## Combined explanation of the B -anomalies

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Presented by Jacky Kumar
Université de Montréal
Based on JK, David London, Ryoutaro Watanabe,
Phys. Rev. D 99, 015007


## B-Anomalies

- Discrepancies in $b \rightarrow s \mu^{+} \mu^{-}$data and SM: Angular Observables in $B \rightarrow K^{*} \mu^{+} \mu^{-}$, Branching Ratio in $B_{s} \rightarrow \phi \mu^{+} \mu^{-}$: (Combined Significance 4-5 $\sigma$.


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- Discrepancies in Lepton Flavor Universality Ratios in $b \rightarrow s \ell$ :

$$
\mathbf{R}_{\mathbf{K}^{(*)}}=\frac{\mathcal{B}\left(\mathbf{B} \rightarrow \mathbf{K}^{(*)} \mu^{+} \mu^{-}\right)}{\mathcal{B}\left(\mathbf{B} \rightarrow \mathbf{K}^{(*)} \mathbf{e}^{+} \mathbf{e}^{-}\right)}
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- Discrepancies in Lepton Flavor Universality Ratios in $b \rightarrow c \nmid \bar{\nu}$ :

$$
\mathbf{R}_{\mathbf{D}^{(*)}}=\frac{\mathcal{B}\left(\mathbf{B} \rightarrow \mathbf{D}^{(*)} \tau \bar{\nu}\right)}{\mathcal{B}\left(\mathbf{B} \rightarrow \mathbf{D}^{(*)} \ell \bar{\nu}_{\ell}\right)}
$$

$$
\mathbf{R}_{\mathbf{J} / \psi}=\frac{\mathcal{B}\left(\mathbf{B}_{\mathbf{c}} \rightarrow \mathbf{J} / \psi \tau \bar{\nu}\right)}{\mathcal{B}\left(\mathbf{B}_{\mathbf{c}} \rightarrow \mathbf{J} / \psi \mu \bar{\nu}_{\mu}\right)}
$$

Both $R_{D}$ and $R_{D^{*}}$ are measured to be above the SM value, the combined significance is $\sim 4.0 \sigma . R_{J / \psi}$ is measured to be $\sim \mathbf{2} \sigma$ above the SM.

## Individual Explanations: EFT Approach

- The NP can be parameterized in terms of the Wilson Coefficients.

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\begin{gathered}
\mathcal{H}_{\mathrm{eff}}=\sum C_{i} O_{i} \\
C_{X}=C_{X}(\mathrm{SM})+C_{X}(\mathrm{NP})
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$b \rightarrow s \mu \mu:$

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$\mathbf{b} \rightarrow \mathbf{c} \tau \bar{\nu}_{\tau}:$

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$$

(Global Fits):

- $\mathrm{b} \rightarrow \mathrm{s} \mu \mu$ :

$$
C_{9}^{\mu \mu}(N P)=-C_{10}^{\mu \mu}(N P) \simeq-0.53 .
$$

- $\mathrm{b} \rightarrow \mathrm{c} \tau \bar{\nu}$ :

$$
C_{V}^{\tau \tau}(N P) \simeq 0.10
$$

See e.g. Phys. Rev. D 96, 095009, JHEP 1809 (2018) 152

## Combined Explanation: EFT Approach

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$$
\left(\bar{Q}_{i L} \gamma_{\mu} Q_{j L}\right)\left(\bar{L}_{k L} \gamma^{\mu} L_{l L}\right), \quad\left(\bar{Q}_{i L} \gamma_{\mu} \sigma^{I} Q_{j L}\right)\left(\bar{L}_{k L} \gamma^{\mu} \sigma^{I} L_{l L}\right)
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Operator $\left(\bar{Q}_{i L} \gamma_{\mu} \sigma^{I} Q_{j L}\right)\left(\bar{L}_{k L} \gamma^{\mu} \sigma^{I} L_{l L}\right)$ relates $b \rightarrow s \ell$ to $b \rightarrow c \ell \bar{\nu}$ transitions.
[Phys.Lett. B742 (2015) 370-374]

