almost **all** (narrow width) parameter space of most *pessimistic* scenario

Z

 μ^{\dagger}

 \overline{s}

$$\mathcal{L}_{Z'f} = \left(\overline{\mathbf{Q}'_{\mathbf{L}i}} \lambda_{ij}^{(Q)} \gamma^{\rho} \mathbf{Q}'_{\mathbf{L}j} + \overline{\mathbf{L}'_{\mathbf{L}i}} \lambda_{ij}^{(L)} \gamma^{\rho} \mathbf{L}'_{\mathbf{L}j}\right) Z'_{\rho},$$

$$\mathcal{L}_{Z'f} = \left(\overline{\mathbf{Q}'_{\mathbf{L}i}}\lambda_{ij}^{(Q)}\gamma^{\rho}\mathbf{Q}'_{\mathbf{L}j} + \overline{\mathbf{L}'_{\mathbf{L}i}}\lambda_{ij}^{(L)}\gamma^{\rho}\mathbf{L}'_{\mathbf{L}j}\right)Z'_{\rho},$$
$$V = V_{u_L}^{\dagger}V_{d_L}, \qquad U = V_{\nu_L}^{\dagger}V_{e_L}$$

$$\mathcal{L}_{Z'f} = \left(\overline{\mathbf{Q}'_{\mathbf{L}i}}\lambda_{ij}^{(Q)}\gamma^{\rho}\mathbf{Q}'_{\mathbf{L}j} + \overline{\mathbf{L}'_{\mathbf{L}i}}\lambda_{ij}^{(L)}\gamma^{\rho}\mathbf{L}'_{\mathbf{L}j}\right)Z'_{\rho},$$

$$V = V_{u_{L}}^{\dagger}V_{d_{L}}, \qquad U = V_{\nu_{L}}^{\dagger}V_{e_{L}}$$

$$\mathcal{L} = \left(\overline{\mathbf{u}_{\mathbf{L}}}V\Lambda^{(Q)}V^{\dagger}\gamma^{\rho}\mathbf{u}_{\mathbf{L}} + \overline{\mathbf{d}_{\mathbf{L}}}\Lambda^{(Q)}\gamma^{\rho}\mathbf{d}_{\mathbf{L}} + \overline{\mathbf{n}_{\mathbf{L}}}U\Lambda^{(L)}U^{\dagger}\gamma^{\rho}\mathbf{n}_{\mathbf{L}} + \overline{\mathbf{e}_{\mathbf{L}}}\Lambda^{(L)}\gamma^{\rho}\mathbf{e}_{\mathbf{L}}\right)Z'_{\rho},$$

$$\Lambda^{(Q)} \equiv V_{d_{L}}^{\dagger}\lambda^{(Q)}V_{d_{L}}, \qquad \Lambda^{(L)} \equiv V_{e_{L}}^{\dagger}\lambda^{(L)}V_{e_{L}}$$

$$\begin{aligned} \mathcal{L}_{Z'f} &= \left(\overline{\mathbf{Q}'_{\mathbf{L}i}} \lambda_{ij}^{(Q)} \gamma^{\rho} \mathbf{Q}'_{\mathbf{L}j} + \overline{\mathbf{L}'_{\mathbf{L}i}} \lambda_{ij}^{(L)} \gamma^{\rho} \mathbf{L}'_{\mathbf{L}j} \right) Z'_{\rho}, \\ \mathbf{V} &= V_{u_L}^{\dagger} V_{d_L}, \qquad U = V_{\nu_L}^{\dagger} V_{e_L} \\ \mathcal{L} &= \left(\overline{\mathbf{u}_{\mathbf{L}}} V \Lambda^{(Q)} V^{\dagger} \gamma^{\rho} \mathbf{u}_{\mathbf{L}} + \overline{\mathbf{d}_{\mathbf{L}}} \Lambda^{(Q)} \gamma^{\rho} \mathbf{d}_{\mathbf{L}} + \overline{\mathbf{n}_{\mathbf{L}}} U \Lambda^{(L)} U^{\dagger} \gamma^{\rho} \mathbf{n}_{\mathbf{L}} + \overline{\mathbf{e}_{\mathbf{L}}} \Lambda^{(L)} \gamma^{\rho} \mathbf{e}_{\mathbf{L}} \right) Z'_{\rho}, \\ \Lambda^{(Q)} &\equiv V_{d_L}^{\dagger} \lambda^{(Q)} V_{d_L}, \qquad \Lambda^{(L)} \equiv V_{e_L}^{\dagger} \lambda^{(L)} V_{e_L} \end{aligned}$$

The 'mixed up-muon' (MUM) model

$$\Lambda^{(Q)} = g_{bs} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, \qquad \Lambda^{(L)} = g_{\mu\mu} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

The 'mixed down-muon' (MDM) model

$$\Lambda^{(Q)} = g_{tt} V^{\dagger} \cdot \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \cdot V, \qquad \Lambda^{(L)} = g_{\mu\mu} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Allanach, Corbett, Dolan, TY [1810.02166]

• Improved MC study including **large widths** and **two benchmark flavour scenarios**:



Allanach, Corbett, Dolan, TY [1810.02166]

• **Indirect effects** from effective operators in LHC di-muon tail would point towards **fat Z'** at higher energies



9000

Leptoquark Sensitivity

• Extrapolate current 8 TeV LQ di-muon+di-jet search: group



- Pair production for scalar LQ depends only on QCD coupling
- Upper limit from Bs mixing constraint

 $\partial L \hat{Q}$

Leptoquark Sensitivity

• Extrapolate current 8 TeV LQ di-muon+di-jet search: growthe search: growthe search with the search of the search



- Pair production for scalar LQ depends only on QCD coupling
- Upper limit from Bs mixing constraint

Leptoquark single production







Take-Home Message: Z' sensitivity

- Drell-Yan, $p p \rightarrow Z' \rightarrow \mu^+ \mu^-$ for two flavour assumptions:
 - Mixed-Up Model:
 - HL-LHC: *No sensitivity*
 - **HE-LHC**: $M_{Z'}$ up to 13 TeV, width up to 30%
 - **FCC-hh**: M_Z , up to 22 TeV, width up to 30%
 - Mixed-Down Model:
 - HL-LHC: M_Z , up to 5 TeV, width up to 10%
 - HE-LHC: $M_{Z'}$ up to 10 TeV, width up to 60%
 - FCC-hh: *M_Z*, up to 20 TeV *(entire parameter space)*

Take-Home Message: LQ sensitivity

• Pair production, $p p \rightarrow LQ LQ \rightarrow \mu^+ \mu^- j j$



• Single production, $p p \rightarrow LQ \rightarrow \mu^+ \mu^- j$



Conclusion

- If confirmed by LHCb and Belle II...
- B anomalies could shed light on many BSM questions
- Accessible scale of new physics
- First studies of direct search potential at future colliders
- Even if anomalies vanish, motivates *interplay* between **direct** discovery potential of future hadron colliders and **indirect** sensitivity from precision physics