

**A NEW ( $\tau$  AND LATTICE-BASED)  
DETERMINATION OF  $|V_{us}|$**

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## CONTEXT (Pre-2016)

- The  $> 3\sigma$  low inclusive FB  $\tau$  FESR  $|V_{us}|$  puzzle

$ V_{us} $	Source
0.2258(9)(?)	3-family unitarity, HT14 $ V_{ud} $
$0.2231(4)_{exp}(7)_{latt}$	$K_{\ell 3}$ , 2+1+1 lattice $f_+(0)$
$0.2250(4)_{exp}(9)_{latt}$	$\Gamma[K_{\mu 2}]/\Gamma[\pi_{\mu 2}]$ , lattice $f_K/f_\pi$
$0.2176(19)_{exp}(10?)_{th}$	Inclusive FB kinematic wt $\tau$ FESR (Passemar CKM14)

- $\tau$  result  $3.6\sigma$  low: lepton flavor non-universality??

[Recall  $\sim 4\sigma$   $R[D^{(*)}]$  discrepancy c.f. SM]

## BASICS: HADRONIC $\tau$ DECAYS IN THE SM

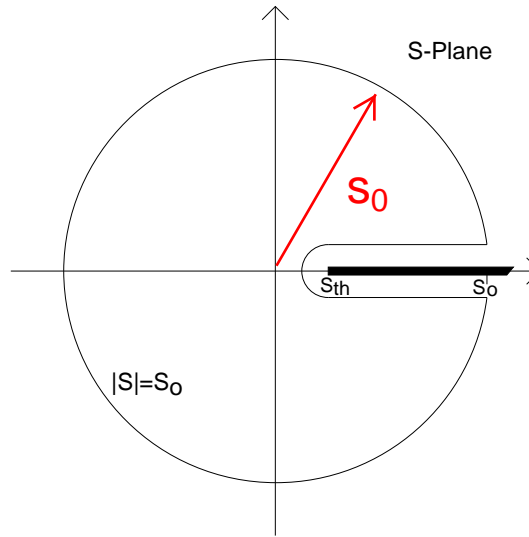
- $R_{ij;V/A} \equiv \Gamma[\tau \rightarrow \nu_\tau \text{ hadrons}_{ij;V/A}(\gamma)] / \Gamma[\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e(\gamma)]$
- With  $y_\tau \equiv s/m_\tau^2$ , flavor  $ij$  decays in SM [Tsai PRD4 (1971) 2821]

$$\frac{dR_{ij;V+A}}{ds} = \frac{12\pi^2 |V_{ij}|^2 S_{EW}}{m_\tau^2} [1 - y_\tau]^2 \tilde{\rho}_{ij;V+A}(s)$$

$$\tilde{\rho}_{ij;V+A}(s) \equiv \left[ (1 + 2y_\tau) \rho_{ij;V+A}^{(J=1)}(s) + \rho_{ij;V+A}^{(J=0)}(s) \right]$$

$$\text{kinematic weight : } w_\tau(y) = (1 - y)^2 (1 + 2y)$$

- Finite energy sum rules (FESRs) (Cauchy's theorem)



$$\int_{s_{th}}^{s_0} ds w(s) \rho(s) = \frac{-1}{2\pi i} \oint_{|s|=s_0} ds w(s) \Pi(s)$$

*expt'l data* *OPE*

- $s_0 \rightarrow \infty$ : generalized dispersion relation

- $R_{ij;V/A}^w(s_0)$ : re-weighted  $R_{ij;V/A}$  analogue

$$R_{ij;V/A}^w(s_0) \sim \int_{th}^{s_0} ds \frac{dR_{ij;V/A}}{ds} \frac{w(s/s_0)}{w_\tau(s/m_\tau^2)}$$

- FB differences:  $\delta R^w(s_0) \equiv \frac{R_{ud;V+A}^w(s_0)}{|V_{ud}|^2} - \frac{R_{us;V+A}^w(s_0)}{|V_{us}|^2}$

- FB FESR: OPE for  $\delta R^w(s_0)$ , input  $|V_{ud}| \Rightarrow$

$$|V_{us}| = \sqrt{\frac{R_{us;V+A}^w(s_0)}{\frac{R_{ud;V+A}^w(s_0)}{|V_{ud}|^2} - [\delta R^w(s_0)]^{OPE}}}$$