## EFT to Models: Leptoquarks

Scalar Triplet: $S_{3}(\mathbf{3}, \mathbf{3},-\mathbf{2} / \mathbf{3})$
Vector Triplet: $U_{3}(\mathbf{3}, \mathbf{3}, \mathbf{4} / \mathbf{3})$
Vector Singlet: $U_{1}(\mathbf{3}, \mathbf{1}, \mathbf{4} / \mathbf{3})$


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$$
\begin{aligned}
\Delta \mathcal{L}_{S_{3}} & =h_{i j}^{S_{3}}\left(\bar{Q}_{i L} \sigma^{I} i \sigma^{2} L_{j L}^{c}\right) S_{3}^{I}+\text { h.c., } \quad \text { (We allow General Couplings) } \\
\Delta \mathcal{L}_{U_{3}} & =h_{i j}^{U_{3}}\left(\bar{Q}_{i L} \gamma^{\mu} \sigma^{I} L_{j L}\right) U_{3 \mu}^{I}+\text { h.c. }, \\
\Delta \mathcal{L}_{U_{1}} & =h_{i j}^{U_{1}}\left(\bar{Q}_{i L} \gamma^{\mu} L_{j L}\right) U_{1 \mu}+\text { h.c. }
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\Delta \mathcal{L}_{U_{1}} & =h_{i j}^{U_{1}}\left(\bar{Q}_{i L} \gamma^{\mu} L_{j L}\right) U_{1 \mu}+\text { h.c. }
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$$

Under the assumption that NP Couples to only II and III Generations we have 4 Free(Real) parameters for each Model:

$$
h_{22}, h_{33}, h_{23}, h_{32}
$$

## Observables

## Six Minimal + Five Lelpton Flavor Violating (LFV) constraints.

| Observable | Measurement or Constraint |
| :---: | :---: |
| minimal |  |
| $\begin{gathered} b \rightarrow s \mu^{+} \mu^{-}(\text {all }) \\ R_{D}^{\tau *} /\left(R_{D^{*}}^{\tau / \ell}\right)_{\mathrm{SM}} \\ R_{D}^{\tau / \ell} /\left(R_{D}^{\tau / \ell}\right)_{\mathrm{SM}} \\ R_{D}^{e / \mu} /\left(R_{D^{*}}^{e / \mu}\right)_{\mathrm{SM}} \\ R_{J / \psi}^{\tau / \mu} /\left(R_{J / \psi}^{\tau / \mu}\right)_{\mathrm{SM}} \\ \mathcal{B}\left(B \rightarrow K^{(*)} \nu \bar{\nu}\right) / \mathcal{B}\left(B \rightarrow K^{(*)} \nu \bar{\nu}\right)_{\mathrm{SM}} \end{gathered}$ | $\begin{aligned} & C_{9}^{\mu \mu}(\mathrm{LQ})=-C_{10}^{\mu \mu}(\mathrm{LQ})=-0.68 \pm 0.12[17] \\ & 1.18 \pm 0.06[18-21] \\ & 1.36 \pm 0.15[18-21] \\ & 1.04 \pm 0.05[68] \\ & 2.51 \pm 0.97[22] \\ &-13 \sum_{i=1}^{3} \operatorname{Re}\left[C_{L}^{i i}(\mathrm{LQ})\right]+\sum_{i, j=1}^{3}\left\|C_{L}^{i j}(\mathrm{LQ})\right\|^{2} \leq 248[69] \end{aligned}$ |
| LFV |  |
| $\begin{gathered} \hline \mathcal{B}\left(B^{+} \rightarrow K^{+} \tau^{-} \mu^{+}\right) \\ \mathcal{B}\left(B^{+} \rightarrow K^{+} \tau^{+} \mu^{-}\right) \\ \mathcal{B}\left(\Upsilon(2 S) \rightarrow \mu^{ \pm} \tau^{\mp}\right) \\ \mathcal{B}(\tau \rightarrow \mu \phi) \\ \mathcal{B}\left(J / \psi \rightarrow \mu^{ \pm} \tau^{\mp}\right) \\ \hline \end{gathered}$ | $\begin{array}{\|c} (0.8 \pm 1.7) \times 10^{-5} ; \end{array} \quad<4.5 \times 10^{-5}(90 \% \text { C.L. })[70] ~\left[\begin{array}{rl} \hline(-0.4 \pm 1.2) \times 10^{-5} ; & <2.8 \times 10^{-5}(90 \% \text { C.L. })[70] \\ (0.2 \pm 1.5 \pm 1.3) \times 10^{-6} ; & <3.3 \times 10^{-6}(90 \% \text { C.L. })[71] \\ <8.4 \times 10^{-8}(90 \% \text { C.L. })[72] \\ <2.0 \times 10^{-6}(90 \% \text { C.L. })[73] \\ \hline \end{array}\right.$ |