Self-consistency:  $|V_{us}|$  independent of  $s_0, w$

- **The conventional implementation** [Gamiz et al. JHEP03(2003)060]

- $s_0 = m_\tau^2$ ,  $w = w_\tau$  only ( $\Rightarrow$  OPE up to  $D = 8$ )
- ASSUME  $D > 4$  OPE small
- Variable  $w$ ,  $s_0$  tests  $\Rightarrow$  self-consistency problems

- **An alternate implementation** [HLMZ17, arXiv:1702.01767]

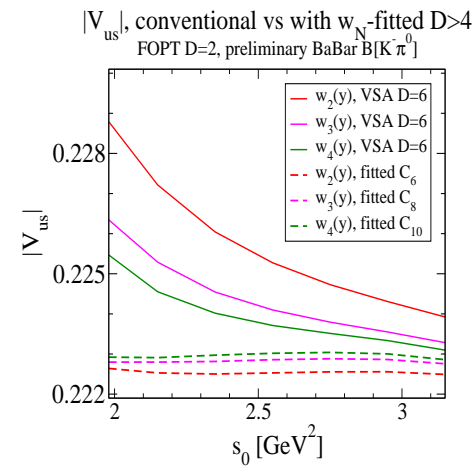
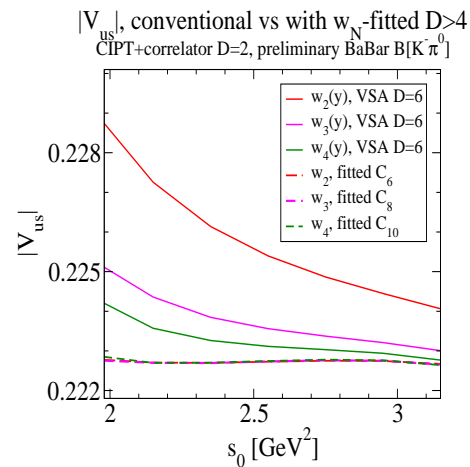
- No  $D > 4$  assumptions: effective condensates  $C_{D>4}$  from fits to data (variable  $s_0$  **required**)
- 3-loop-truncated FOPT  $D = 2$ , standard  $D = 2 + 4$  error estimates [from comparison to lattice]
- $w$ - and  $s_0$ -stability tests

- The alternate FB FESR implementation results

- $|V_{us}| = 0.2229(22)_{exp}(4)_{th}$

- Resolves long-standing low  $\tau$   $|V_{us}|$  puzzle

- Unphysical  $s_0$ -,  $w(y)$ -dependence problems solved



- $|V_{us}|$  increased by  $\sim 0.0020$  with fitted  $C_{D>4}$

- Similar increase from newer  $us$  BF normalizations
- Error budget, 3-weight, Adametz  $B[K^- \pi^0 \nu_\tau]$ , 3-loop-truncated FOPT  $D = 2$  fit

Source	$\delta V_{us} $ ( $w_2$ FESR)	$\delta V_{us} $ ( $w_3$ FESR)	$\delta V_{us} $ ( $w_4$ FESR)
$\delta\alpha_s$	0.00001	0.00004	0.00004
$\delta m_s(2 \text{ GeV})$	0.00017	0.00019	0.00019
$\delta\langle m_s \bar{s}s \rangle$	0.00035	0.00035	0.00035
$\delta(\text{long corr})$	0.00009	0.00009	0.00009
$ud$ exp	0.00027	0.00028	0.00028
<b><math>us</math> exp</b>	<b>0.00226</b>	<b>0.00227</b>	<b>0.00227</b>

- $us$  spectral integrals dominant current error source



- Limitation for near-term improvement:  $\sim 25\%$  higher-multiplicity “residual mode” contribution error (sub-%  $us$  spectral integral error to be fully competitive)

Relative exclusive mode  $R_{us:V+A}^w$  contributions

$Wt$	$s_0$ [GeV <sup>2</sup> ]	$K$	$K\pi$	$K\pi\pi$ (B-factory)	Residual
$w_2$	2.15	0.496	0.426	0.062	0.010
	3.15	0.360	0.414	0.162	0.065
$w_3$	2.15	0.461	0.446	0.073	0.019
	3.15	0.331	0.415	0.182	0.074
$w_4$	2.15	0.441	0.456	0.082	0.021
	3.15	0.314	0.411	0.194	0.081