## $S_{3}$ and $U_{3}$ Leptoquarks Models

The Fit of $S_{3}$ and $U_{3}$ to the Minimal set of Constraints yields:

$$
\chi^{2} / \text { dof }=7.5\left(S_{3}\right), \quad 10\left(U_{3}\right),
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- For the $S_{3}$ LQ, we have

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\begin{gathered}
h_{33} h_{23}=-0.28 \pm 0.08\left(R_{D^{(*)}}\right) \\
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- Similarly, the $U_{3}$ LQ has

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& h_{33} h_{23}=-0.14 \pm 0.04\left(R_{D^{(*)}}\right) \\
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The constraint from $B \rightarrow K^{(*)} \nu \bar{\nu}$ is not compatible with the $R_{D^{(*)}}$.

## $U_{1}$ Leptoquark Model

- No contributions to $b \rightarrow s \nu \bar{\nu}$ (at Tree Level) since

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9 Observables:

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\text { Minimal : } & b \rightarrow s \mu \mu, R_{D}^{\tau / / \ell}, R_{D^{*}}^{\tau / \ell}, R_{D^{*}}^{e / \mu}, R_{J / \psi}^{\tau / \mu} \\
\text { LFV : } & \mathcal{B}\left(B \rightarrow K \tau^{ \pm} \mu^{\mp}\right), \mathcal{B}(\tau \rightarrow \phi \mu), \mathcal{B}\left(\Upsilon \rightarrow \mu^{ \pm} \tau^{\mp}\right) .
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\end{aligned}
$$

4 Free Parameters:

$$
\begin{gathered}
h_{22}, h_{33}, h_{23}, h_{32} \Longrightarrow \text { d.o.f }=5 . \\
\chi_{\min }^{2} / \operatorname{dof}=1
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Therefore $U_{1}$ LQ can explain both the charged and neutral current B-anomalies simultaneously.

## LQ Couplings: Pattern \& LFV Constraints

Using Minimal Observables only product of LQ couplings are constrained but the individual couplings remain unconstrained.

$$
\begin{array}{rll}
b \rightarrow s \mu^{+} \mu^{-} & : h_{32} h_{22} \\
b \rightarrow c \tau \bar{\nu} & : V_{c s} h_{33} h_{23}+V_{c b} h_{33}^{2} \\
B^{+} \rightarrow K^{+} \tau^{-} \mu^{+} & : & h_{32} h_{23} \\
B^{+} \rightarrow K^{+} \tau^{+} \mu^{-} & : & h_{33} h_{22} \\
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\tau \rightarrow \mu \phi & : & h_{23} h_{22} .
\end{array}
$$

Lepton Flavor Violating Observables put additional constraints:

$$
\begin{aligned}
& |h 22| \leq 0.12,|h 32| \leq 0.7 \\
& |h 23| \leq 0.9,|h 33| \geq 0.1 .
\end{aligned}
$$




## LQ Couplings: Pattern \& LFV Constraints

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\begin{array}{ccc}
R_{D^{(*)}} & R_{K^{(*)}} \\
A=(a, c): h_{33}=O(1.0), & h_{23}=O(0.1), & h_{32}=O(0.01), \quad h_{22}=O(0.1) \\
B=(b, c): h_{33}=O(0.1), & h_{23}=O(1.0), & h_{32}=O(0.01), \\
C=(a, d): h_{22}=O(0.1) \\
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\end{array}
$$

| $h_{33}$ | $\chi_{\min , S M+U_{1}}^{2}$ | $h_{23}$ |
| :---: | :---: | :---: |
| 1.0 | 5.0 | $0.10 \pm 0.04$ |
| 0.5 | 5.2 | $0.26 \pm 0.07$ |
| 0.2 | 6.8 | $0.60 \pm 0.15$ |
| 0.1 | 11.3 | $0.70 \pm 0.20$ |

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LFV constraints prefer a large value of $h_{33}$ coupling. A sizable $h_{23} \sim \mathcal{O}(0.1)$ is needed to fit the data.

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$$

| $h_{33}$ | $\chi_{\min , S M+U_{1}}^{2}$ | $h_{23}$ |
| :---: | :---: | :---: |
| 1.0 | 5.0 | $0.10 \pm 0.04$ |
| 0.5 | 5.2 | $0.26 \pm 0.07$ |
| 0.2 | 6.8 | $0.60 \pm 0.15$ |
| 0.1 | 11.3 | $0.70 \pm 0.20$ |

LFV constraints prefer a large value of $h_{33}$ coupling. A sizable $h_{23} \sim \mathcal{O}(0.1)$ is needed to fit the data.



## Predictions for $U_{1}$ model

- Enhancement of same size in $b \rightarrow u \ell \bar{\nu}$ modes is predicted:

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R_{\pi \ell \bar{\nu}}^{\tau / \mu}=\frac{\mathcal{B}(B \rightarrow \pi \tau \bar{\nu})}{\mathcal{B}(B \rightarrow \pi \ell \bar{\nu})} \simeq R_{D^{*}}^{\tau / \ell} \simeq 1.20 .
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- More than two orders of enhancement is expected in the $b \rightarrow s \tau \tau$ modes!

$$
\mathcal{B}(B \rightarrow K \tau \bar{\tau}) \simeq 250 \times \mathcal{B}(B \rightarrow K \tau \bar{\tau})_{S M}
$$



## Vector Boson (VB) Triplet Model

- An SM-like VB $\left(W^{\prime}, Z^{\prime}\right)$ which transforms as $(\mathbf{1}, \mathbf{3}, \mathbf{0})$ under the SM Gauge group is another possibility.

Six Couplings: $\quad\left(g_{\mu \mu}, g_{\tau \tau}, g_{\mu \tau}\right),\left(g_{s s}, g_{b b}, g_{s b}\right)$


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- In addition to the Semi-Leptonic operators required to explain the B-Anomalies the four Fermion are also generated at the Tree Level.
- Additional constraints like $B_{s}-\bar{B}_{s}$ Mixing, $\tau \rightarrow 3 \mu, \tau \rightarrow \ell \nu \bar{\nu}$ come into play.


## VB Triplet Model: Results

- Due to the constraints from $\tau \rightarrow \ell \nu \bar{\nu}$ and $B_{s}$-Mixing the $g_{\tau \tau} \sim \mathcal{O}(0.01-0.1)$ is small, so the NP effect in $b \rightarrow c \tau \bar{\nu}$ are limited.


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- So, we conclude that the VB model is excluded.


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Thanks for your attention!