## A LATTICE+ $\tau$ -BASED ALTERNATIVE

- Work with J. Hudspith, R. Lewis (York) and T. Izubuchi, H. Ohki, C. Lehner + ... (RBC/UKQCD)
- Basic idea: generalized dispersion relations for products of combination  $\tilde{\Pi}$  of  $J = 0, 1$   $us$   $V+A$  polarizations with weights having poles at Euclidean  $Q^2$ 
  - $\tilde{\Pi}(Q^2)$ : polarization sum with spectral function  $\tilde{\rho}(s)$  (experimental  $dR_{us;V+A}/ds$ )
  - Theory: Lattice  $us$  2-point function data (no OPE)
  - Weights tunable, allow suppression of larger-error, higher-multiplicity  $us$  spectral contributions

## More on the lattice-inclusive $us\ \tau$ approach

- $|V_{us}|^2 \tilde{\rho}_{us;V+A}(s)$  from experimental  $dR_{us;V+A}/ds$

$$\tilde{\rho}_{us;V+A}(s) \equiv \left(1 + 2\frac{s}{m_\tau^2}\right) \rho_{us;V+A}^{(J=1)}(s) + \rho_{us;V+A}^{(J=0)}(s)$$

(no continuum  $us$   $J = 0$  subtraction required)

- Associated (kinematic-singularity-free) polarization

$$\tilde{\Pi}_{us;V+A}(Q^2) \equiv \left(1 - 2\frac{Q^2}{m_\tau^2}\right) \Pi_{us;V+A}^{(J=1)}(Q^2) + \Pi_{us;V+A}^{(J=0)}(Q^2)$$

- $\tilde{\rho}_{us;V+A}(s) \sim s$  as  $s \rightarrow \infty$

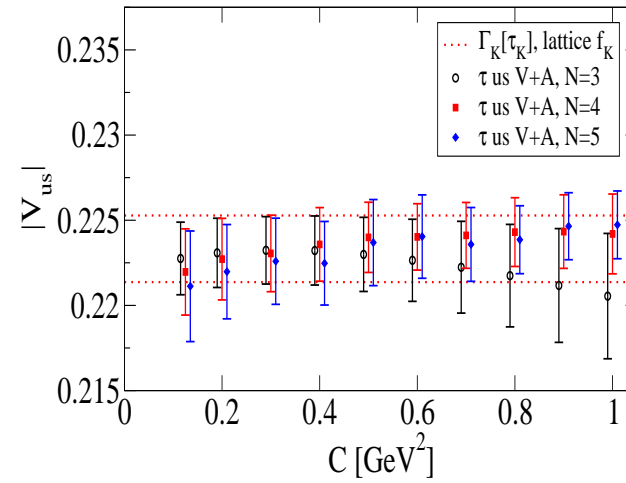
- For weights  $w_N(s) \equiv \frac{1}{\prod_{k=1}^N (s+Q_k^2)}$ ,  $Q_k^2 > 0$ ,  $N \geq 3$ , obtain convergent, unsubtracted 'dispersion relation'

$$\int_{th}^{\infty} ds w_N(s) \tilde{\rho}_{us;V+A}(s) = \sum_{k=1}^N \frac{\tilde{\Pi}_{us;V+A}(Q_k^2)}{\prod_{j \neq k} (Q_j^2 - Q_k^2)}$$

- Lattice data for  $\tilde{\Pi}_{us;V+A}(Q_k^2)$  on RHS
- LHS from experimental  $dR_{us;V+A}/ds$ , up to  $|V_{us}|^2$
- $w_N(s)$ : rapid fall-off if all  $Q_k^2 < 1 \text{ GeV}^2$   
 $\Rightarrow$   **$K$ ,  $K\pi$  dominate LHS, near-endpoint multi-particle,  $s > m_\tau^2$  contributions strongly suppressed**
- Optimization: increasing  $\{Q_k^2\}$  decreases RHS lattice error, increases LHS experimental error

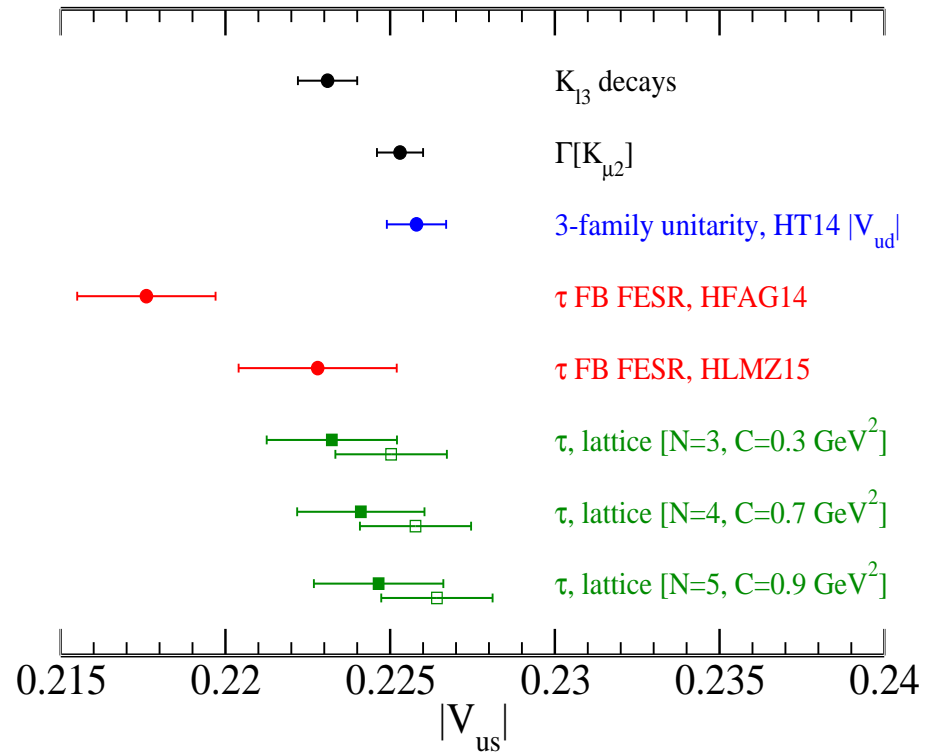
- *PRELIMINARY inclusive lattice  $us$   $V+A$  results*
  - $N = 4, C = 0.7 \text{ GeV}^2, \Gamma[K_{\mu 2}]: |V_{us}| = 0.2258(10)_{exp}(13)_{th}$
  - $N = 4, C = 0.7 \text{ GeV}^2, B_K^T: |V_{us}| = 0.2241(14)_{exp}(13)_{th}$
  
- Advantages of lattice-based vs. FB FESR approach
  - $K$  essentially saturates  $J = 0, A$  contribution  $\Rightarrow |V_{us}|$  determinations possible with or without  $K$  pole
  - Reduced expt'l error *without theory error blowup*
  - Theory side: lattice in place of OPE  $\Rightarrow$  theory errors systematically improvable

- Stability e.g.



$N = 3, 4, 5, B_K^\tau$  choice for  $K$  pole contribution

## Comparison to $|V_{us}|$ from other sources



## SUMMARY

- Old  $3\sigma$  low inclusive FB  $\tau$  FESR  $|V_{us}|$  problem resolved
  - Alternate, no-assumptions implementation:  $|V_{us}|$  higher by  $\sim 0.0020$ , compatible with other determinations
  - Near-term improvements feasible through improvements in  $us$  exclusive mode BFs
  - Highly favorable theoretical error situation
  - However, for competitive  $|V_{us}|$  need improvements to old ALEPH higher-multiplicity, low-statistics data [unlikely in the near-term]



- Advantage of new lattice-inclusive  $us V + A \tau$  approach
  - Theory:
    - \* Lattice in place of OPE; no  $us J = 0$  subtraction; improvement through increased statistics
    - \* Parasitic on lattice  $a_\mu$  effort (major effort in lattice community)
  - Spectral integrals:
    - \* Theory errors still small for weights strongly suppressing higher multiplicity contributions
    - \* Strong  $K, K\pi$  dominance of spectral integral
    - \* Significant experimental improvements possible through just improved  $K\pi$  BFs, distributions